



High-beta Cavity Design

K.Saito

**High Energy Accelerator Research Organization (KEK),
Accelerator Lab**

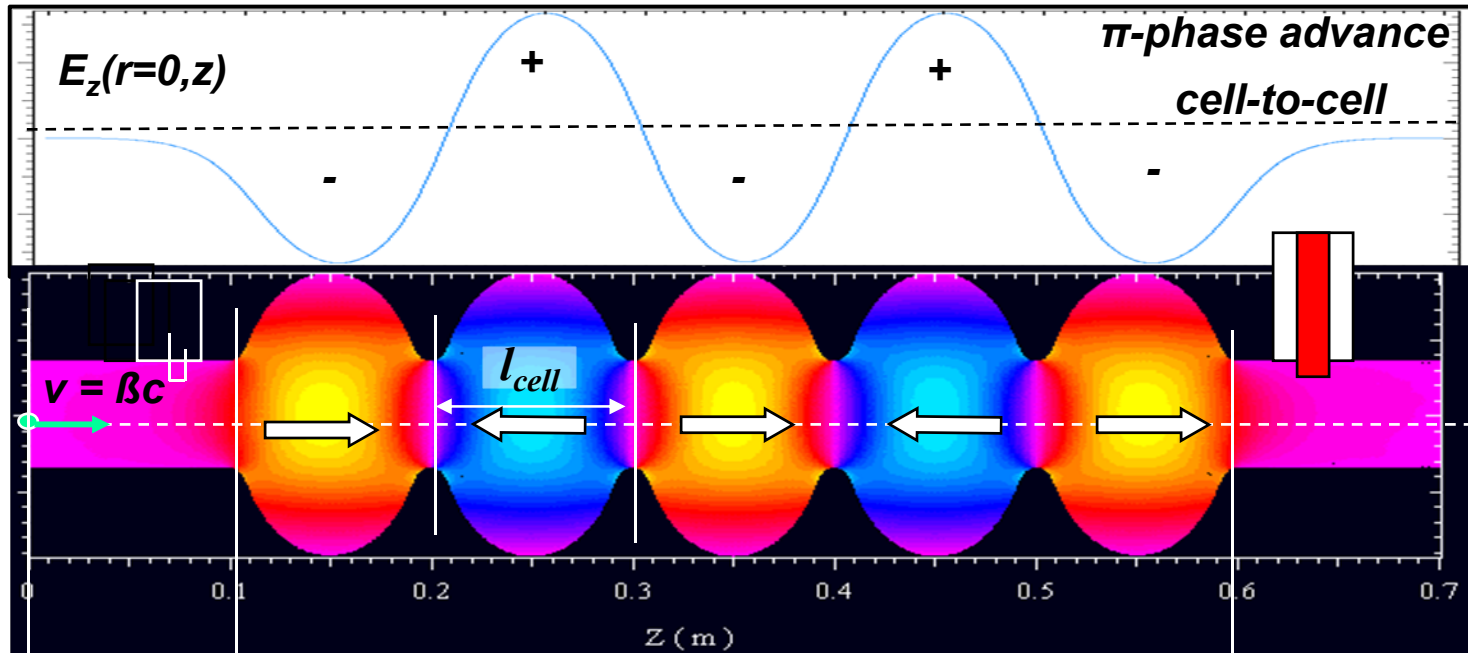
- 1. SW Operation with SRF RF Cavity**
- 2. Pill Box Cavity**
- 3. Figure of merits of SRF Cavity Design**
- 4. Criteria for Multi-Cell Structures**
- 5. Example of SRF High-beta Cavities**

Acknowledgement:

In this lecture, I used many slides by Dr. J.Sektowitz@ DESY, which he made nice lecture in China. I would like to thank for him very much.

1. Sanding Wave (SW) Operation and Standing Wave in SRF Cavity

SW Scheme in SRF Cavity Operation



TM_{010}
 π -mode

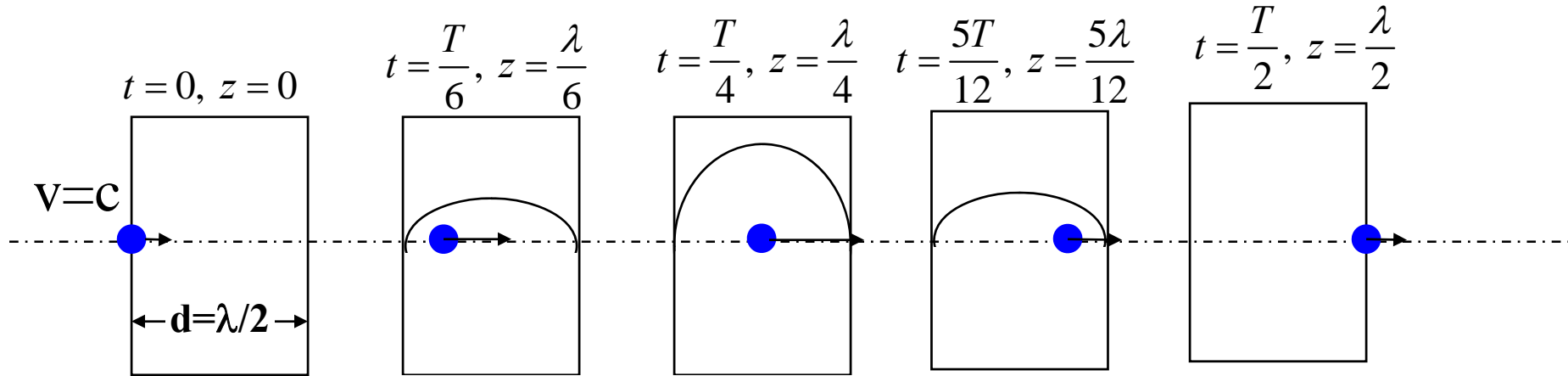
Standing wave (CW) is used in SRF cavity acceleration !

Synchronic acceleration and max of $(R/Q)_{acc} \leftrightarrow I_{active} = NI_{cell} = Nc\beta/(2f)$ and the injection takes place at an optimum phase φ_{opt} which ensures that particles will arrive at the mid-plane of the first cell when E_{acc} reaches its maximum (+q passing to the right) or minimum (-q passing to the right).

$$L_{cell} = \frac{\beta \cdot \lambda}{2}, \quad \lambda = \frac{c}{f} \quad \beta = v/c$$

for electron $\beta=1$, proton $\beta < 1$

Transit Time Factor Due to SW Operation



$$V = \left| \int_0^d E_z(r=0, z) e^{i\omega t} dz \right| = \left| \int_0^d E_z(r=0, z) e^{i\omega \frac{z}{c}} dz \right| = E_0 \left| \int_0^d e^{i\omega \frac{z}{c}} dz \right| = E_0 d \frac{\sin\left(\frac{\omega d}{2c}\right)}{\frac{\omega d}{2c}} = E_0 d \cdot T$$

$z = c \cdot t$

T : Transit time factor

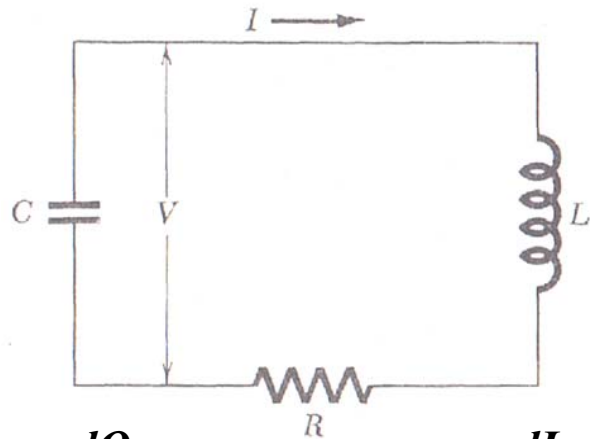
$$T = \frac{2}{\pi} = 0.637 \quad (\text{for Pill Box Cavity})$$

$$E_{acc} \equiv \frac{V}{d} = E_0 T$$

Acceleration efficiency is automatically reduced by ~ 40% in the SW scheme.

2. Pill Box Cavity

LC Equivalent Circuit of Cavity



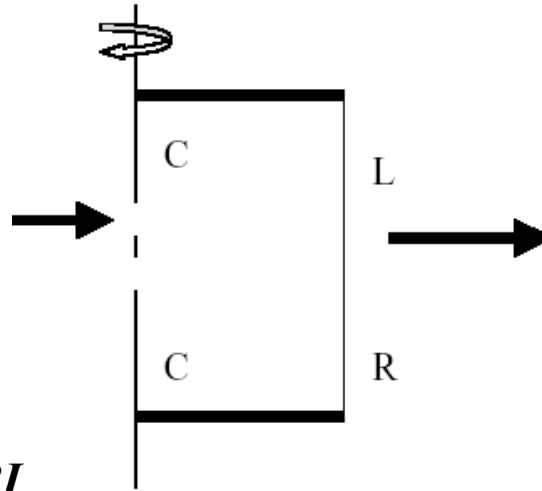
$$I = -\frac{dQ}{dt}, \quad Q = CV, \quad V = L\frac{dI}{dt} + RI$$

$$\frac{d^2V}{dt^2} + \left(\frac{R}{L}\right)\frac{dV}{dt} + \left(\frac{1}{LC}\right)V = 0, \quad V(t) = V_0 \exp(-\alpha + i\omega)t$$

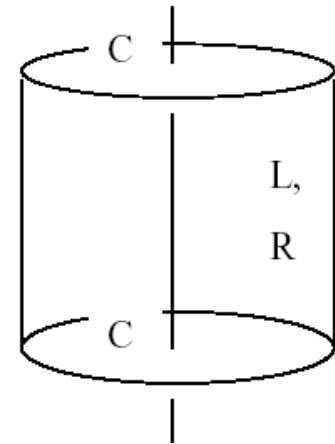
$$(-\alpha + i\omega)^2 + (-\alpha + i\omega)\left(\frac{R}{L}\right) + \left(\frac{1}{LC}\right) = 0,$$

$$\alpha = \frac{R}{2L}, \quad \omega^2 = \frac{1}{LC} - \frac{R^2}{4L^2}$$

$$R \ll L, \quad \omega_0^2 = \frac{1}{LC} \Rightarrow f = \frac{1}{2\pi\sqrt{LC}}$$



Pill Box Cavity



Q-value of the circuit

$$Q \equiv \omega \cdot \frac{\text{Stored Energy}}{\text{Powerloss/sec}} = \omega \cdot \frac{P(t)}{dP(t)/dt} = \omega \cdot \frac{L}{R}$$

$$= \frac{\omega}{2\alpha}$$

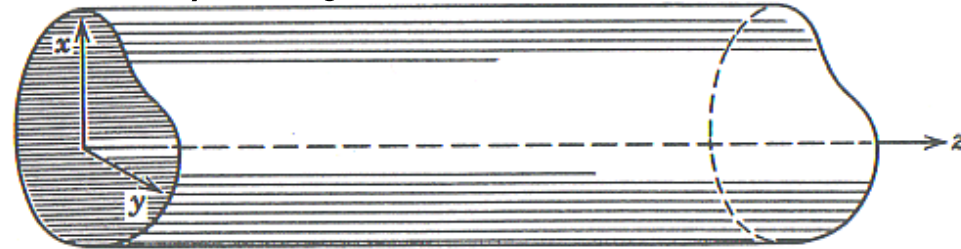
Q : proportional to 1/R

Electro-magnetic field in a waveguide

Maxwell equations in a waveguide

$$\nabla \times \mathbf{E} = i \frac{\omega}{c} \mathbf{B}, \quad \nabla \cdot \mathbf{B} = 0, \quad \nabla \times \mathbf{B} = -i \mu \varepsilon \frac{\omega}{c} \mathbf{E}, \quad \nabla \cdot \mathbf{E} = 0, \quad \rho = 0, \quad \mathbf{j} = 0$$

$$\left(\nabla^2 + \mu \varepsilon \frac{\omega^2}{c^2} \right) \begin{Bmatrix} \mathbf{E} \\ \mathbf{B} \end{Bmatrix} = 0,$$



$$\mathbf{E}(x, y, z, t) = \mathbf{E}(x, y) \exp(\pm ikz - i\omega t), \quad k: \text{wavevector},$$

$$\mathbf{B}(x, y, z, t) = \mathbf{B}(x, y) \exp(\pm ikz - i\omega t),$$

$$\left[\nabla_t^2 + \left(\varepsilon \mu \frac{\omega^2}{c^2} - k^2 \right) \right] \begin{Bmatrix} \mathbf{E} \\ \mathbf{B} \end{Bmatrix} = 0, \quad \nabla_t^2 \equiv \nabla^2 - \frac{\partial^2}{\partial z^2}, \quad \mathbf{E} = E_z \mathbf{e}_z + \mathbf{E}_t, \quad \mathbf{B} = B_z \mathbf{e}_z + \mathbf{B}_t$$

$$\mathbf{B}_t = \frac{1}{\left(\varepsilon \mu \frac{\omega^2}{c^2} - k^2 \right)} \left[\nabla_t \left(\frac{\partial B_z}{\partial z} \right) + i \varepsilon \mu \frac{\omega}{c} \mathbf{e}_z \times \nabla_t E_z \right],$$

$$\mathbf{E}_t = \frac{1}{\left(\varepsilon \mu \frac{\omega^2}{c^2} - k^2 \right)} \left[\nabla_t \left(\frac{\partial E_z}{\partial z} \right) - i \frac{\omega}{c} \mathbf{e}_z \times \nabla_t B_z \right]$$

Homework I.
Lead this formula.

TM mode: $B_z = 0, E_z \neq 0$

TE mode: $B_z \neq 0, E_z = 0$

TM- Modes

TM-mode: $B_z = 0, E_z \neq 0 \rightarrow$ Can accelerate beam
Beam

$$\mathbf{B}_t = \frac{i\epsilon\mu \frac{\omega}{c}}{\left(\epsilon\mu \frac{\omega^2}{c^2} - k^2\right)} [\mathbf{e}_z \times \nabla_t E_z],$$

$$\mathbf{E}_t = \frac{1}{\left(\epsilon\mu \frac{\omega^2}{c^2} - k^2\right)} \nabla_t \left(\frac{\partial E_z}{\partial z} \right),$$

$$\left[\nabla_t^2 E_z + \left(\epsilon\mu \frac{\omega^2}{c^2} - k^2\right) \right] E_z = 0,$$

Solve the eigenvalue problem,
get k and E_z

Boundary condition $\mathbf{E}_\perp|_S = 0$ ($\because \mathbf{n} \times \mathbf{E} = 0$ on the surface of perfect conductor)

$$\frac{\partial \mathbf{B}_z}{\partial \mathbf{n}}|_S = 0 \quad (\because \mathbf{n} \cdot \mathbf{B} = 0 \quad \text{on the surface,})$$

but automatically satisfied by the TM mode condition)

Similar for TE modes, which kicks beam.....

Eigen-value problem

$\psi(x, y) = E_z(x, y)$ for TM-mode or $B_z(x, y)$ for TE-mode

$$(\vec{\nabla}_t^2 + \gamma^2)\psi = 0, \quad \psi|_S = 0 \text{ (for TM-Modes)} \quad \text{or} \quad \frac{\partial}{\partial n}\psi|_S = 0 \text{ (for TE-Modes)}$$

$$\gamma^2 = \epsilon\mu \cdot \frac{\omega^2}{c^2} - k^2 \geq 0$$

From the boundary condition,

$$\gamma^2 = \gamma_\lambda^2, \quad \psi = \psi_\lambda \quad (\lambda = 1, 2, \dots) \quad \Rightarrow \quad k_\lambda^2 = \epsilon\mu \frac{\omega^2}{c^2} - \gamma_\lambda^2$$

If $\omega < c \frac{\gamma_\lambda}{\sqrt{\epsilon\mu}}$, then k_λ is an imaginal number.

The wave is damped/trapped in the waveguide, if surface resistance large/small.

$$\omega_\lambda = c \frac{\gamma_\lambda}{\sqrt{\epsilon\mu}} \dots \text{cutoff frequency}$$

When $\omega \geq \omega_\lambda$, wave number k_λ is a real number,
then the wave can propagate into the waveguide.

$TM_{m,n,p}$ – mode

$$E_z = E_0 \cos(kz) J_m \left(\frac{\rho_{m,n}}{a} r \right) \exp(-im\theta), \quad B_z = 0$$

$$E_r = \frac{iE_0 p \pi}{\gamma_{m,n,p}} \cos\left(\frac{p\pi}{d} z\right) \frac{\partial J_m(\rho)}{\partial \rho} \exp(-im\theta), \quad B_r = -\frac{E_0 m \varepsilon \mu \omega_{m,n,p}}{\gamma_{m,n,p}} \cos(kz) J_m \left(\frac{\rho_{m,n}}{a} r \right) \exp(-im\theta)$$

$$E_\theta = \frac{E_0 m p \pi}{\gamma_{m,n,p}^2 d c} \cos\left(\frac{p\pi}{d} z\right) J_m \left(\frac{\rho_{m,n}}{a} r \right) \exp(-im\theta), \quad B_\theta = \frac{iE_0 \varepsilon \mu \omega_{m,n,p}}{\gamma_{m,n,p} c} \cos(kz) \exp(-im\theta) \frac{\partial J_m(\rho)}{\partial \rho}$$

resonance frequency $\omega_{m,n,p} \Rightarrow f_{m,n,p} = \frac{c}{\sqrt{\varepsilon \mu}} \sqrt{\frac{\rho_{m,n}^2}{a^2} + \frac{p^2 \pi^2}{d^2}}$

TE_{mnp} Modes

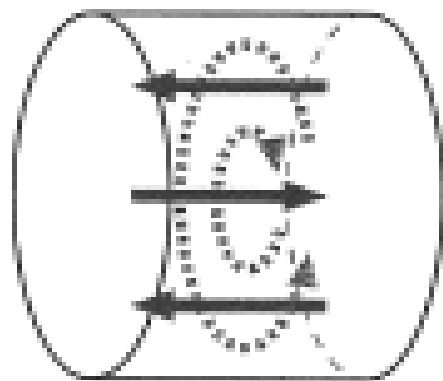
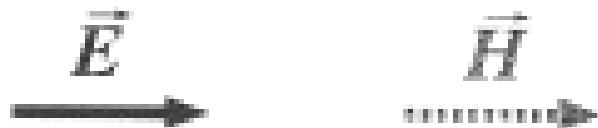
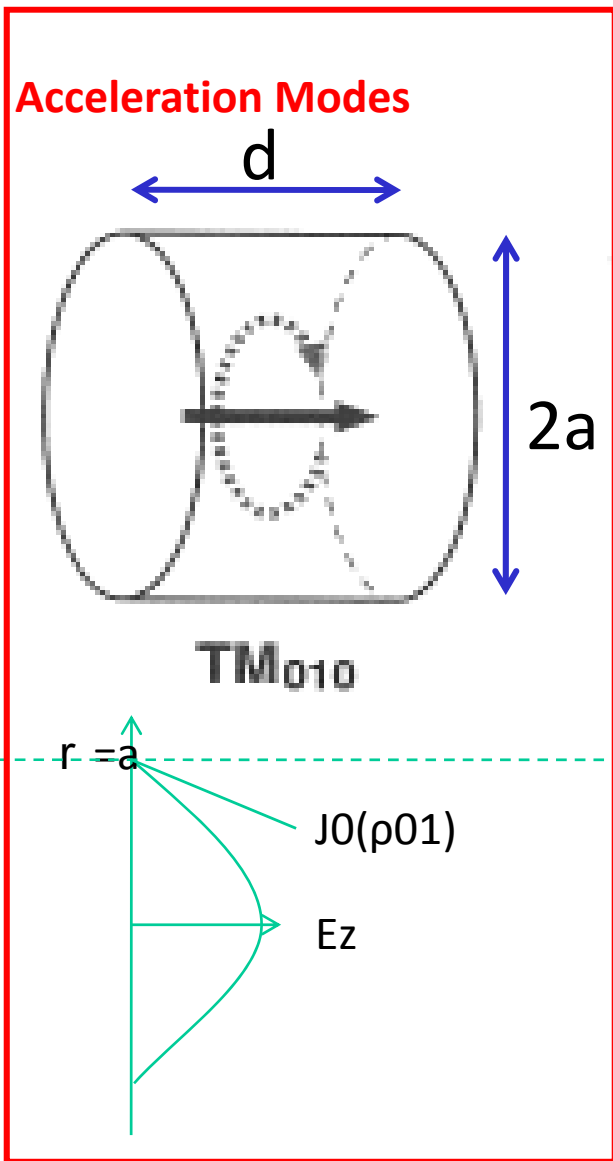
$$E_z = 0 \quad H_z = E_0 J_m \left(\frac{\rho'_{mn}}{a} r \right) \cos(m\theta) \sin\left(\frac{p\pi z}{d}\right)$$

$$E_r = \frac{i\omega \varepsilon}{k^2} \frac{m}{r} E_0 J_m \left(\frac{\rho'_{mn}}{a} r \right) \sin(m\theta) \sin\left(\frac{p\pi z}{d}\right) \quad H_r = \frac{1}{k^2} \frac{p\pi}{d} E_0 J_m \left(\frac{\rho'_{mn}}{a} r \right) \cos(m\theta) \cos\left(\frac{p\pi z}{d}\right)$$

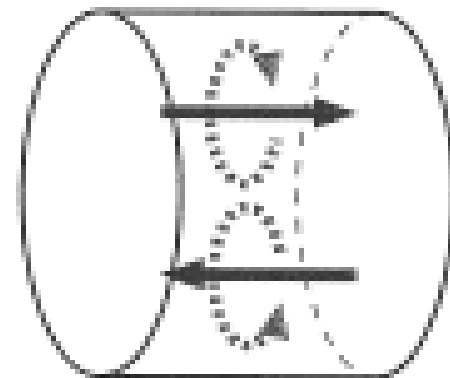
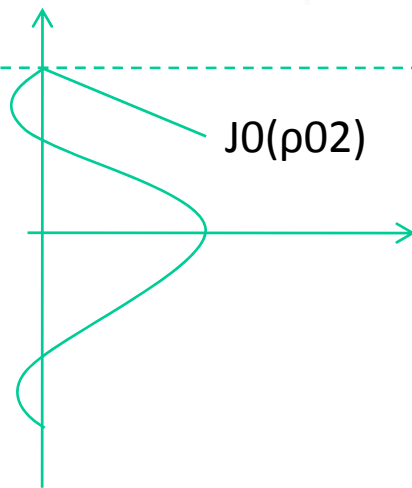
$$E_\theta = \frac{i\omega \varepsilon}{k^2} \frac{\rho'_{mn}}{a} E_0 J_m \left(\frac{\rho'_{mn}}{a} r \right) \cos(m\theta) \sin\left(\frac{p\pi z}{d}\right) \quad H_\theta = -\frac{1}{k^2} \frac{p\pi}{d} \frac{m}{r} E_0 J_m \left(\frac{\rho'_{mn}}{a} r \right) \sin(m\theta) \cos\left(\frac{p\pi z}{d}\right)$$

Resonance frequency $\omega^2 \varepsilon_0 \mu_0 = \left(\frac{\rho'_{mn}}{a} \right)^2 + \left(\frac{p\pi}{d} \right)^2 \Rightarrow f = \frac{c}{2\pi} \sqrt{\left(\frac{\rho'_{mn}}{a} \right)^2 + \left(\frac{p\pi}{d} \right)^2}$ 10

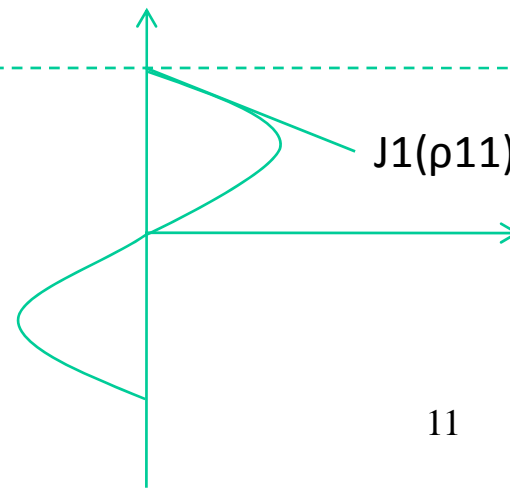
TM_{mnp} Modes



TM₀₂₀

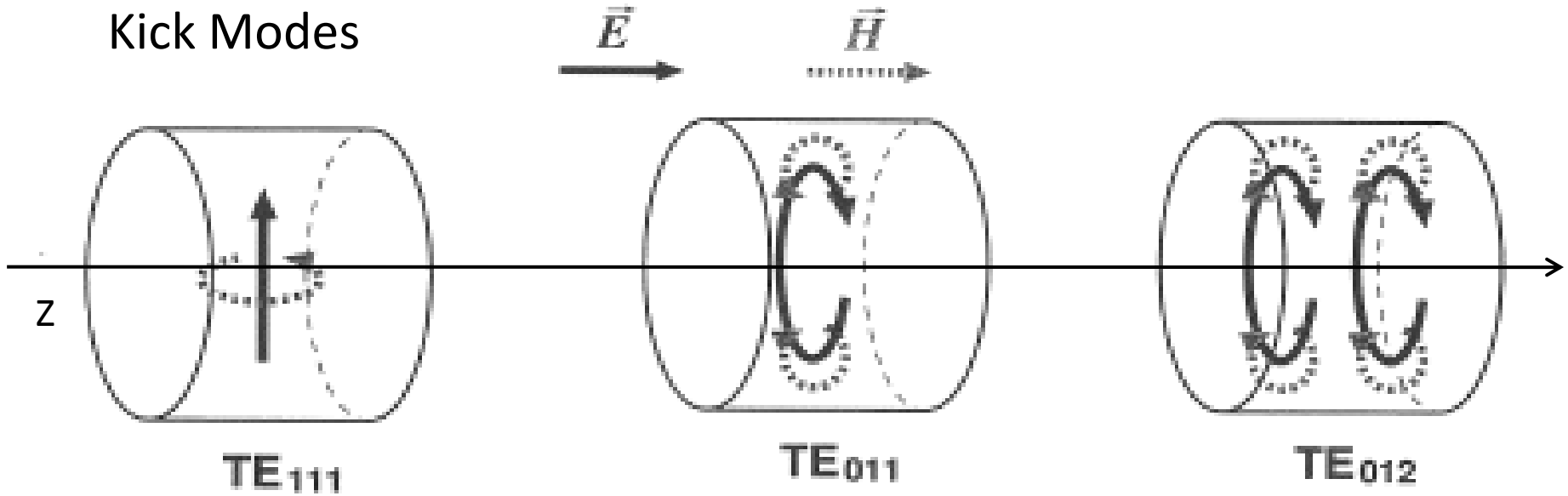


TM₁₁₀



TE_{mnp} Modes

Kick Modes



4. Figure of merits of Cavity Design

Figure of Merits of The RF Cavity Design

	Items	Notation
1)	Surface resistance	$R_s [\Omega]$
2)	RF dissipation	$P_{\text{loss}} [\text{W}]$
3)	Acceleration gradient	$V [\text{V}]$
3)	Unloaded Q	Q_o
4)	Geometrical factor	$\Gamma [\Omega]$
5)	Shunt Impedance	$R_{\text{sh}} [\Omega]$
6)	R over Q	$(R/Q) [\Omega]$
7)	Acceleration gradient	$E_{\text{acc}} [\text{V/m}]$
8)	Ratio of the surface peak electric field vs E_{acc}	E_p/E_{acc} [Oe/(MV/m)]
9)	Ratio of the surface magnetic field vs E_{acc}	H_p/E_{acc} $k_{//}, k_{\perp}$
10)	HOM loss factor	
11)	Field flatness factor	a_{ff}
12)	Cell to cell coupling	k_{CC}

Surface Resistance

Maxwell Equations.

$$\nabla \cdot \vec{B} = 0, \quad \nabla \times \vec{E} + \mu \frac{\partial \vec{H}}{\partial t} = 0$$

$$\nabla \cdot \vec{D} = 0, \quad \nabla \times \vec{H} - \varepsilon \frac{\partial \vec{E}}{\partial t} - \sigma \vec{E} = 0$$

$$\vec{J} = \sigma \vec{E} \quad (\text{Ohm's Law})$$

$$\vec{E}(\vec{x}, t) = \vec{E}_\ell(\vec{x}, t) + \vec{E}_t(\vec{x}, t),$$

$$\vec{H}(\vec{x}, t) = \vec{H}_\ell(\vec{x}, t) + \vec{H}_t(\vec{x}, t)$$

For the transversers,

$$\text{Plane wave : } \vec{E}_t(\vec{x}, t) = \vec{E}_t(0) \cdot \exp(i\vec{k} \cdot \vec{x} - \omega t)$$

$$\vec{H}_t(\vec{x}, t) = \frac{1}{\mu\omega} [\vec{k} \times \vec{E}_t(\vec{x}, t)],$$

$$[k^2 - (\varepsilon\mu\omega^2 + i\mu\omega\sigma)] \begin{Bmatrix} \vec{E}_t(\vec{x}, t) \\ \vec{H}_t(\vec{x}, t) \end{Bmatrix} = 0$$

Rs: Surface resistance

$$Z \equiv R_s + iX_s \equiv \left. \frac{E_t}{H_t} \right|_{\text{Surface}} = \frac{\mu\omega}{k}$$

$$R_s = \sqrt{\frac{\mu\omega}{2\sigma}} = \frac{1}{\sigma} \sqrt{\frac{\mu\sigma\omega}{2}} = \frac{1}{\sigma\delta}$$

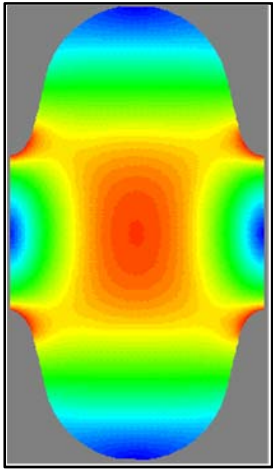
$$P_{\text{loss}} = \frac{1}{2} R_s \cdot \int_S H_s^2 dS$$

Homework II

**Lead the formula of Rs
for good electric conductor**

Charismatic RF Parameter of Cavity

$U \equiv$ stored energy, P_{loss} : Surface RF Heating



Unloaded Q_0

$$U = \frac{1}{2} \cdot \epsilon \int_V E^2 dV = \frac{1}{2} \cdot \mu \int_V H^2 dV$$

$$P_{loss} = \frac{1}{2} \cdot R_S \int_S H_S^2 dS$$

$$Q_0 \equiv \frac{\omega U}{P_{loss}}$$

For acceleration mode (TM010), the higher Q_0 is better.

Acceleration Voltage

$$V[\text{V}] = \int_0^{L_{eff}} E dz \quad \text{on the beam axis.}$$

Shunt Impedance

$$R_{sh} [\Omega] \equiv \frac{V^2}{P_{loss}}$$

This means the efficiency of the acceleration.

Charismatic RF Parameter of Cavity, continued.

(R/Q)

$$(R/Q) \equiv (R_{sh}/Q_O) = \frac{V^2}{\omega U}$$

This means how much energy is concentrated on beam axis.

This does not depend on material but only cavity shape.

This means the goodness of the cavity shape.

Geometrical factor

$$\Gamma[\Omega] \equiv Q_O \cdot R_S = \frac{\omega \int_V H^2 dV}{\frac{1}{2} \int_S H_S^2 dS}$$

When you know the Γ using a computer code, you can calculate the surface resistance R_S from the measured Q_O value.

Γ is about 270Ω for $\beta=1$ cavities with elliptical or spherical shape.

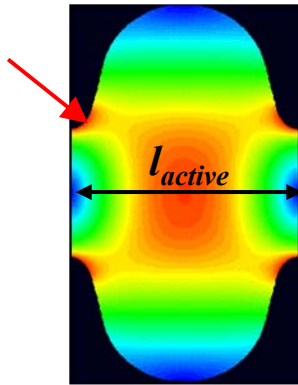
Charismatic RF Parameter of Cavity

Gradient

$$E_{acc} = \frac{\sqrt{(R/Q)}}{L_{eff}} \cdot \sqrt{P_{loss} \cdot Q_0} = Z \cdot \sqrt{P_{loss} \cdot Q_0}$$

E_{peak} on the surface

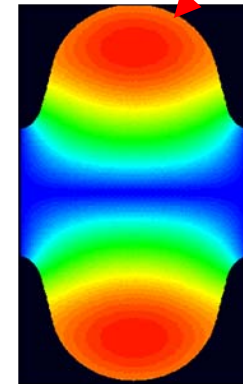
$$\frac{E_{peak}}{E_{acc}}$$



Contour plot of /E/

B_{peak} on the surface

$$\frac{B_{peak}}{E_{acc}}$$



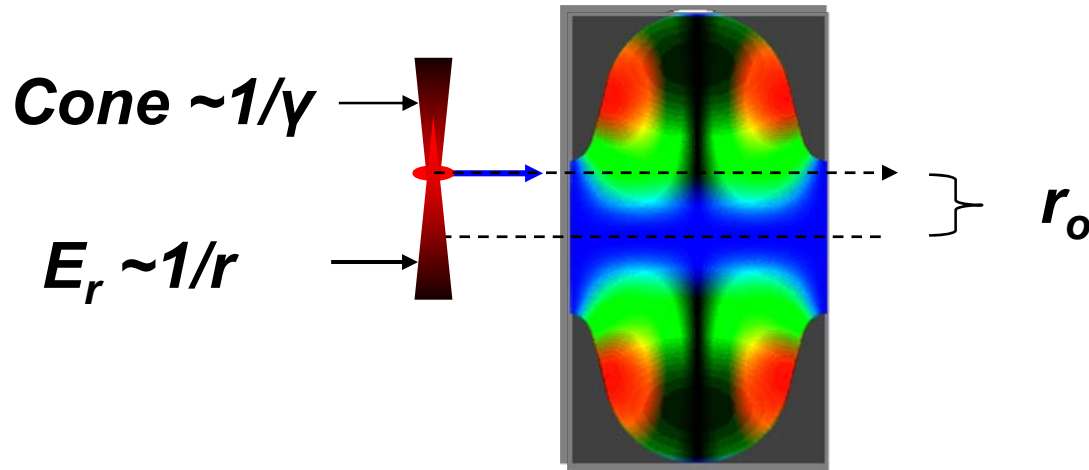
Contour plot of /B/

Smaller value is preferred against electron emission phenomenon.

Ratio shows the limit in E_{acc} due to the break-down of superconductivity (Nb ~2000 Oe).

Longitudinal and Transverse Loss Factors

When ultra relativistic point charge q passes **empty cavity**,



HOM modes are excited in the cavity.

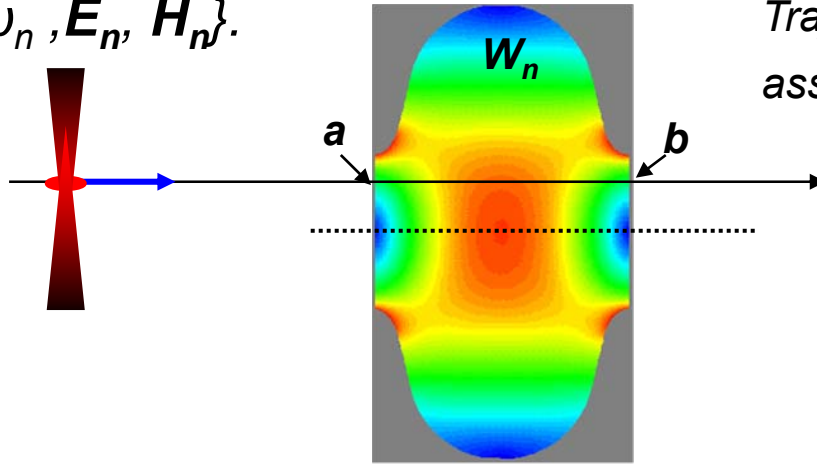
If those are damped fast enough, the following beam feel them and can be kicked.

Thus the beam can be degraded.

$(R/Q)_n$ of n-th HOM

HOM nth $(R/Q)_n$, a “measure” of the energy exchange between point charge and mode n .

mode $n : \{\omega_n, \mathbf{E}_n, \mathbf{H}_n\}$.



Trajectory of the point charge q , assumed here to be a straight line.

$$V_n = \sqrt{\left(\int_{z_a}^{z_b} E_{n,z} \sin\left(\frac{\omega_n}{\beta c} (z - z_a)\right) dz \right)^2 + \left(\int_{z_a}^{z_b} E_{n,z} \cos\left(\frac{\omega_n}{\beta c} (z - z_a)\right) dz \right)^2}$$

$$(R/Q)_n \equiv \frac{V_n^2}{\omega_n W_n}$$

Longitudinal and Transverse Loss Factors

The amount of energy lost by charge q to the cavity is:

J.Sekutwitz's Slide

$$\Delta U_q = k_{\parallel} \cdot q^2 \quad \text{for monopole modes (max. on axis)}$$

$$\Delta U_q = k_{\perp} \cdot q^2 \quad \text{for non monopole modes (off axis)}$$

where k_{\parallel} and $k_{\perp}(\mathbf{r})$ are loss factors for the monopole and transverse modes respectively.

The induced **E-H field (wake)** is a superposition of cavity eigenmodes (monopoles and others) having the $\mathbf{E}_n(\mathbf{r}, \varphi, \mathbf{z})$ field along the trajectory.

For individual mode n and point-like charge:

$$k_{\parallel, n}^p = \frac{\omega_n \cdot (R/Q)_n}{4}$$

Similar for other loss factors.....

These harmful HOM power should be take out from the cavity. 21
Dr. Noguchi will make a lecture about the HOM coupler design.

Field Flatness Factor a_{ff}

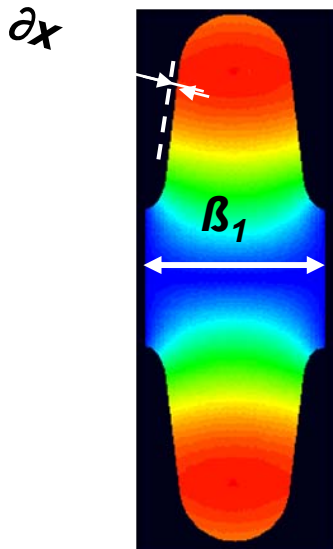
J.Sekutwitz's Slide

Field flatness factor for elliptical cavities with arbitrary $\beta=v/c$

$$\frac{\Delta A}{A} = a_{ff} \cdot \frac{\Delta f}{f}$$

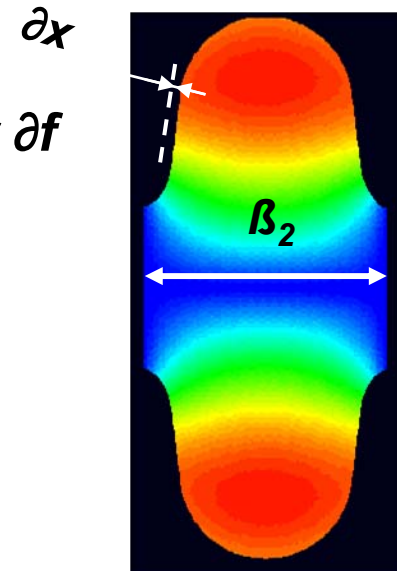
$$a_{ff} = \frac{N^2}{k_{cc} \cdot \beta}$$

This is an empirical correction, based on intuition.



The same error ∂x causes bigger ∂f when a cell is thinner

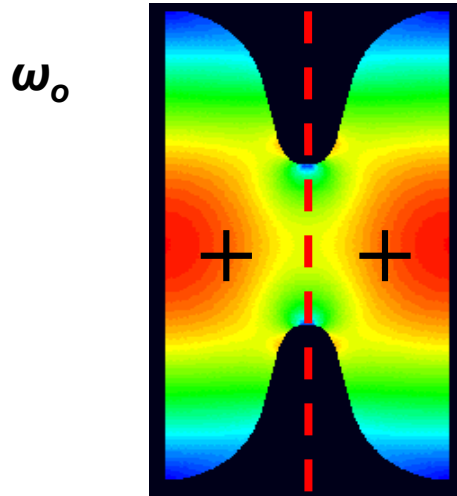
$$\frac{\partial f_1}{\partial f_2} = \frac{\partial x \cdot \beta_2}{\partial x \cdot \beta_1}$$



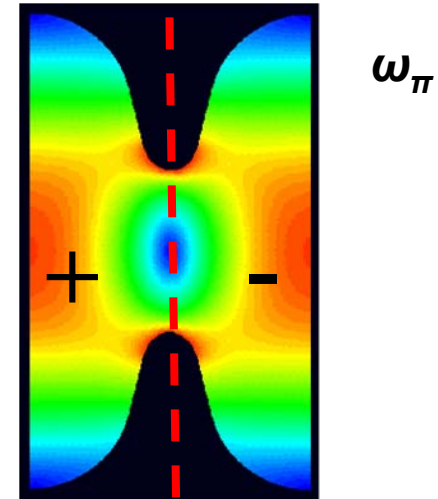
Cells which geometric $\beta < 1$ are more sensitive to shape errors

Cell to cell coupling k_{CC}

The last parameter, relevant for multi-cell accelerating structures, is the coupling k_{cc} between cells for the accelerating mode pass-band (Fundamental Mode pass-band).



no E_r (in general transverse E field) component at the symmetry plane



no H_ϕ (in general transverse H field) component at the symmetry plane

The normalized difference between these frequencies is a measure of the Poynting vector (energy flow via the coupling region)

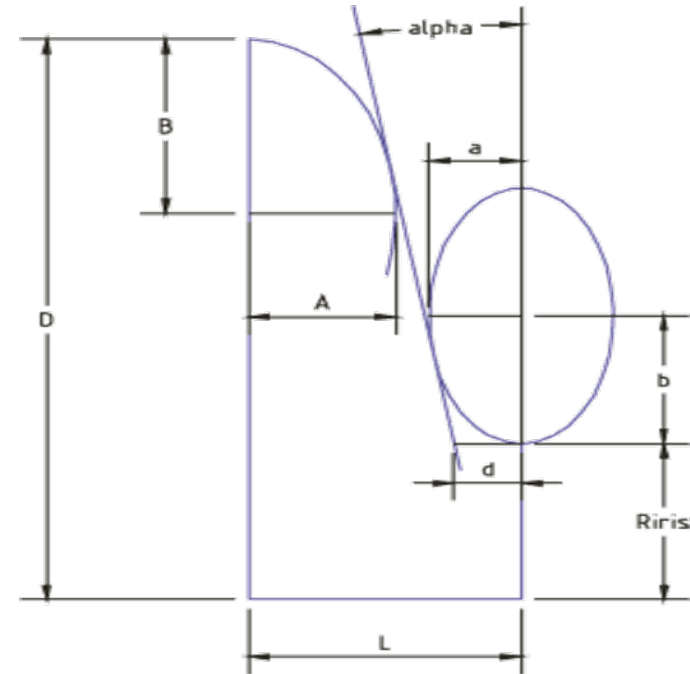
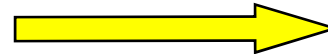
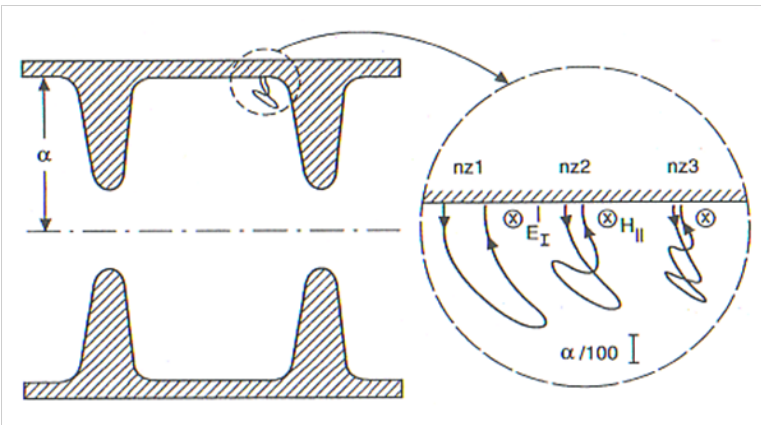
$$k_{cc} = \frac{\omega_\pi - \omega_0}{\omega_\pi + \omega_0} \cdot 2$$

Spherical/Elliptical Shape in SRF cavity design

Design of SRF cavity is usually done with spherical or Elliptical shapes.

Choose spherical shape to reduce multipacting

Multipacting



*R. Parodi (1979) presented first **spherical** C-band cavity with much less multipacting barrier than other cavities at that time.*

*P. Kneisel (early 80's) proposed for the DESY experiment the **elliptical shape** of 1 GHz cavity preserving good performance of the spherical one and stiffer mechanically.*

The RF cavity design tool will be given by Dr. U. van Rienen In this tutorial.

Optimization of Cell Shape

We begin with inner cells design because these cells “dominate” parameters of a multi-cell superconducting accelerating structure.

RF parameters summary:

J.Sekutwitz's Slide

FM : $(R/Q), G, E_{peak}/E_{acc}, B_{peak}/E_{acc}, k_{cc}$

HOM : k_{\perp}, k_{\parallel} .

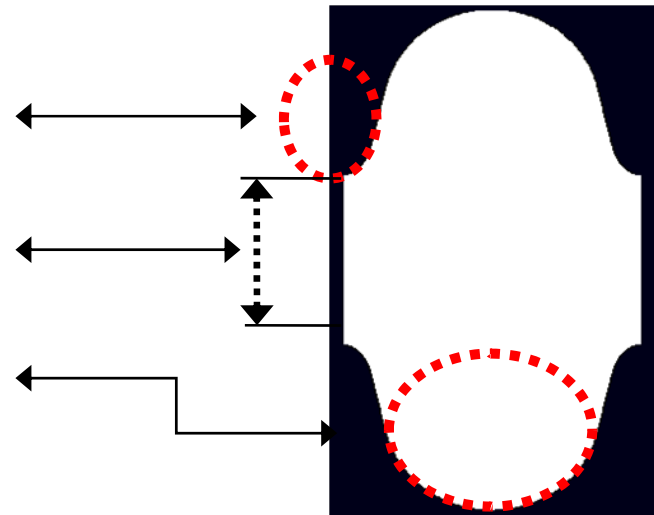
There are 7 parameters we want to optimize for a inner cell

Geometry :

iris ellipsis : half-axis h_r, h_z

iris radius : r_i

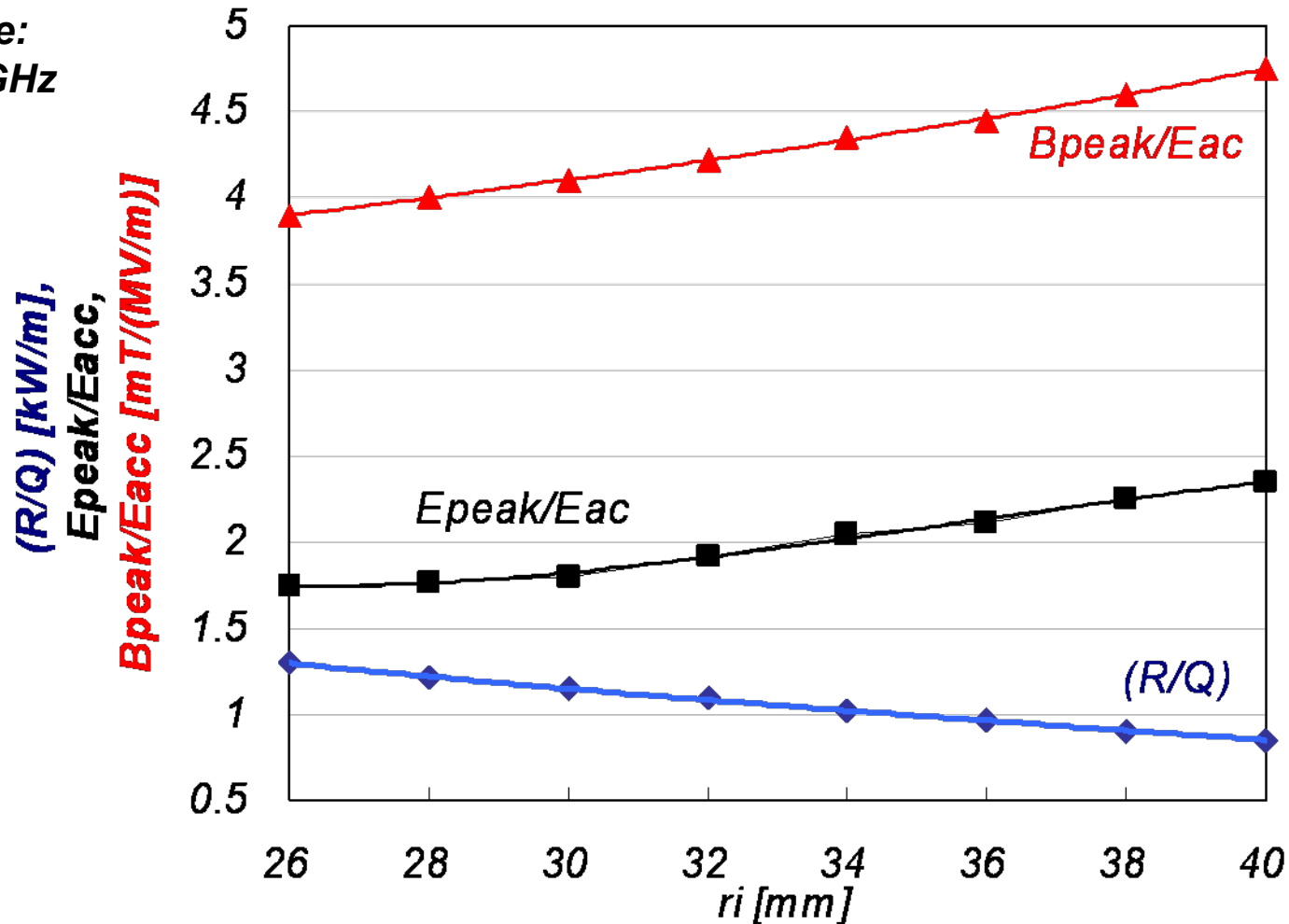
equator ellipsis : half-axis h_r, h_z



There is some kind of conflict 7 parameters and only 5 variables to “tune”

Effect of Cavity Aperture on RF Parameters

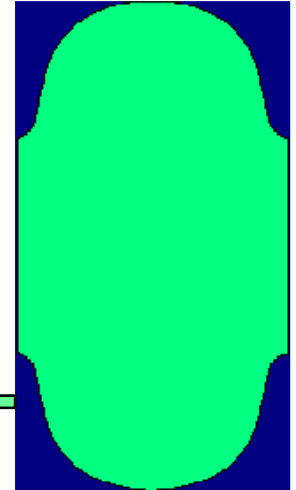
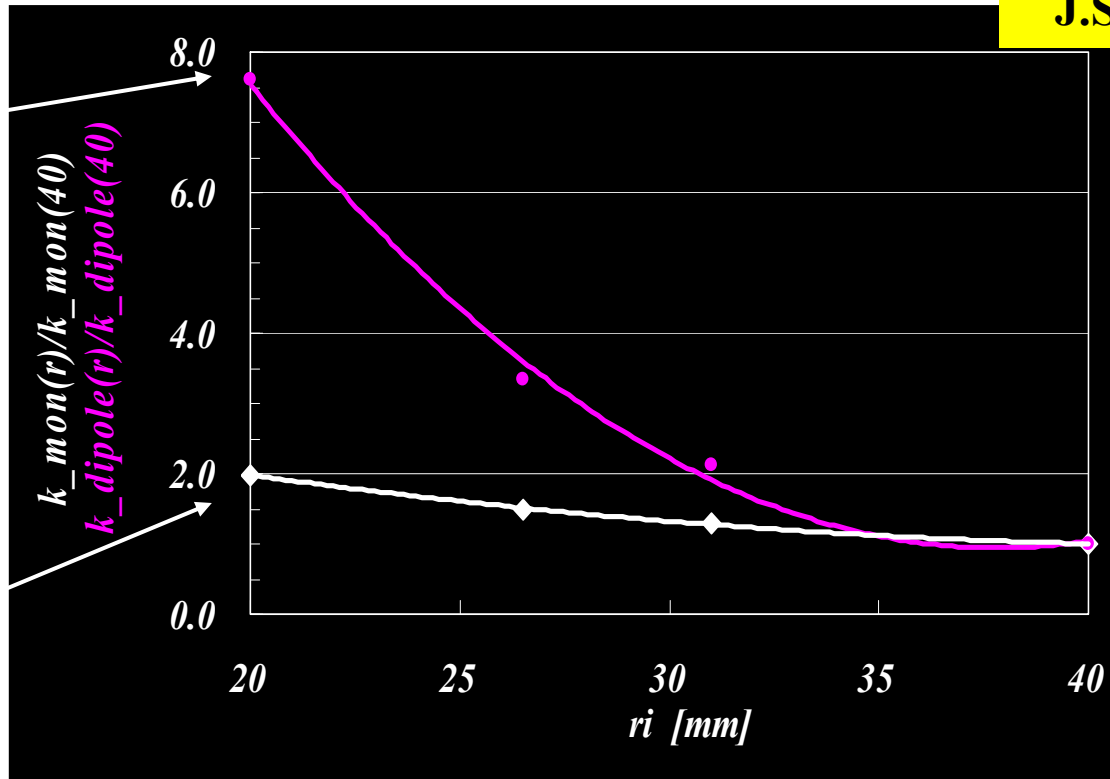
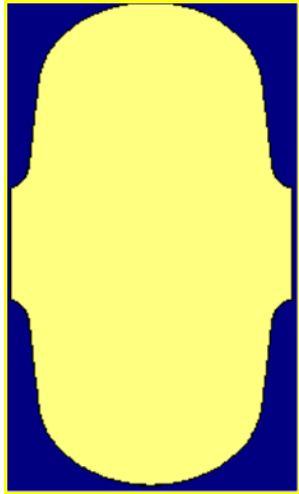
Example:
 $f = 1.5 \text{ GHz}$



A. Mosnier, E. Haebel, SRF Workshop 1991

Aperture Effects on $\kappa_{//}$ and κ_{\perp}

J.Sekutwitz's Slide



$$(R/Q) = 152 \Omega$$

$$B_{peak} / E_{acc} = 3.5 \text{ mT}/(\text{MV}/\text{m})$$

$$E_{peak} / E_{acc} = 1.9$$

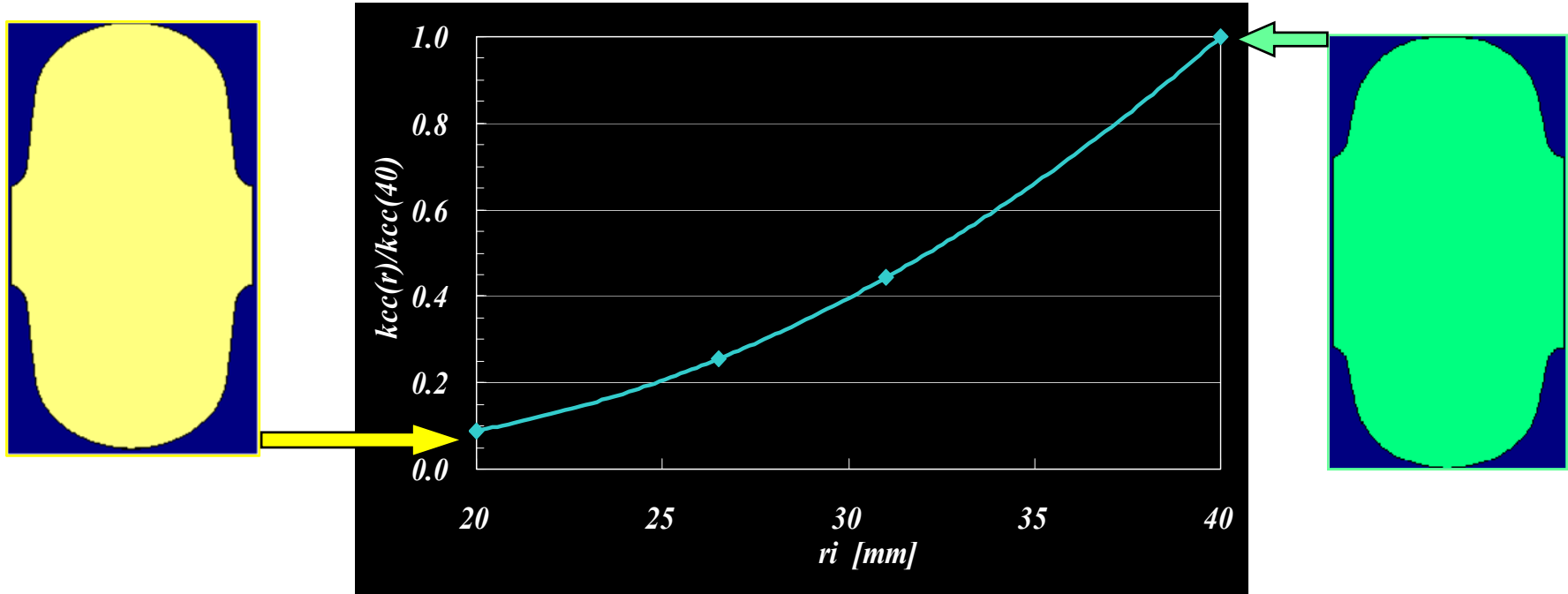
$$(R/Q) = 86 \Omega$$

$$B_{peak} / E_{acc} = 4.6 \text{ mT}/(\text{MV}/\text{m})$$

$$E_{peak} / E_{acc} = 3.2$$

Aperture Effect on Cell to Cell Coupling (κ_{CC})

J.Sekutwitz's Slide



$$(R/Q) = 152 \Omega$$

$$B_{peak} / E_{acc} = 3.5 \text{ mT}/(\text{MV}/\text{m})$$

$$E_{peak} / E_{acc} = 1.9$$

$$(R/Q) = 86 \Omega$$

$$B_{peak} / E_{acc} = 4.6 \text{ mT}/(\text{MV}/\text{m})$$

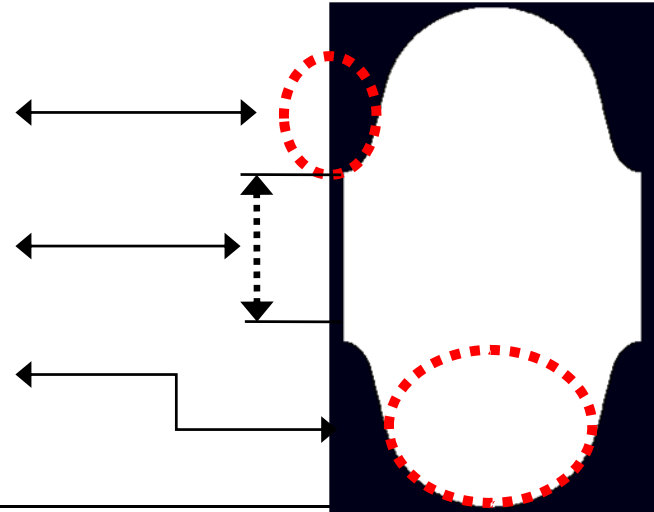
$$E_{peak} / E_{acc} = 3.2$$

General Trends of Cavity Optimization on RF Geometrical Parameters

iris ellipsis : half-axis h_r, h_z

iris radius : r_i

equator ellipsis : half-axis h_r, h_z



<i>Criteria</i>	<i>RF-parameter</i>	<i>Improves when</i>	<i>Cavity examples</i>
<i>Operation at high gradient</i>	E_{peak}/E_{acc} B_{peak}/E_{acc} ↓	r_i ↓ <i>Iris, Equator shape</i>	<i>TESLA,</i> <i>HG CEBAF-12 GeV</i>
<i>Low cryogenic losses</i>	$(R/Q) \cdot \Gamma$ ↑	r_i ↓ <i>Equator shape</i>	<i>LL CEBAF-12 GeV</i> <i>LL- ILC cavity</i>
<i>High $I_{beam} \leftrightarrow$ Low HOM impedance</i>	k_{\perp}, k_{\parallel} ↓	r_i ↑	<i>B-Factory</i> <i>RHIC cooling</i>

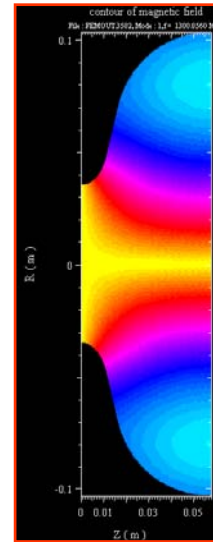
Choice of the RF Frequency

What about accelerating mode frequency of a superconducting cavity?

J.Sekutwitz's Slide



$\times 2 =$



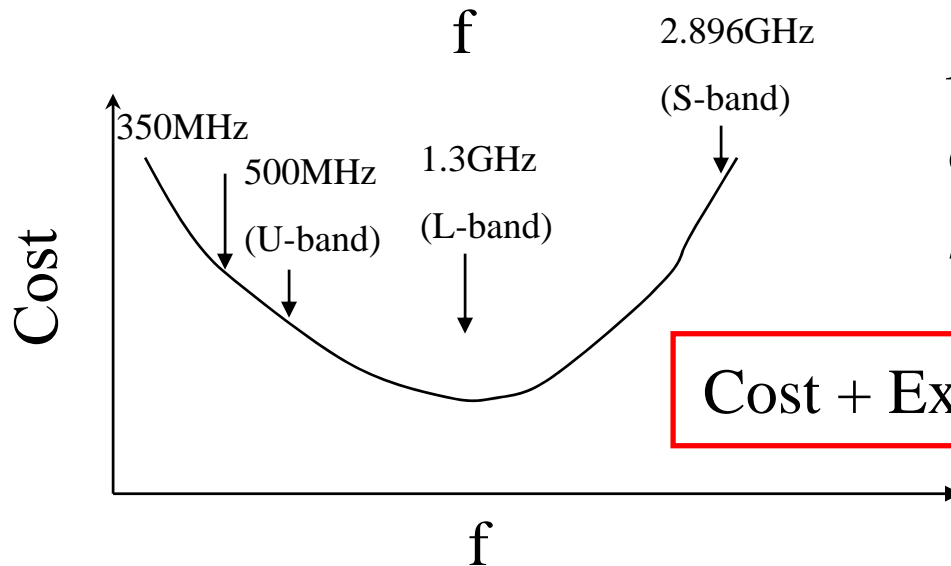
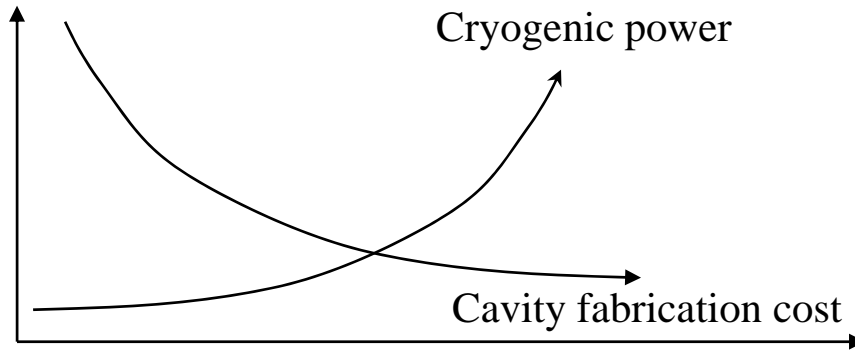
f_{π}	[MHz]	2600
R/Q	[Ω]	57
$r/q=(R/Q)/l$	[Ω/m]	2000
G	[Ω]	271

f_{π}	[MHz]	1300
R/Q	[Ω]	57
$r/q=(R/Q)/l$	[Ω/m]	1000
G	[Ω]	271

$$r/q=(R/Q)/l \sim f$$

Frequency Choice of the SRF Cavity

$$P_{loss} = \frac{1}{2} R_s \int_S H_S dS \propto E_{acc}^2, \quad R_S (BCS) = \frac{A}{T} \cdot f^2 \cdot \exp\left(-\frac{\Delta}{k_B T}\right)$$



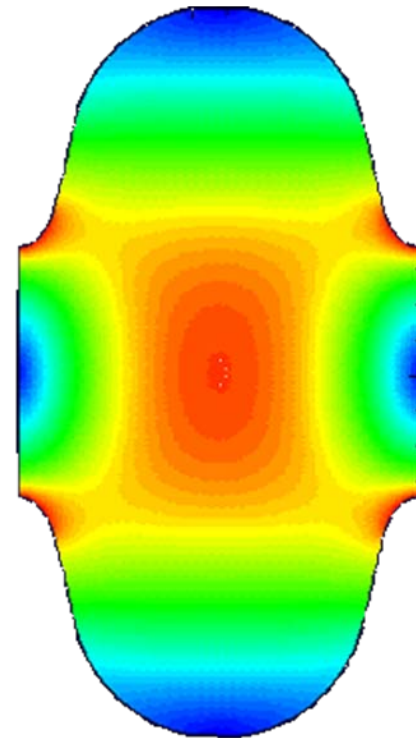
L-band cavity is preferred and it is operated at 2K to reduce the cryogenic power.

TESLA Cavity Inner Cell Shape

The inner cell geometry was optimized with respect to: low E_{peak}/E_{acc} and coupling k_{cc} .

At that time (1992) the field emission phenomenon and field flatness were of concern, no one was thinking about reaching the magnetic limit.

f_{π}	[MHz]	1300.0
r_{iris}	[mm]	35
k_{cc}	[%]	1.9
E_{peak}/E_{acc}	-	1.98
B_{peak}/E_{acc}	[mT/(MV/m)]	4.15
R/Q	[Ω]	113.8
G	[Ω]	271
$R/Q * G$	[$\Omega * \Omega$]	30840



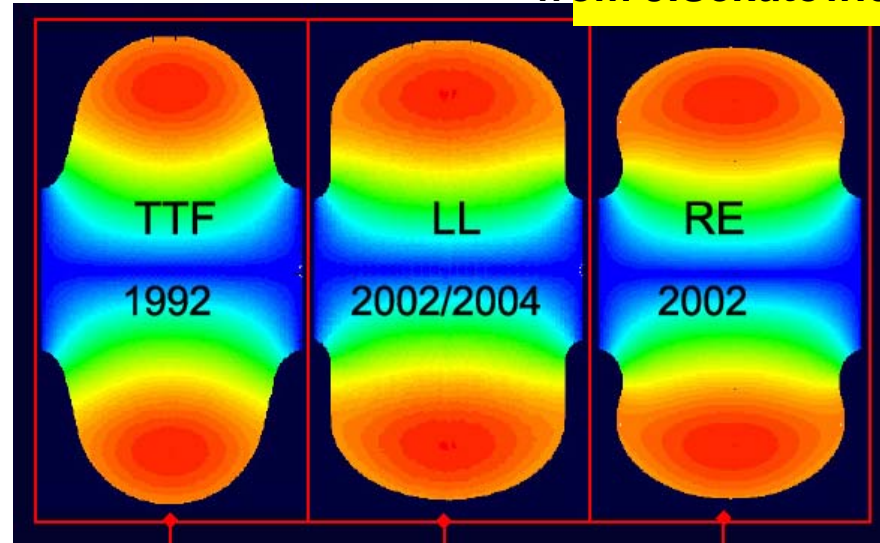
Inner cell; Contour of E field

High Gradient Shapes

Cavity shape designs with low H_p/E_{acc}

from J.Sekutowicz lecture Notes

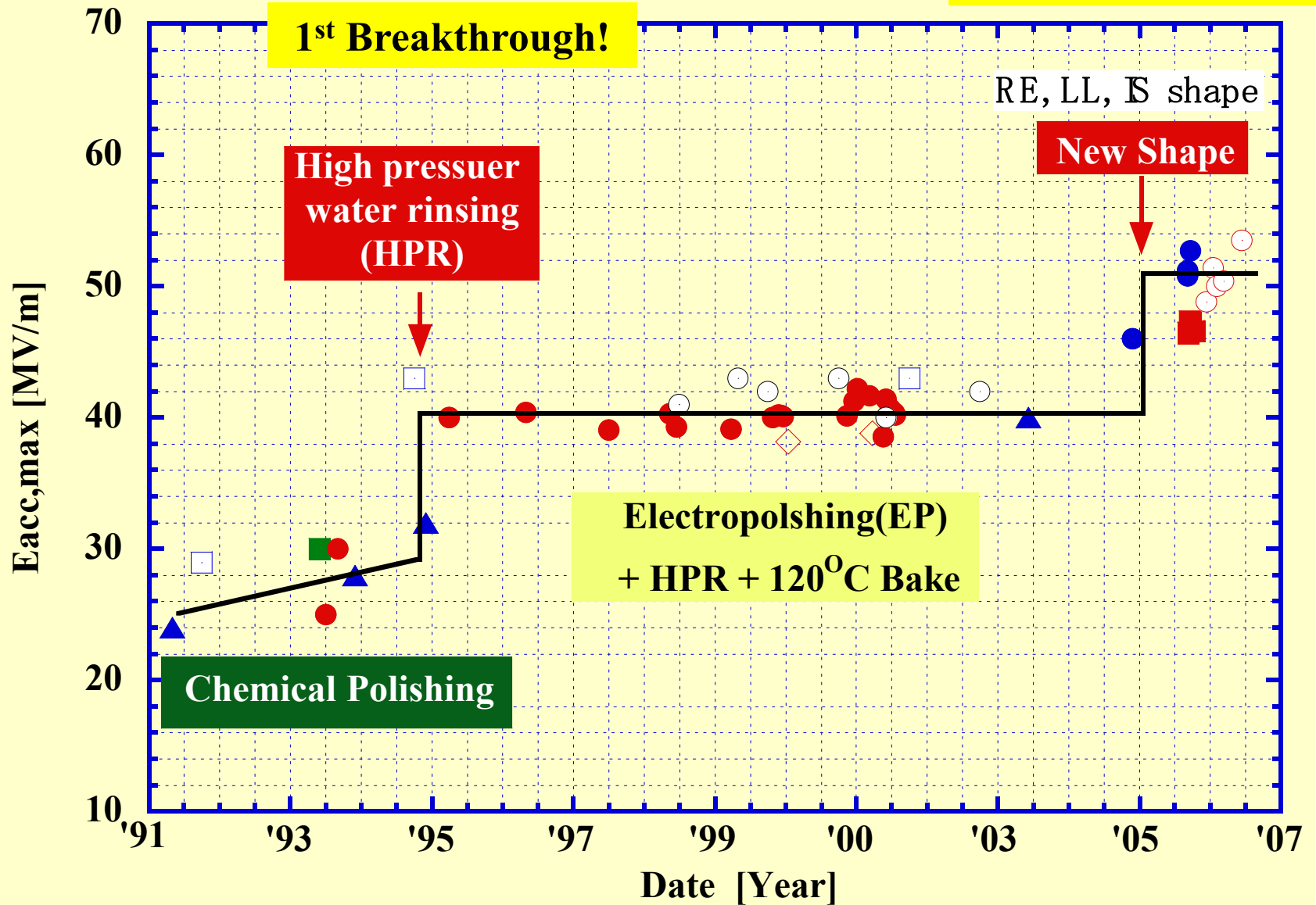
TTF: TESLA shape
Reentrant (RE): Cornell Univ.
Low Loss(LL): JLAB/DESY
Ichiro-Single (IS) : KEK



	TESLA	LL	RE	IS
Diameter [mm]	70	60	66	61
E_p/E_{acc}	2.0	2.36	2.21	2.02
H_p/E_{acc} [Oe/MV/m]	42.6	36.1	37.6	35.6
R/Q [W]	113.8	133.7	126.8	138
G[W]	271	284	277	285
E_{acc} max	41.1	48.5	46.5	49.2

Eacc vs. Year

2nd Breakthrough!



4. Criteria for Multi-cell Structures

Single-cell is attractive from the RF-point of view:

- *Easier to manage HOM damping*
- *No field flatness problem.*
- *Input coupler transfers less power*
- *Easy for cleaning and preparation*
- *But it is expensive to base even a small linear accelerator on the single cell. We do it only for very high beam current machines.*

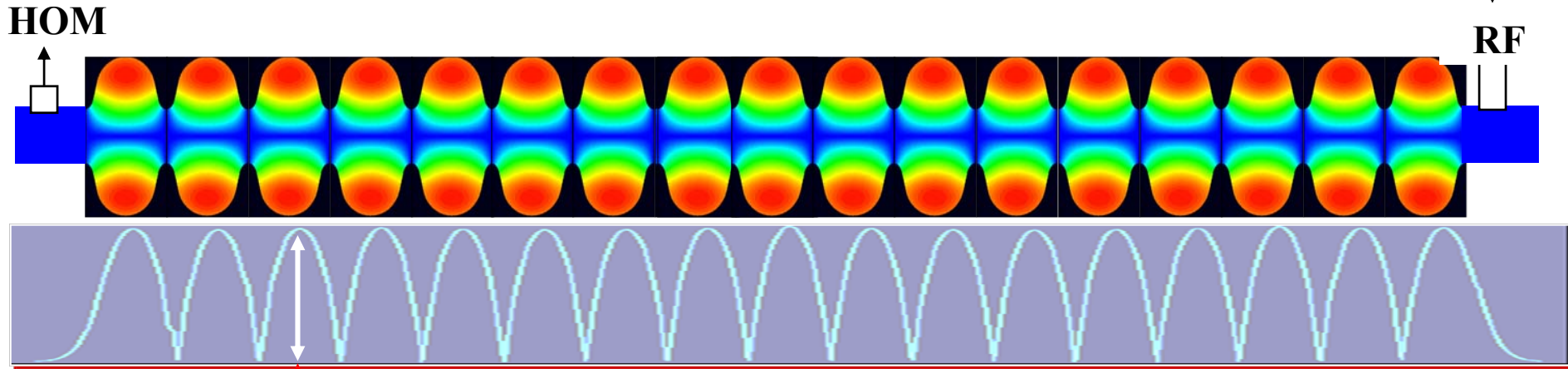
A multi-cell structure is less expensive and offers higher real-estate gradient but:

- *Field flatness (stored energy) in cells becomes sensitive to frequency errors of individual cells*
- *Other problems arise: HOM trapping...*

How to decide the number of cells ?



Parameters considered with N



$$\frac{\Delta A_i}{A_i} = a_{ff} \frac{\Delta f_i}{f_i} = \frac{N^2}{k_{cc}} \cdot \frac{\Delta f_i}{f_i}$$

Beam pipe has no acceleration beam.
BP reduce the efficiency.
Multi-cell is more efficient.

Field flatness factor : $a_{ff} = \frac{N^2}{k_{cc}}$

HOM trapped mode : More serious within creasing N

Input handling power : $P_{INPUT} \propto N$

Handling : Easy to bend, Difficult to preparation

Effect of N on Field Flatness Sensitivity Factor

J.Sekutwitz's Slide

- Field flatness vs. N

	<i>Original Cornell N = 5</i>	<i>High Gradient N = 7</i>	<i>Low Loss N = 7</i>	<i>TESLA N = 9</i>	<i>SNS $\beta=0.61$ N = 6</i>	<i>SNS $\beta=0.81$ N = 6</i>	<i>RIA $\beta=0.47$ N = 6</i>	<i>RHIC N = 5</i>
<i>year</i>	<i>1982</i>	<i>2001</i>	<i>2002</i>	<i>1992</i>	<i>2000</i>	<i>2000</i>	<i>2003</i>	<i>2003</i>
<i>a_{ff}</i>	<i>1489</i>	<i>2592</i>	<i>3288</i>	<i>4091</i>	<i>3883</i>	<i>2924</i>	<i>5040</i>	<i>850</i>



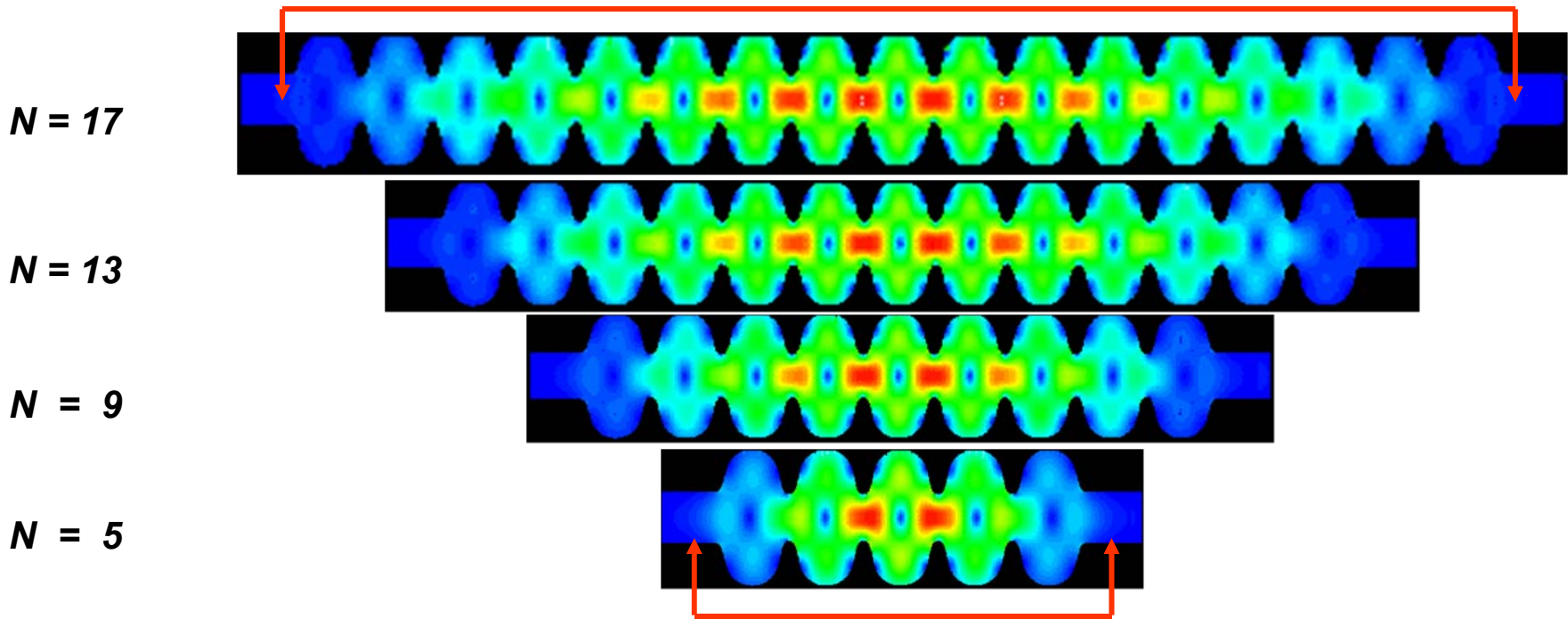
Many years of experience with: heat treatment, chemical treatment, handling and assembly allows one to preserve tuning of cavities, even for those with bigger N and weaker k_{cc}

For the TESLA cavities: field flatness is better than 95 %

HOM trapping vs. N

J.Sekutwitz's Slide

No fields at HOM couplers positions, which are always placed at end beam tubes



e-m fields at HOM couplers positions

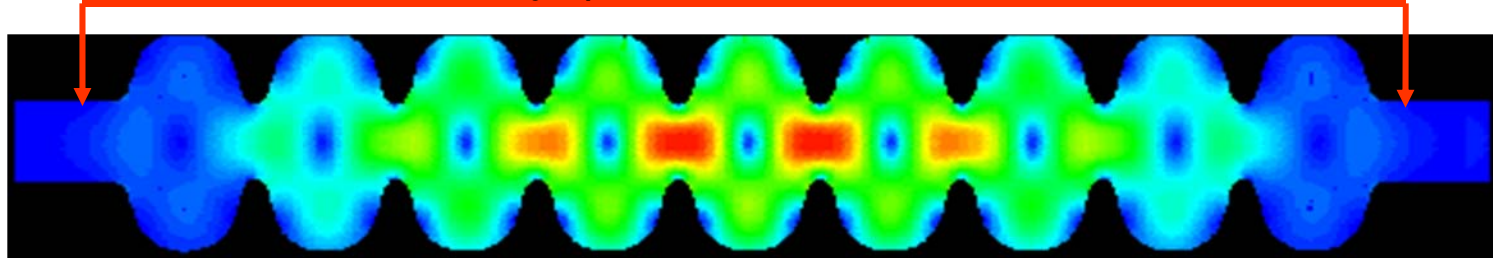
Cure for the trapped mode: Make the bore radius larger.
Break mirror symmetry.

HOM Issue with Multi-Cell Structure

HOM couplers limit RF-performance of sc cavities when they are placed on cells

no E-H fields at HOM couplers positions, which are always placed at end beam tubes

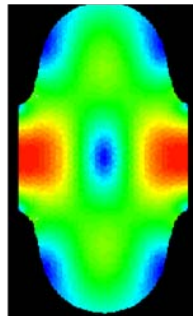
J.Sekutwitz's Slide



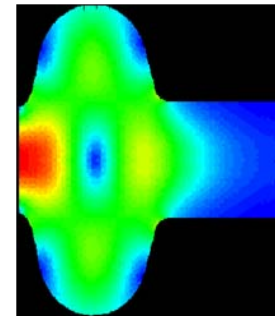
The HOM trapping mechanism is similar to the FM field profile unflatness mechanism:

- *weak coupling HOM cell-to-cell, $k_{cc,HOM}$*
- *difference in HOM frequency of end-cell and inner-cell*

$f = 2385 \text{ MHz}$



That is why they hardly resonate together

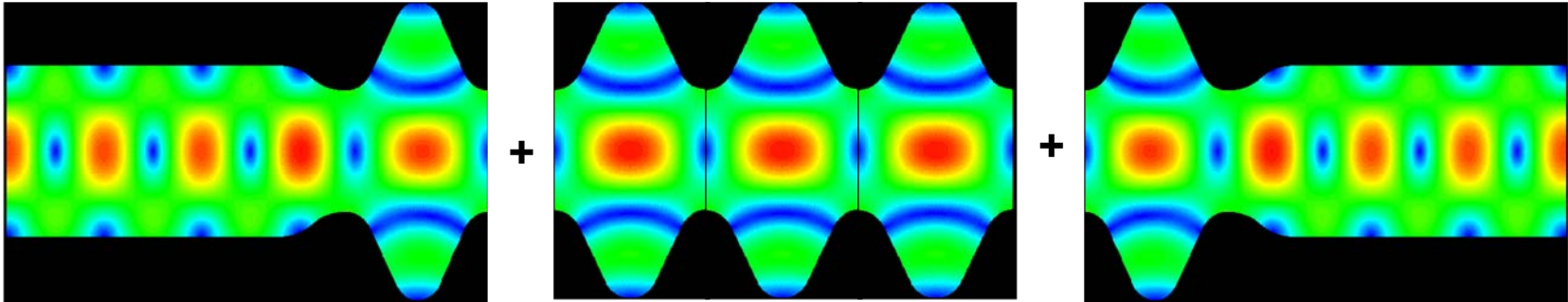


$f = 2415 \text{ MHz}$

Adjustment of End-Cells

J.Sekutwitz's Slide

The geometry of end-cells differs from the geometry of inner cells due to the attached beam tubes



Their function is multi-folded and their geometry must fulfill three requirements:

- ➔ *field flatness and frequency of the accelerating mode*
- ➔ *field strength of the accelerating mode at FPC location enabling operation with matched Q_{ext}*
- ➔ *fields strength of dangerous HOMs ensuring their required damping by means of HOM couplers or/and beam line absorbers.*

All three make design of the end-cells more difficult than inner cells.

Capable Input Power on N

- Power capability of fundamental power couplers vs. N

When I_{beam} and E_{acc} are specified and a superconducting multi-cell structure does not operate in the energy recovery mode:

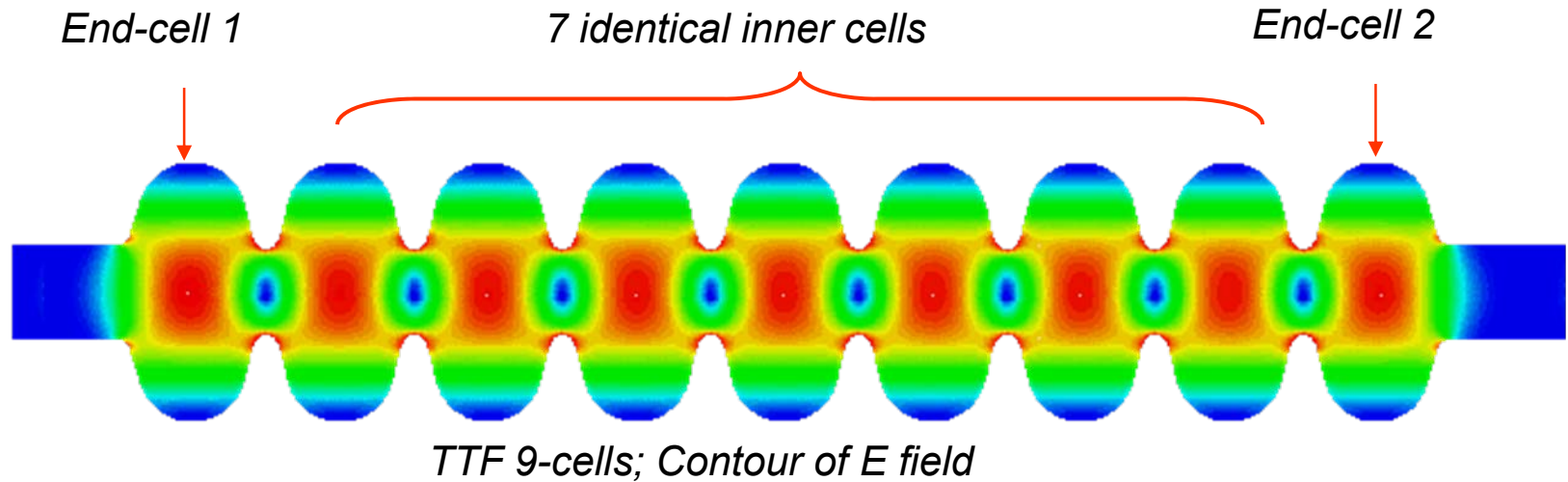
$$P_{INPUT} \sim N$$

**Coupler handling power limits the small N
in the intensity frontier machine like KEK B(N=1).**

KEKB: Acceleration beam current > 1 A.

TESLA/ILC Baseline Cavity

The cavity was designed in 1992 (A. Mosnier, D. Proch and J.S.).



f_{π}	[MHz]	1300.00
$f_{\pi-1}$	[MHz]	1299.24
R/Q	[Ω]	1012
G	[Ω]	271
Active length	[mm]	1038

Overview of Cavities

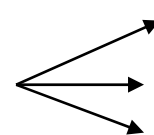
J.Sekutwitz's Slide

Examples of Inner cells

			new	new				new	new
		<i>CEBAF Original Cornell</i> $\beta=1$	<i>CEBAF -12 High Gradient</i> $\beta=1$	<i>CEBAF -12 Low Loss</i> $\beta=1$	<i>TESLA</i> $\beta=1$	<i>SNS</i> $\beta=0.61$	<i>SNS</i> $\beta=0.81$	<i>RIA</i> $\beta=0.47$	<i>RHIC Cooler</i> $\beta=1$
f_o	[MHz]	1448.3	1468.9	1475.1	1278.0	792.8	792.8	793.0	683.0
f_π	[MHz]	1497.0	1497.0	1497.0	1300.0	805.0	805.0	805.0	703.7
k_{cc}	[%]	3.29	1.89	1.49	1.9	1.52	1.52	1.52	2.94
E_{peak}/E_{acc}	-	2.56	1.96	2.17	1.98	2.66	2.14	3.28	1.98
B_{peak}/E_{acc}	[mT/(MV/m)]	4.56	4.15	3.74	4.15	5.44	4.58	6.51	5.78
R/Q	[Ω]	96.5	112	128.8	113.8	49.2	83.8	28.5	80.2
G	[Ω]	273.8	266	280	271	176	226	136	225
$R/Q * G$	[$\Omega * \Omega$]	26421	29792	36064	30840	8659	18939	3876	18045
$k_\perp (\sigma_z=1mm)$	[V/pC/cm ²]	0.22	0.32	0.53	0.23	0.13	0.11	0.15	0.02
$k_\parallel (\sigma_z=1mm)$	[V/pC]	1.36	1.53	1.71	1.46	1.25	1.27	1.19	0.85

β vs RF parameters

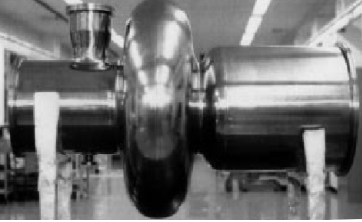


β smaller






E_p/E_{acc} larger
 H_p/E_{acc} larger
 (R/Q) smaller

5. Example of SRF High-beta Cavities

Cavities operating with highest I_{beam} or E_{acc}

<i>Type /No. of cavities</i>			<i>$P_{beam}/cavity$ [kW]</i>	<i>$P_{HOM}/cavity$ [kW]</i>
<p>KEK-B 0.5 GHz</p>  <p style="text-align: right;"><i>x 8</i></p>	<p><i>Single-cell with max I_{beam}</i></p>	<p><i>$I_{beam} = 1.34 A$ 1389 bunches cw</i></p>	<p style="text-align: center;"><i>350</i></p>	<p style="text-align: center;"><i>16</i></p>
<p>HERA 0.5 GHz</p>  <p style="text-align: right;"><i>x 16</i></p>	<p><i>Multi-cell with max I_{beam}</i></p>	<p><i>$I_{beam} \leq 40 mA$ 180 bunches cw</i></p>	<p style="text-align: center;"><i>60</i></p>	<p style="text-align: center;"><i>0.13</i></p>
<p>TTF-I, 1.3 GHz</p>  <p style="text-align: right;"><i>x 1</i></p>	<p><i>Multi-cell with max E_{acc}</i></p>	<p><i>$E_{acc} = 35 MV/m$ 1.3ms/pulse 1Hz PRF</i></p>	<p style="text-align: center;"><i>~100</i> <i>Almost no beam loading</i></p>	<p style="text-align: center;"><i>0</i></p>

Cavities which will operate with high I_{beam} in the near future

<i>Type /No. cavities</i>			$P_{beam}/cavity$ [kW]	$P_{HOM}/cavity$ [kW]
<p>SNS $\beta= 0.61, 0.805$ GHz</p>  <p>x 33</p>	<p><i>Multi-cell with max I_{beam}</i></p>	<p>$I_{beam}=38 (59) mA$ $1.3ms/pulse$ $DF = 6 \%$</p>	<p>240 (366)</p>	<p>0.06 peak</p>
<p>SNS $\beta= 0.805$ GHz</p>  <p>x 48</p>			<p>482</p>	<p>0.06 peak</p>
<p>TTF-II ep , 1.3 GHz</p>  <p>x 8</p>	<p><i>Multi-cell with max E_{acc}</i></p>	<p>$E_{acc}= 35 MV/m$ $1.3ms/pulse 10Hz$ PRF</p>	<p>146</p>	<p>< 0.02 ></p>