

# Site-Selective Dissociation upon Sulfur L-edge X-ray Absorption in a Gas-Phase Protonated Peptide

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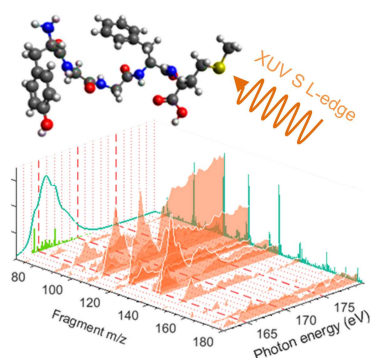
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## ABSTRACT

Site-selective dissociation induced by core photoexcitation of biomolecules is of key importance for the understanding of radiation damage processes and dynamics, and for its promising use as “chemical scissors” in various applications. However, identifying products of site-selective dissociation in large molecules is challenging at the carbon, nitrogen and oxygen edges because of the high recurrence of these atoms and related chemical groups. In this paper we present the observation of site-selective dissociation at the sulfur L-edge in the gas-phase peptide methionine enkephalin which contains only a single sulfur atom. Near-edge X-ray absorption mass spectrometry has revealed that the resonant  $S\ 2p \rightarrow \sigma^*_{C-S}$  excitation of the sulfur contained in the methionine side chain leads to site-selective dissociation, which is not the case after core ionization above the sulfur L-edge. The prospects of such results for the study of charge dynamics in biomolecular systems are discussed.

## TOC GRAPHICS



## KEYWORDS

Peptide • Electronic structure • Site-selective dissociation • X-ray absorption spectroscopy • Mass spectrometry

The interest in studies on site-selective dissociation (SSD) upon core excitation and ionization of small molecules of increasing complexity has emerged several decades ago<sup>1</sup> with the availability of advanced light sources in the XUV and soft X-ray regime. Core-hole states of molecules composed of atoms with low atomic numbers usually decay via Auger processes, yielding a single or multiple valence hole state. In the case of core excitation, SSD may be expected if the resulting valence hole is localized in the vicinity of the initial core hole and if the fragmentation precedes energy redistribution.<sup>2-8</sup> On the contrary, it has been shown that the character of the dicationic state of a molecule after core ionization prevents SSD.<sup>3</sup> Site-specific absorption as such has brought deep insight into charge dynamics with the help of X-ray free-electron laser (FEL) facilities.<sup>9,10</sup>

Studies on larger, biologically relevant systems like peptides, proteins and oligonucleotides in the frame of radiation damage and the utilization of SSD as “chemical scissors” are of great importance. The former can help understanding the early ultrafast nanoscale physical processes in cancer treatment based on ionizing radiation, while the latter finds application in protein and DNA sequencing allowing for targeted bond cleavage near specific chemical groups or atoms. Earlier experiments by Weinkauff and Schlag,<sup>11-13</sup> followed by Lifshitz *et al.*<sup>14</sup> have been performed using nanosecond UV lasers to resonantly photoionize neutral peptides locally on an absorbing aromatic side chain (Trp, Tyr). On peptides, it has been shown that dissociation on a side opposite to the absorbing chromophore may happen and that two dissociation processes are competing: a fast (fs) hole-driven bond-selective dissociation (although not localized near the absorption site) and a slower statistical fragmentation which follows the intramolecular vibrational redistribution of the energy (IVR, ps timescale) by up to milliseconds.

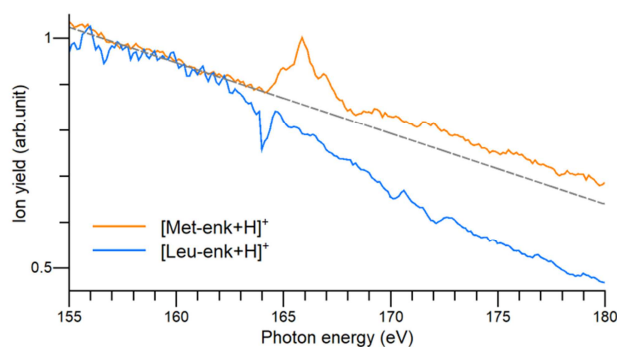
Gas-phase investigations on photoionization of large mass-selected biomolecules have been carried out only recently by interfacing tandem mass-spectrometers, based on electrospray ionization (ESI) sources and radio-frequency ion traps, with synchrotron and FEL<sup>15</sup> beamlines in both the VUV<sup>16-20</sup> and soft X-ray<sup>21-28</sup> energy range. Near edge X-ray absorption mass spectrometry (NEXAMS)<sup>29</sup>, where a mass spectrum of the photofragments is recorded for each photon energy step while scanning across an atomic absorption edge, has proven to be an excellent technique for investigating electronic and geometric structures as well as fragmentation pathways of peptides and proteins. In such systems, carbon, nitrogen and oxygen are the most abundant atoms and are distributed across the whole molecule. Consequently, the bottleneck of investigating SSD upon core excitation or ionization is that the localization of the absorption site cannot be determined. For example, despite the strong resonant excitations to  $\pi^*_{C=C}$  in aromatic side chains and to  $\pi^*_{C=O}$  in the carbonyl groups of the peptide backbone from C, N and O K-shells observed in peptides,<sup>21,25</sup> the information retrieved is averaged over different absorption sites and the fragmentation pathways are similar.<sup>26</sup> Probing the inner shells of a single sulfur atom included in cysteine or methionine residues as the only excitation site is a promising approach to overcome this obstacle and could pave the way for the study of charge dynamics in large biomolecular systems using soft X-rays in pump-probe experiments.

To test the feasibility of this hypothesis we explored for the first time a singly protonated methionine containing pentapeptide [Met<sup>5</sup>]-enkephalin (Tyr-Gly-Gly-Phe-Met, C<sub>27</sub>H<sub>36</sub>N<sub>5</sub>O<sub>7</sub>S, m/z = 574, Met-enk in the following) by means of NEXAMS across the sulfur L-edge at the UE52-PGM nanocluster trap beamline<sup>30,31</sup> at the BESSY II synchrotron (HZB Berlin, Germany) (see experimental section). Met-enk is an opioid peptide, only differing from leucine-enkephalin (Tyr-Gly-Gly-Phe-Leu, C<sub>28</sub>H<sub>37</sub>N<sub>5</sub>O<sub>7</sub>, Leu-enk) by its C-terminal residue. Leu-enk has been

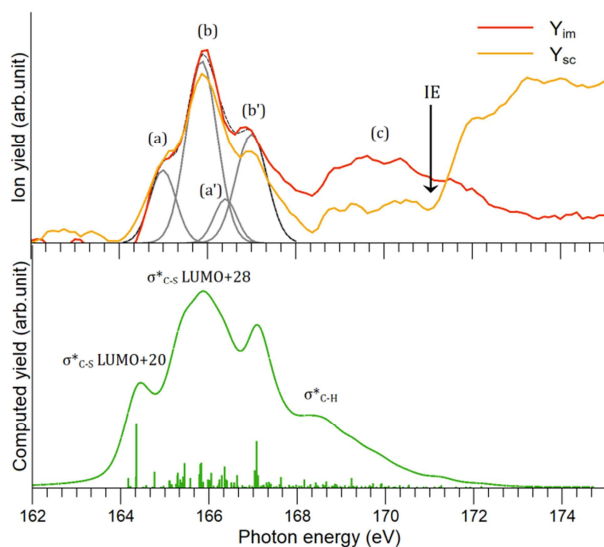
widely studied as a model peptide in many spectrometric studies.<sup>16,19,29,32-37</sup> The structural similarity of Met-enk to Leu-enk makes it the ideal choice for the investigation of SSD through sulfur core excitation and ionization. However, as the total absorption cross section of a molecule increases with its size, the contrast in cross section for the photon absorption by a single sulfur atom is drastically decreased. The contrast in absorption strength at the S 2p orbital against another site in the peptide can be estimated considering the sum of the tabulated atomic cross sections<sup>38</sup> for C, N, O and H,  $\sigma(\text{C}_{27}\text{H}_{36}\text{N}_5\text{O}_7)$ , and the reported absolute absorption cross section for H<sub>2</sub>S at the S 2p inner-shell region,  $\sigma(\text{S})$ . One can expect  $\sigma(\text{S}):\sigma(\text{C}_{27}\text{H}_{36}\text{N}_5\text{O}_7)$  ratios of 9:91 for the absorption strength for the S 2p resonant excitation at 166.7 eV and of 11:89 for the S 2p ionization at 180 eV. In the following, we will first discuss the features observed in the NEXAMS spectrum of Met-enk across the sulfur L-edge and then detail the mass spectra where evidence for SSD has been found. The trapping parameters of the nanocluster trap were optimized to a m/z of 60 to 200, where most of the fragments for this peptide are expected.<sup>16,19,29,39</sup> As all fragments cannot be detected at once, we discuss the absorption spectra in the following in terms of ion yield rather than total ion yield.

**Figure 1** shows the ion yield spectra for both enkephalins at the sulfur L-edge, between photon energies of 155 and 180 eV. The spectra were obtained by summing up the yields of the observed fragments in the mass spectra and were normalized by the photon flux (see **Figure 4**). For Leu-enk, which contains no sulfur atom, the ion yield shows no resonance and declines due to the decrease in photoabsorption cross section by a factor 2 over the studied photon energy range. Met-enk though, shows resonant excitation peaks centered at 166 eV. These features are superimposed on a baseline caused by the fragmentation of the peptide following non-resonant photoabsorption by the valence orbitals. At 166 eV, the contribution to the fragmentation yield

resulting from the absorption at the sulfur atom is ~13 % which is in agreement with the anticipated ratios calculated from the absorption cross sections.



**Figure 1.** Ion yield spectra for protonated  $[\text{Met-enk+H}]^+$  and  $[\text{Leu-enk+H}]^+$  around the sulfur L-edge, adjusted at 160 eV. The ion yield is obtained by integrating all fragment peaks in the mass spectra for each photon energy. The dip observed at 164 eV for  $[\text{Leu-enk+H}]^+$  is due to precursor ion intensity fluctuations. The dashed line is a linear extrapolation of the pre-edge trend of  $[\text{Met-enk+H}]^+$  as used for baseline correction.



**Figure 2.** Top: baseline-corrected ion yield spectrum for the tyrosine immonium ( $Y_{im}$ ) and side chain ( $Y_{sc}$ ) fragments. Gaussian fitting leads to four main transitions of width 0.8 eV (FWHM)

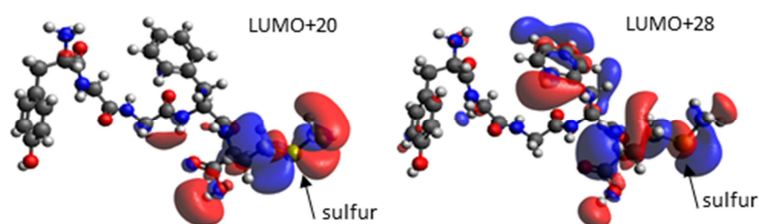
centered at the energies (a) 164.9 eV ( $2p_{3/2} \rightarrow \text{LUMO}+20$ ), (a') 166.4 eV ( $2p_{1/2} \rightarrow \text{LUMO}+20$ ), (b) 165.9 eV ( $2p_{3/2} \rightarrow \text{LUMO}+28$ ) and (b') 167.0 eV ( $2p_{1/2} \rightarrow \text{LUMO}+28$ ). Bottom: DFT/ROCIS calculation of the transitions from sulfur 2p to unoccupied molecular orbitals. The green curve is obtained by a 0.68 eV (FWHM) pseudo-Voigt function (0.3 eV FWHM Lorentzian with 0.5 eV Gaussian functions) broadening of the transition lines. The computed spectrum is shifted by 6.7 eV to fit the experimental data.

Examining the ion yields of individual fragments allows for extracting more detailed information. The top panel of **Figure 2** shows the baseline-corrected ion yield spectra for the tyrosine immonium ( $Y_{\text{im}}$ ) and side chain ( $Y_{\text{sc}}$ ) fragments (see **Figure 4** and discussion). For each fragment ion yield, the baseline subtraction is made using a linear extrapolation of the pre-edge trend of the individual spectrum, as shown in example in the **Figure 1**. Here, data from  $[\text{Leu-enk}+\text{H}]^+$  cannot be used as baseline because the nature of the fragments and their yield may differ from  $[\text{Met-enk}+\text{H}]^+$  due to electronic and geometric structural differences. While only resonant excitations (peaks (a), (b) and (c)) contribute to the ion yield of  $Y_{\text{im}}$ , the yield of  $Y_{\text{sc}}$  shows a strong increase upon core ionization. An ionization energy of  $\sim 170.9$  eV is extracted from these data, consistent with measurements on small sulfur-containing molecules.<sup>40</sup> In order to assess the probed molecular orbitals of the resonant excitations, density functional theory / restricted open-shell single excitation configuration interaction (DFT/ROCIS) calculations, including spin-orbit coupling (see computational section), were conducted for protonated Met-enk. In these calculations, the additional proton was positioned at the N-terminus, as was determined for protonated Leu-enk<sup>35</sup>. Gaussian fitting of the experimental data gives rise to four transitions of full width at half maximum (FWHM) 0.8 eV and centered at 164.9, 165.9, 166.4 and 167 eV. As the lifetime broadening is expected to be 0.3 eV<sup>41</sup> and the photon energy



bandwidth is of 0.077 eV, the ~0.4 eV remaining broadening of the transitions mainly originates from additional vibrational excitations. Therefore the broadening of the transition lines of the calculated absorption spectrum have been made with pseudo-Voigt functions from the linear combination of 0.3 eV width Lorentzian with 0.5 eV width Gaussian functions. **Figure 2** shows a good agreement between the ion yield spectrum (top panel) and the convoluted absorption spectrum (bottom panel). By comparison with the calculations four main resonances have been assigned to transitions from the two spin orbit split of the S 2p levels  $2p_{3/2}$  and  $2p_{1/2}$  to two unoccupied molecular orbitals mainly involving antibonding  $\sigma^*_{C-S}$  states:  $2p_{3/2} \rightarrow \text{LUMO}+20$  (a),  $2p_{1/2} \rightarrow \text{LUMO}+20$  (a'),  $2p_{3/2} \rightarrow \text{LUMO}+28$  (b) and  $2p_{1/2} \rightarrow \text{LUMO}+28$  (b') where the LUMO is the Lowest Unoccupied Molecular Orbital and the +20 or +28 relate to the molecular orbital number relative to the LUMO. The average energy separation between the two  $2p_{3/2}$  and  $2p_{1/2}$  states is ~1.3 eV, close to the expected value of 1.16 eV for the spin-orbit splitting of the S 2p orbital.<sup>42</sup> At higher energies, the broader feature centered at 168-170 eV cannot be assigned to a specific antibonding orbital but mainly contains molecular orbitals of  $\sigma^*_{C-H}$  character in the proximity of the methionine side chain. The LUMO+20 and LUMO+28, depicted in **Figure 3**, are the least energetic unoccupied molecular orbitals involving the sulfur atom. Contributions of  $\sigma^*_{O-H}$  located in the COOH group of the LUMO+20 and  $\sigma^*_{C-H}$  in the phenyl ring of the phenylalanine of the LUMO+28 are weaker than the  $\sigma^*_{C-S}$  contribution by a tenth and a half respectively. Contrary to the observations at the C, N and O K-edges for peptides<sup>24</sup>, here no strong transitions to  $\pi^*_{C=O}$  or  $\pi^*_{C=C}$  states of neighboring chemical groups have been identified in our calculations. It is worth noting that in the case of gas-phase<sup>43</sup> and condensed-phase<sup>44</sup> methionine upon sulfur K-edge excitation, two molecular orbitals of similar character are probed, corresponding in these cases to the LUMO and LUMO+1. Additionally, our ion yield spectrum

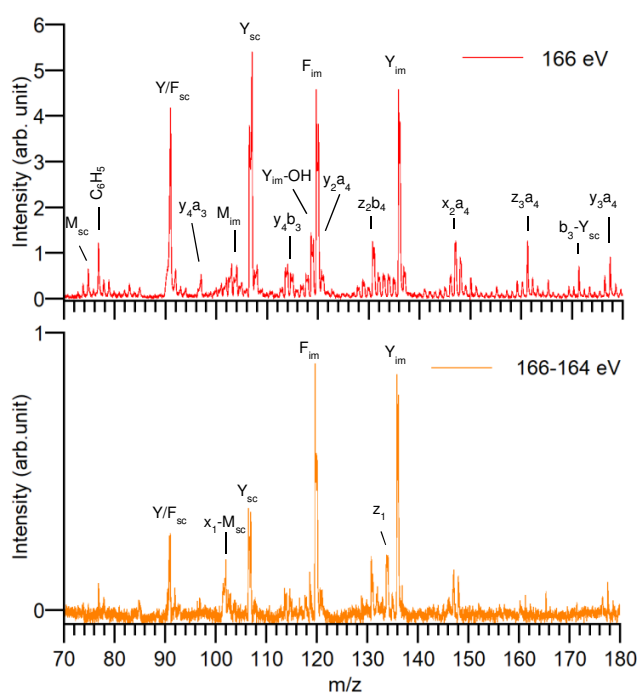
for  $[\text{Met-enk+H}]^+$  is in very good agreement with the X-ray absorption spectrum of methionine powder<sup>45</sup> and the photo-electron yield spectrum of isolated  $\text{H}_2\text{S}$  molecule<sup>46</sup> measured at the sulfur L-edge. This suggests that the proximity of other residues has only little effect on the transitions probed at the sulfur L-edge.



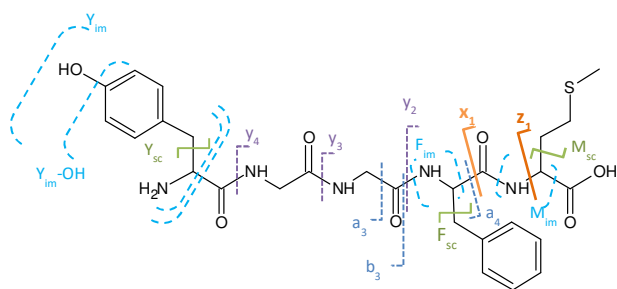
**Figure 3.** Calculated electronic densities of the two main molecular orbitals probed at the sulfur L-edge. The LUMO+20 and LUMO+28 correspond to the first two unoccupied molecular orbitals involving the sulfur atom of the methionine side chain. Details of the atomic orbital contributions can be found in the supporting information.

Despite the small contrast in absorption cross section, evidence for the photoabsorption at a sulfur atom in a large peptide has been observed and the different transitions accessible upon sulfur L-edge photoabsorption in  $[\text{Met-enk+H}]^+$  have been identified. Still, evidence for site-selective dissociation can only be found in the photodissociation mass spectra. The top panel of **Figure 4** shows the mass spectrum obtained at 166 eV photon energy *i.e.* on the  $(2p_{3/2} \rightarrow \text{LUMO}+28)$  (b) resonance. Fragments are labelled according to the peptide fragment nomenclature established by Roepstorff and Fohlman<sup>47</sup> and Biemann<sup>48</sup> (see **Figure 5**). Briefly, the most intense peaks correspond to the immonium ions (*im*) and side chain fragments (*sc*) of the aromatic residues phenylalanine ( $F_{\text{im}}$ ,  $m/z = 120$  and  $F_{\text{sc}}$ ,  $m/z = 91$  and  $77$ ) and tyrosine ( $Y_{\text{im}}$ ,  $m/z = 136$  and  $Y_{\text{sc}}$ ,  $m/z = 107, 91$  and  $77$ ). The immonium ion and protonated side chain

fragments of the methionine ( $M_{im}$ ,  $m/z = 104$  and  $M_{sc}$ ,  $m/z = 75$ ) are also produced, as well as several internal fragments. This spectrum resembles the spectra of protonated  $[Leu-enk+H]^+$  obtained upon carbon K-edge photoabsorption<sup>29</sup> and keV ion-induced dissociation.<sup>39</sup> Indeed, as the relative increase in fragmentation following the photon absorption at the sulfur is only ~13 % most of the fragmentation is the result of non-resonant absorption.



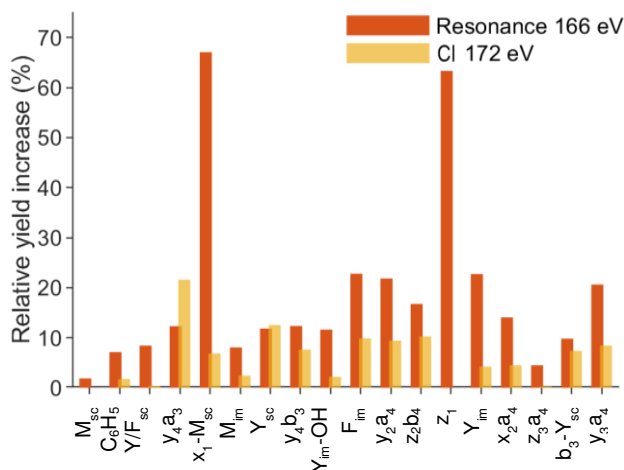
**Figure 4.** Top: photodissociation mass spectrum of  $[Met-enk+H]^+$  at a photon energy of 166 eV. Bottom: difference spectrum after subtraction of the mass spectra at 166 eV (on resonance) – 164 eV (below resonance). For intensity comparison, the same scale is used in both spectra. Letters M, Y and F stand for methionine, tyrosine and phenylalanine respectively. “sc” is for side chain fragment and “im” is for immonium ion.



**Figure 5.** Met-enk chemical structure and cleavage sites.

Without coincidence measurements with the Auger electron emitted after S 2p electron excitation, one way to differentiate the result of the sulfur absorption from non-resonant absorption is to rule out the contribution of the non-resonant absorption. For this purpose **Figure 4-bottom** shows the difference spectrum calculated by subtracting the below-resonance (164 eV) spectrum from the on-resonance (166 eV) spectrum. The obtained mass spectrum includes only dissociation resulting from resonant absorption at the sulfur L-edge. Interestingly, the fragmentation pattern appears similar to the one at 166 eV, the main fragments still being the tyrosine and phenylalanine immonium ions. To evaluate how the absorption at the sulfur specifically contributes to the production of each fragment, in **Figure 6** is plotted the increase of the fragmentation yield induced by the resonant photon absorption at the sulfur L-edge (the difference spectrum) relatively to the yield of the non-resonant absorption (164 eV spectrum) for each fragment. The same operation has been made for the spectrum obtained after core ionization (CI), at 172 eV. At the resonance, the first striking result is the ~65 % increase in yield for the two fragments  $z_1$  and  $x_1$ - $M_{sc}$  (loss of 74 Da,  $C_3H_6S$ ), both situated at the C-terminal side of the peptide, where the methionine residue, and thus the sulfur atom, is located. In contrast, all the other fragment yields exhibit an averaged relative increase of  $(13 \pm 6)$  %. Moreover, despite the absorption at the sulfur atom, the methionine immonium ion  $M_{im}$  and side chain fragment  $M_{sc}$

show a particularly weak yield increase of 8 % and 2 %, respectively. Whereas the LUMO+28 has some  $\sigma^*_{\text{C-H}}$  character on the phenyl ring of the phenylalanine residue, the yields of phenylalanine-related fragments do not show any significant increase in comparison to other fragments. Similarly, we do not observe fragments related to the cleavage of the  $\text{CH}_2\text{-S}$  nor  $\text{S-CH}_3$  bonds although we are probing the C-S anti-bonding orbitals which would weaken these bounds in the methionine side chain. Interestingly, upon core ionization, the  $z_1$  fragment is absent and the yield of the  $x_1\text{-M}_{\text{sc}}$  fragment decreases by a factor of ten. These two fragments seem to be exclusive to the resonant S 2p excitation followed by Auger decay and the processes by which they are produced are hindered in the double-hole state after core ionization and Auger decay. At 172 eV, the relative yield increase drops to  $(6 \pm 5)$  % in average.



**Figure 6.** Increase of the fragmentation yield for each fragment due to absorption at the sulfur at the resonance (166 eV, orange bars) and above the sulfur L-edge, *i.e.* after core ionization (CI) (172 eV, yellow bars) relative to the fragmentation yield of the pre-edge non-resonant absorption.

The high yield increase for the  $z_1$  and  $x_1\text{-M}_{\text{sc}}$  fragments in comparison to all other fragments proves local site-selective dissociations. Peptide fragments of  $z$ - and  $x$ -type are known to be

formed *via* radical (valence hole)-induced processes in radical cations,<sup>49–51</sup> in contrast to the regular b/y-type fragments observed in collision-induced dissociation (CID).<sup>52</sup> Similarly, side chain losses following photo ionization have been identified recently as being caused by the valence hole remaining on the peptide after ionization.<sup>26</sup> In particular the neutral loss of C<sub>3</sub>H<sub>6</sub>S from the methionine side chain has already been observed for methionine-containing peptides upon VUV photoionization<sup>53</sup> and radical-induced dissociation.<sup>54–56</sup> This side chain loss involves the cleavage of the C<sub>α</sub>-C<sub>β</sub> bond and leaves the radical on the backbone of the peptide which can thus undergo further radical-induced dissociations in a cascade fashion. In the present case, this would suggest that the valence hole resulting from the Auger decay initially remains in the vicinity of the sulfur atom and then migrates from the side chain towards the peptide backbone, consequently producing the z<sub>1</sub> and x<sub>1</sub>-M<sub>sc</sub> fragments, the latter involving cascade dissociations. The other fragments show, however, a different behavior. Their similar relative yield increase and the fact that fragments remote from the sulfur, such as Y<sub>im</sub> on the opposite side of the peptide, are equally produced, indicate that the peptide may also undergo a statistical fragmentation following a fast IVR.

In conclusion, we have shown that resonant sulfur 2p photon absorption on a pentapeptide in the gas phase can be observed with a contrast of about ten percent to non-resonant absorption and also leads to site-selective dissociation pathways. Near edge X-ray absorption mass spectrometry proved to be a suited method for suchlike investigations by providing information about both the electronic structure in the proximity of the probed sulfur atom and the subsequent dissociation pathways. We have evidenced that two processes are competing following the excitation of a sulfur 2p electron: SSD induced by the valence hole in the vicinity of the sulfur atom and statistical fragmentation following IVR. However, SSD is not observed upon core

ionization. In this case, the presence of multiple positive charges might induce a strong coulombic repulsion leading to different fragmentation pathways. The following open questions have now to be addressed: To which extent in terms of peptide size can site-specific dissociation still be observed? How do neighboring residues affect the local radical migration and does the radical preferentially migrate towards the N- or C- terminus or another specific chemical group? Which processes govern the fragmentation of such systems after core ionization? Are different timescales involved for SSD and IVR? Testing the region-specificity of these processes in custom-made peptides with the methionine residue placed at different positions will be addressed in future experiments. It is worth noting here that a similar approach could also be envisioned for the study of oligonucleotides by probing the phosphorus atom of the backbone. This topic is now opened to more systematic studies to shed light on these processes. Eventually, this should pave the way for pump-probe studies on large biomolecules in the gas phase at XUV and soft X-ray FEL facilities. Valuable information on the dynamics of charge migration (proton or radical) could thus be obtained, yielding a considerable step forward in the understanding of the physical processes occurring in biomolecules.

## EXPERIMENTAL SECTION

The experiments have been carried out at the Nanocluster trap endstation at the UE52\_PGM beamline<sup>30,31</sup> at the BESSY II synchrotron (HZB Berlin, Germany). The sample ions were produced by electrospray ionization and guided through a radio-frequency hexapole ion guide to a quadrupole mass filter where the precursor ions were mass-to-charge selected. After passing a 90° bender, the ions were buffer-gas cooled, accumulated and stored in a cryogenic linear ion trap at 18 K. The ions were exposed to soft X-ray radiation for 100 ms in the trap and the cation products were mass analyzed in a reflectron time-of-flight mass spectrometer. Due to the

specificity of the setup, the mass resolution is not constant over the whole mass range. The mass resolution  $m/\Delta m$  is of 350 at  $m/z = 77$ , 207 at  $m/z = 120$  and 700 at  $m/z = 177$ . The exit slit of the monochromator was set to an energy bandwidth of 77 meV. Energy scans were performed by steps of 125 meV. Leucine-enkephalin and [Met<sup>5</sup>]-enkephalin acetate salt hydrate (purity  $\geq 95\%$ ), purchased from Sigma-Aldrich, were used without further purification. The peptides were prepared at 30  $\mu\text{M}$  in 1:1 water/methanol solutions with 1% volume formic acid.

## COMPUTATIONAL METHOD

The calculations were carried out with the ORCA program package.<sup>57</sup> Molecular geometry optimizations were performed using the B3LYP<sup>58,59</sup> density functional method with the Ahlrichs TZVP<sup>60</sup> basis set. Transition energies and moments for sulfur L-edge were calculated with DFT/ROCIS using the same basis set. For DFT/ROCIS calculations, the B3LYP functional together with the parameters  $c1 = 0.18$ ,  $c2 = 0.20$ , and  $c3 = 0.40$  was applied.<sup>61</sup> During the calculations, the resolution of identity<sup>62-66</sup> approximation was used employing the Autoaux generation procedure.<sup>67</sup> Numerical integrations during the DFT calculations were performed on a dense grid (ORCA grid4). In all calculations, relativistic effects were taken into account using zero<sup>th</sup>-order regular approximation (ZORA).<sup>68</sup> The geometry calculations had no symmetry constraint. Vibronic effects were not taken into account in the calculations.

## ASSOCIATED CONTENT

Calculated molecular orbitals LUMO+20 and LUMO+28 with corresponding atomic contribution, as well as fragment masses and attributions can be found in the **Supporting Information**.

## AUTHOR INFORMATION



The authors declare no competing financial interests.

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