DOI: 10.1002/ ((please add manuscript number)) Article type: Progress Report

Influence of Radiation on the Properties and the Stability of Hybrid Perovskites

Felix Lang*, Oleksandra Shargaieva, Viktor V. Brus, Heinz C. Neitzert, Jörg Rappich, and Norbert H. Nickel

F. Lang, O. Shargaieva, Dr. V. V. Brus, Dr. J. Rappich, Prof. Dr. N. H. Nickel, Helmholtz-Zentrum Berlin für Materialien und Energie GmbH, Institut für Silizium Photovoltaik, Kekuléstr. 5, 12489 Berlin, Germany. E-mail: felix.lang@helmholtz-berlin.de

Prof. Dr. H. C. Neitzert Department of Industrial Engineering (DIIn), Salerno University, Via Giovanni Paolo II 132, 84084 Fisciano, SA, Italy.

Keywords: perovskite, solar cells, radiation, stability, defects

Organic-inorganic perovskites are well suited for optoelectronic applications. In particular, perovskite single and perovskite tandem solar cells with silicon are close to their market entry. Despite their swift rise in efficiency to more than 21 %, solar cell lifetimes are way below the needed 25 years. In fact, comparison of the time when the device performance has degraded to 80 % of its initial value (T₈₀ lifetime) of numerous solar cells throughout literature reveals a strongly reduced stability under illumination. The various detrimental effects are discussed. Most notably, moisture and heat related degradation can be mitigated easily by now. Recently however, several photo-induced degradation mechanisms have been observed. Under illumination alloyed perovskites tend to phase segregate, while further, oxygen catalyzes deprotonation of the organic cations. Additionally, during illumination photo-generated charge can be trapped in the N-H antibonding orbitals causing the dissociation of the organic cation. On the other hand, organic-inorganic perovskites exhibit a high radiation hardness that is superior to crystalline silicon. This progress report thoroughly reviews proposed degradation mechanisms reported in literature and discusses the microscopic mechanisms and their implications for solar cells.

1. Introduction

Exceptional electronic and optical properties attracted great interest in organic-inorganic perovskites within the last years. Because of their direct band-gap, large absorption coefficient ^[1], low defect concentration ^[2] and large diffusion lengths ^[3,4] this material class is well suited for optoelectronics. Applications range from efficient light emitting diodes ^[5–13] and lasers ^[14–23] to photodetectors ^[24–39] and solar cells. Today's best performing solar cells exceed power conversion efficiencies (η) of 21%. ^[40,41] In fact, organic-inorganic perovskites outperform silicon solar cells already in terms of their potential loss (difference between Voc and band-gap), as they are not limited by Auger recombination.^[40] The tunable band-gap energy makes the material class further highly attractive for tandem solar cells. ^[42] The combination of the high band-gap perovskite and a low band-gap semiconductor such as silicon allows to compensate for thermalization losses. Efficiencies are close to outperform the respective single junctions ^[43–45] and realistic estimates with over 28 % power conversion efficiency are within reach.^[46] However, organic-inorganic perovskites tend to degrade under numerous conditions. ^[47–52] To be more specific, degradation of organic-inorganic perovskites has been observed under illumination ^[53–58], at elevated temperatures ^[59–65], under bias ^[66–68], in moisture ^[63,69–86], in oxygen ^[53,57,76,87,88], and under high-energy irradiation ^[89–92]. The reported solar cell lifetimes are typically less than 2000 hrs. ^[93] Commercially available silicon solar cells, on the other hand, maintain over 85 % of their initial efficiency over a course of 25 years. For a true commercial success, this time span has to be reached for any new material class. In particular, stability is essential for high power perovskite single junctions and tandem solar cells combining perovskites with silicon.

In this review we revisit the degradation of organic-inorganic perovskites. First, we present an overview of the current state of device stability. Considering over 50 publications, we are able to observe two trends: (i) An increase in the reported device stability within the last years and (ii) a severe mismatch of reported lifetimes between storage in the dark or under illumination. Then, we focus on the degradation mechanisms of organic-inorganic perovskites prepared from methylammonium (MA) and formamidinium (FA) lead iodide, which are the two compounds most commonly used. We will show that moisture and heat are two important degradation channels that are both well understood and resolved. Then, the focus will shift on two recently identified degradation channels that are observed during irradiation with light and high-energy particles. The exact details of the degradation mechanisms are still being

2

discussed controversially. Based on recent publications we will show that photo-induced defect formation is the Achilles heel of organic-inorganic perovskite solar cells. As there are no brilliant solutions yet, we will propose routes to go beyond organic-inorganic perovskites.

2. Instability of Perovskite Solar Cells

The very first perovskite solar cells were based on methylammonium (MA) lead iodide (MAPbI₃) quantum dots within a nanocrystalline TiO₂ matrix. The used liquid electrolyte induced a swift dissolution of the perovskite. ^[94] Substitution with a solid-state hole conductor enabled a first improvement in stability from minutes to hours.^[95] With the vast increase of research effort, perovskite solar-cell efficiencies along with their stability increased vastly. Figure 1a depicts the evolution of reported solar-cell lifetimes as a function of their efficiency and their publication year.^[40–42,45,63,72,80,83,94–140] More precisely, the graph depicts the T_{80} lifetime to allow accurate comparison. This lifetime is defined as the time until the efficiency of the device decreases to 80 % of its initial value. ^[141] Blue and orange columns correspond to devices stored under dark and light conditions (100 mWcm⁻²), respectively. Figure 1 reflects the leaping trend of the efficiency vs. year. ^[142] Considering devices stored under dark conditions only, a vast increase of the T₈₀ lifetime over time is obvious (blue columns). Record T₈₀ lifetimes reach 4500 hours. ^[103] This achievement is a result of improved device encapsulation and/or superior barrier layers. The improved devices often include hydrophobic hole transport layers ^[72,83,96,98,102,106,108,111,119,120,125,143,144], inorganic barrier layers ^[103,130,134], PEDOT:PSS additives ensuring a neutral pH value ^[109,123], inorganic charge selective contacts ^[107,116,120,122,128,134,136,137] and perovskite additives ^[72,110] to name a few. Apart from that, compositional changes of the perovskite itself are important. Indeed, various organic and inorganic cations, as well as combinations thereof, have been explored recently. [40,80,145,146]

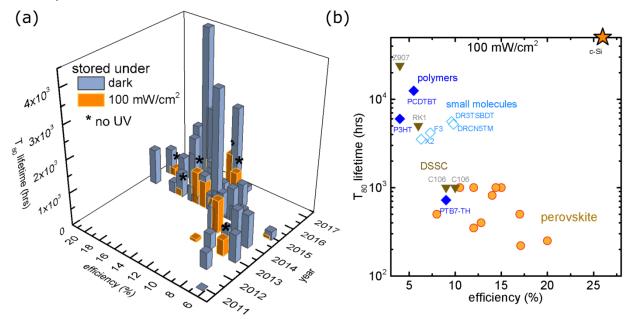


Figure 1. (a) Evolution of the T_{80} lifetime of perovskite solar cells as a function of their initial efficiency and the publication year. Blue and orange columns refer to devices stored in the dark or under illumination with a light intensity of 100 mWcm⁻², respectively. Asterisks refer to studies performed under reduced UV illumination. ^[40–42,45,63,72,80,83,94–140,147] (b) Comparison of the T_{80} lifetime of perovskite solar cells illumination to dye sensitized^[148–151] and organic solar cells based on polymers^[152,153] and small molecules^[154].

Intriguingly, T₈₀ lifetimes are considerably shorter for devices exposed to continuous illumination with a power density of 100 mWcm⁻² (orange columns in figure 1). Under these conditions record lifetimes reach 1000 hours only.^[45] This is far below the set target of 25 years of outdoor operation with over 50,000 sun hours. [Note that this calculation is based on the yearly sum of the global irradiation in San Francisco.] It has to be mentioned that various record lifetimes were collected under illumination with reduced or no UV irradiation. The corresponding data are marked with an asterisk in figure 1. However, the blue to ultraviolet part of the solar spectrum is not negligible and triggers the dissociation of organic cations.^[138,143,155] Moreover, oxygen vacancies are activated within the widely used TiO₂ layers under UV illumination. ^[138,143] The underlying mechanisms are discussed in chapter 3.3. Therefore, for a truthful evaluation of the long-time stability of solar cells it is necessary to use a light source that resembles solar AM1.5G conditions. **Table 1** gives an overview over several outstanding stable perovskite solar cells, which were tested under continuous illumination. Note that some of the reports stopped the long-term measurements before reaching the T₈₀ lifetime.

Table 1.Overview of some perovskite solar cells tested under continuous illumination of 100
mWcm ⁻² . Here FTO is F:SnO ₂ , spiro-OMeTAD is 2,2',7,7'-tetrakis-(N,N-di-p-
methoxyphenyl-amine)-9,9'-spirobifluorene, BCP is bathocuproine, TSHBC a thiolated
nanographene and PTAA poly[bis(4-phenyl)(2,4,6-trimethylphenyl)amine].

year	η	T ₈₀ lifetime	structure	
	[%)	[hrs]		
2017	14.4	>1000 ^[147]	glass/FTO/La:BaSnO3/MAPbI3/NiO/FTO	*UV filter,
2016	12	>1000 ^[45]	glass/ITO/NiO Cs _{0.17} FA _{0.83} PbBr _{0.17} I _{0.83} /LiF/PCBM/SnO ₂ /ITO	*sulfur lamp, MPP
2016	17.7	>500 ^[40]	glass/FTO/TiO2/RbaCsbMAcFAdPblxBry/PTAA/Au	*LED, MPP, 85°C
2016	20	>250 ^[41]	glass/FTO/TiO2/Cs5(MA0.17FA0.83)95Pb(I0.83Br0.17)3/spiro-OMeTAD/Au	*white LED, MPP
2015	15	>1000 ^[156]	glass/FTO/NiMgLiO/MAPbI3/PCBM/Ti(Nb)Ox/Ag	*UV cutoff filter
2015	17.1	220 ^[80]	glass/FTO/TiO2/FA0.9Cs0.1PbI3/spiro-OMeTAD	*sulfur lamp
2015	12	>350 ^[134]	glass/FTO/TiO2/MAPbI3/Al2O3/spiro-OMeTAD	

2015	12.8	400 ^[120]	glass/FTO/TiO2/MAPbI3/TSHBC/Au	# 45% humidity
2015	14	816 ^[98]	glass/ITO/PT/MAPbl ₃ /C ₆₀ /BCP/Ag	
2014	10.5	>1000 ^[121]	glass/FTO/TiO2/(5-AVA)x(MA)1-xPbI3/carbon	
2013	8	500 ^[140]	glass/FTO/TiO2/ MAPbI3/spiro-OMeTAD/Au	*white LED, MPP

To put these achievements into context, we compare T_{80} lifetimes of perovskite solar cells to outstanding dye sensitized (DSSC) ^[148–151] and organic solar cells (OSC) in Figure 1b. For the latter we show several examples based on different polymers^[152,153] and small molecules^[154], which are typically used together with PC₆₁₍₇₁₎BM [6,6-phenyl-C61(71)-butyric acid methyl ester]. So far, dye and organic solar cells exceed the lifetime of perovskite solar cells by roughly one order of magnitude. Research on these technologies however started more than a decade earlier. A very promising long term stability, exceeding 10,000 hrs under 100 mW/cm² illumination, has recently been achieved in a perovskite module following the layer stack glass/FTO/TiO₂/(HOOC(CH₂)₄NH₃)_xMA_yPbI₃/ZrO₂/C. ^[157] It is fair to notice, that the authors kept the module at short circuit and observed an increase in efficiency during the first 500 hrs, which they attributed to concomitant effects and device hysteresis ^[157] Degradation is accelerated if devices are kept at their maximum power point (MPP). ^[56,158] (see also section 3.3.2)

Stabilizing organic-inorganic perovskite solar cells is a key issue for a successful commercialization. However, so far the degradation mechanisms in general and light-induced processes in particular are still not completely unraveled. In the following sections, we provide a comprehensive overview of a number of degradation mechanisms comprising the effects of moisture, heat, oxygen, light, and high-energy particles. Insights to the degradations and the governing microscopic mechanisms will be given.

3. Perovskite Degradation

3.1. Moisture Induced Degradation

Typical organic-inorganic perovskites – such as MAPbI₃ – are strongly hydroscopic. As a result, the material degrades quickly in ambient environment. The process of course can be inhibited by proper device encapsulation and/or by hydrophobic contact or barrier layers.^[72,83,96,98,102,103,106,108,111,119,120,125,130,134,144,159] Nevertheless, the underlying mechanism is of general interest. **Figure 2** depicts the general degradation route. The hydration of CH₃NH₃I and PbI terminated surfaces is exothermic.^[73] Therefore, solvation as well as water incorporation can be quite swift. Interestingly, a water molecule can reside within the

6

perovskite cavity stabilized by hydrogen bonds to methylammonium and iodine. ^[73] Recently, the formed hemihydrate CH₃NH₃PbI₃•¹/₂H₂O has been identified by a subtle shift of the N-H stretching vibrations using Fourier transform infrared spectroscopy (FT-IR). ^[75] Concomitantly, O-H stretching vibrations emerge around 3400 cm⁻¹, as shown in figure 2e. ^[75] The hemihydrate preserves the original crystal structure with minor changes in bond length and a subtle increase of the band gap energy.^[73] Although the hemihydrate formation is reversible, two distinct defect states are created.^[79] Figure 2f depicts temperature stimulated current (TSC) measurements before and after the formation of the hemihydrate by exposure to humid air.^[79] The observed defect states reside in the band gap at energies of 0.32 and 0.42 eV. However, it is unclear whether the energies are given with respect to the valence or conduction band. In addition, hemihydration is accompanied by a strong increase in conductivity.^[75]

Upon further exposure to water vapor the monohydrate (CH₃NH₃PbI₃•H₂O) is formed. As a result, the short circuit current J_{SC} is reduced dramatically.^[69,77] In contrast to the perovskite and the hemihydrate, the color of the monohydrate is pale yellow. Figure 2b depicts a sketch of the crystal structure determined from X-ray diffraction (XRD).^[69,77] The formation of the monohydrate is reversible upon dehydrating. In fact, the short circuit current and the external quantum efficiency (EQE) of TiO₂/MAPbI₃/spiro-OMeTAD [2,2',7,7'-tetrakis(N,N-di-p-methoxyphenylamine)-9,9'-spirobifluorene] solar cells recovers completely, as depicted in figure 2g.^[77] Note that prolonged exposure to H₂O subsequently leads to the formation of the colorless dihydrate (CH₃NH₃)₄PbI₆•2H₂O. The corresponding crystal structure is sketched in figure 2c.^[77] Due to the loss of CH₃NH₃I, the reaction becomes irreversible. In case of further excess of water, the reaction is finally driven to PbI₂ under expulsion of CH₃NH₃I, figure 2d. [74–77,79,82,85]

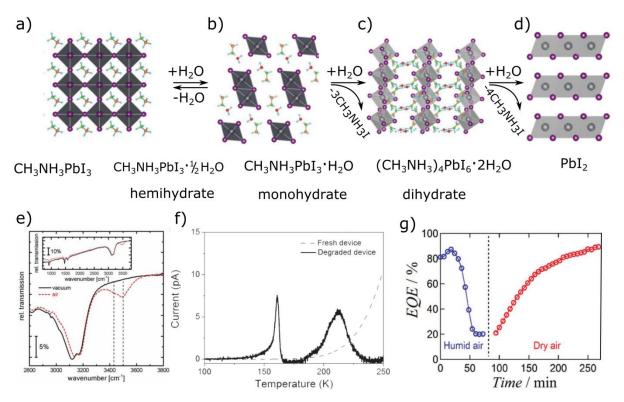


Figure 2. Moisture induced degradation of MAPbI₃. (a), (b), (c), and (d) depict the crystal structure of MAPbI₃, the monohydrate CH₃NH₃PbI₃•H₂O, the dihydrate (CH₃NH₃)₄PbI₆•2H₂O, and lead iodide, respectively. Reproduced with permission. ^[77] Copyright 2016, John Wiley and Sons. (e) FT-IR spectra of MAPbI₃ after exposure to vacuum or air. Reproduced with permission. ^[75] Copyright 2015, American Chemical Society. (f) TSC measurements of freshly prepared and degraded MAPbI₃ based solar cells. Reproduced with permission. ^[79] Copyright 2015, John Wiley and Sons. (g) External quantum efficiency as a function of exposure time to humid or dry air. Reproduced with permission. ^[77] Copyright 2016, John Wiley and Sons

In contrast to the observed degradation, small amounts of water added to the precursor solution can induce a controlled and oriented crystallization of MAPbI₃. As result, charge carrier lifetimes and power conversion efficiencies are enhanced.^[71] Furthermore, impedance spectroscopy and photoluminescence (PL) lifetimes suggest that adsorbed water molecules can effectively suppress certain recombination processes.^[78]

3.2. Thermal Instabilities

In thermal equilibrium the temperature of a solar cell is given by ^[160]

$$T_s = \left[T_0^4 + \frac{1-\eta}{2\cdot\sigma_s} \cdot P_{sun}\right]^{1/4} \tag{1}$$

Neglecting convection, T_S is determined by the Boltzmann constant, σ_S , the efficiency, η , the integrated power of the solar spectrum, P_{sun} , and the temperature of the environment, T_0 . For a solar cell with an efficiency of $\eta = 20$ % and an outdoor temperature of $T_0 = 45^{\circ}C$ (typical

summer conditions), T_s amounts to about 83 °C. Hence, under typical environmental conditions a solar cell needs to withstand this temperature.

In general, perovskites undergo phase transitions depending on the temperature. As an example, MAPbI₃ exhibits two phase transitions. The first one occurs at about 160 K when its phase changes from orthorhombic (Pnma) to tetragonal (I4/mcm). A second phase transition from tetragonal to cubic (Pm3m) is observed at about 330 K.^[161,162] The corresponding crystal structures are depicted in **Figure 3**a.^[162] The band gap energy decreases with increasing temperature for a given crystalline structure and shows an abrupt change at the phase transition (see figure 3b). The high temperature phase transition is of particular interest, since it resides within the operating temperatures of a solar cell. On the other hand, the photovoltaic properties do not change abruptly at the phase transition. ^[163–165] Instead, J_{SC} and V_{OC} decrease monotonically, as shown in figure 3c. The lack of an abrupt change is attributed to the high structural flexibility in the high temperature phase.^[164] Since the temperature of solar cells can vary from night to day between about -40 °C and 85 °C the influence of temperature cycling and phase transitions on the long-term stability are of importance. However, such experiments have not been performed to date.^[166] Note that commercial absorber materials, such as crystalline/amorphous Si, Cu(In,Ga)Se₂, CdTe, GaInAs, GaInP or Ge do not have phase transitions in this temperature range.

Besides structural changes, organic-inorganic perovskites - especially MAPbI₃ - degrade at elevated temperatures. Interestingly, investigations on the temperature stability differ vastly depending on the used technique. On the one hand, thermogravimetric analyses (TGA), as presented in figure 3d, often reveals a thermal stability exceeding T = 250 °C. ^[1,59] On the other hand, numerous other experimental techniques show that MAPbI₃ degrades already at lower temperatures. For instance, XRD shows the reflex of the (001) plane of hexagonal PbI₂ ($2\Theta \approx 12.6^{\circ}$) already at temperatures above 140 °C (figure 3e). ^[64] At the same temperature, the reflex originating from the (110) plane of tetragonal MAPbI₃ ($2\Theta \approx 14.1^{\circ}$) vanishes and at T > 190°C the decomposition of MAPbI₃ to PbI₂ is completed.

Wavelength dependent surface photovoltage (SPV) measurements taken after heating give further insight into the altered charge separation processes (figure 3f).^[65] Up to a temperature of 140 °C the material quality increases since the charge separation by defect states within the band gap is reduced. Above 140°C charge separation by defect states increases due to an additional photovoltage located at about 2.4 eV. This behavior corresponds to charge separation in PbI₂, which is formed due to thermal degradation.

9

Conversely, a temperature of 85°C seems high enough to degrade the surface of MAPbI₃ films. As an example, figure 3g depicts X-ray photoelectron spectroscopy (XPS) measurements of MAPbI₃ films before and after annealing at 85°C in various atmospheres. ^[61] In pristine samples, the C1s core level is dominated by a peak representing the C-N bond of methylamine. Upon prolonged annealing at 85°C in ambient air, pure O₂, or pure N₂ the C-N related peak decreases gradually and the peak related to C-C or C-H bonds increases. The observed surface degradation is sufficient to affect the performance of TiO₂/MAPbI₃/P3HT [poly(3-hexylthiophen-2,5-diyl] devices. In particular, FF and J_{SC} decrease by about 50 % during an anneal at 85°C for 24 hrs in ambient atmosphere. ^[61]

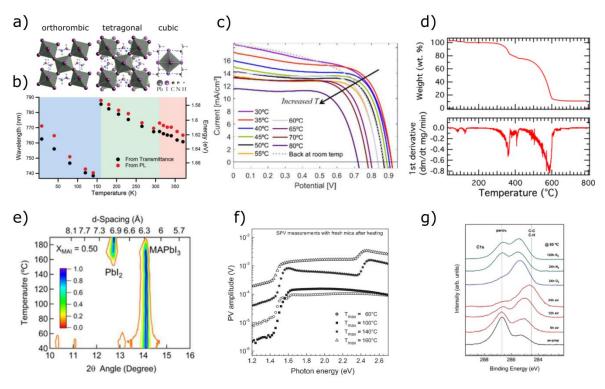


Figure 3. Heat induced degradation of MAPbI₃. (a) Schematic representation of the MAPbI₃ crystal structures. Reproduced with permission. ^[162] Copyright 2015, American Physical Society. (b) Temperature dependent band-gap position of MAPbI₃. Reproduced with permission. ^[167] Copyright 2015, John Wiley and Sons. (c) Temperature dependent *J*-V characteristics of a regular TiO₂/MAPbI₃/spiro-OMeTAD solar cell. Reproduced with permission. ^[165] Copyright 2016, American Chemical Society. (d) TGA heating curves of MAPbI₃. Reproduced with permission. ^[59] Copyright 2014, American Chemical Society. (e) X-ray diffractogram of MAPbI₃ as a function of temperature. Reproduced with permission. ^[64] Copyright 2015, American Chemical Society. (f) SPV measurements of MAPbI₃ layers after annealing at moderate temperatures. Reproduced with permission. ^[65] Copyright 2013, American Institute of Physics. (g) XPS measurements of ITO/TiO₂/MAPbI₃ samples at 85°C in various indicated atmospheres. Reproduced with permission. ^[61] Copyright 2015, John

Compositional engineering allows to strengthen the thermal stability of perovskites. For instance, a replacement of the methylammonium cations with Cs⁺ or formamidinium cations, HC(NH₂)₂⁺, is possible. Cesium lead halides are stable up to temperatures of 300°C.^[168,169] Formamidinium (FA) lead halides are stable up to temperatures of at least 200°C. **Figure 4**a depicts an X-ray diffractogram of FAPbI₃ in the α -phase, acquired at a temperature of 170°C. Notably, reflexes originating from hexagonal PbI₂ are absent. This suggests a higher thermal stability than observed for MAPbI₃. However, when pure FAPbI₃ is cooled to room temperature the perovskite undergoes a phase transition to a hexagonal structure (δ -phase). This phase is yellow and therefore not suitable for solar cells.^[170] This structural transformation can be suppressed by blending formamidinium with methylammonium. ^[170,171] Even more complex blends, comprising rubidium, cesium, formamidinium, and methylammonium, have been explored recently. ^[40,41] These triple or quadruple mixed cation perovskites exhibit high power conversion efficiencies and a superior thermal stability compared to perovskites prepared with pure methylammonium cations. As an example, figure 4b shows the evolution of the power conversion efficiency of a

 $Rb_5(Cs_5MA,FA)_{95}Pb(_{1.83}Br_{0.17})_3$ based solar cell. Although the device is exposed to a temperature of 85°C and continuous illumination with a power density of 100 mWcm⁻² the degradation is less than 5 % over a period of 500 hrs. ^[40] Note that detailed investigations on heating induced defects in triple or quadruple cation perovskites have not been performed to date.

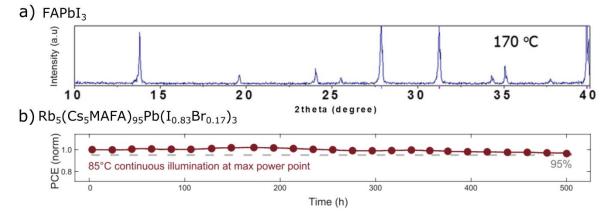


Figure 4. Thermal stability of FA based perovskites. (a) X-ray diffractogram of FAPbI₃ measured at 170°C. Reproduced with permission. ^[171] Copyright 2015, Nature Publishing Group. (b) Power conversion efficiency of a Rb₅(Cs₅MA,FA)₉₅Pb(_{1.83}Br_{0.17})₃ based solar cell as a function of time. The measurements were performed at a temperature of 85°C. The device was continuously illuminated with a light intensity of 100 mWcm⁻² using white LEDs. Reproduced with permission. ^[40] Copyright 2017, The American Association for the Advancement of Science.

3.3. Photo-Induced Degradation

Comparing T₈₀ lifetimes of over 50 publications, we revealed a considerably reduced lifetime of perovskite solar cells under illumination (see figure 1a, chapter 2). Most notably, the sun spectrum contains a considerable amount of high energetic photons, which are particularly harmful to perovskite solar cells.^[138,155] Since solar cells need to be optimized for maximal absorption of a broad spectral range, this is highly problematic. With only a few publications so far, the investigation of photo-induced degradation is still in its infancy. In the following, we critically review recently observed light-induced degradation mechanisms. First, we focus on the interaction of oxygen and light, which results in a swift degradation. ^[53,57,76] However, photo-induced degradation does not require the presence of oxygen. Therefore, we will discuss a number of reversible^[56,158,172] and unfortunately also permanent^[155,158,173] degradation processes reported throughout literature. In particular, we will highlight the observed dissociation of CH₃NH₃⁺, which commences under illumination with $h \nu \ge 2.72$ eV.^[155] Beyond that, mixed cation and mixed halide perovskites are prone to phase segregation under illumination.

3.3.1. Oxygen Catalyzed Photo-Degradation

Mesoporous (mp) TiO₂ sensitized solar cells are susceptible to rapid degradation by ultraviolet light. This effect is a result of oxygen desorption, which activates oxygen vacancies within TiO₂. These act as deep electronic trap states ^[174] and finally quench the performance of perovskite and dye sensitized solar cells. ^[138,143] Replacement of TiO₂ is possible using other inorganic materials such as Al₂O₃^[138], SnO₂ ^[175] or La:BaSnO₃ ^[147] to name only a few.

However, even in the absence of mp-TiO₂ and/or UV-light, the combined action of molecular oxygen and light induces a rapid degradation of organic-inorganic perovskites. ^[53,57,176] **Figure 5**a depicts absorption spectra of MAPbI₃ on glass prior to and after aging in dry air under illumination. ^[57] After a course of 48 hours, the absorption onset clearly shifts from around 800 nm to 500 nm. This corresponds to a complete degradation to PbI₂. It should be noted, that in absence of O₂ or light, no degradation was observed. ^[57] This effect was confirmed by measuring the efficiency of FTO/cp-TiO₂/mp-TiO₂/MAPbI₃/spiro-OMeTAD/Au based devices in various atmospheres, see figure 5b. ^[57] Again, no degradation is observed in absence of O₂ or light. Note that a UV filter was used to prevent the activation of TiO₂ related defects.

The mechanism of the combined oxygen and light induced degradation is currently under debate. ^[53,55,57,155] Super-resolution luminescence micro-spectroscopy, for example, revealed a smooth collapse of the MAPbI₃ crystal structure to PbI₂, starting locally and then spreading over the whole crystal. ^[55] Further, hydroethidine was used recently as a molecular fluorescent probe.^[53] The molecule has a characteristic emission upon exposure to the superoxide (O_2^{-}). Perovskite films that were immersed in the probe solution caused an increase of the characteristic emission when illuminated.^[53] This shows that O_2^{-} triggers the degradation under combined O_2 and light exposure. As the governing mechanism, it was suggested that illumination leads to the formation of the superoxide, which subsequently causes the deprotonation of MAPbI₃ and the formation of methylamine, lead iodide, iodine, and water as byproducts. ^[53]

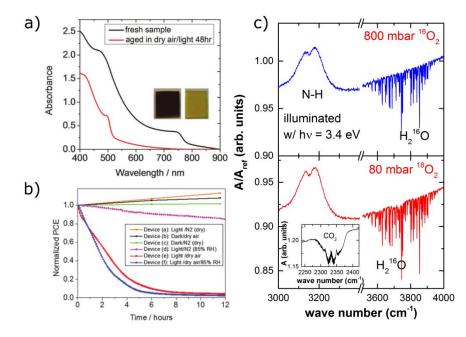


Figure 5. Oxygen catalyzed photo-degradation of perovskites: (a) Absorption spectra of MAPbI₃ on glass prior to and after prolonged illumination. (b) Evolution of the normalized η of FTO/cp-TiO₂/mp-TiO₂/MAPbI₃/spiro-OMeTAD/Au devices in various environments. Illumination was performed with white LEDs that were equipped with a UV filter. (a,b) Reproduced with permission. ^[57] Copyright 2016, The Royal Society of Chemistry. (c), (d) FT-IR spectra of MAPbI₃ that were exposed to UV light and isotopically enriched ¹⁶O₂ or ¹⁸O₂. The inset shows the absorption modes of C¹⁶O₂ and C¹⁸O₂. Reproduced with permission. ^[155] Copyright 2017, unpublished.

The light-induced degradation in the presence of oxygen was recently revisited to elucidate the exact microscopic mechanism. ^[155] Before and after light exposure in the presence of oxygen the samples were characterized using FT-IR measurements. For these experiments isotopically enriched ¹⁶O₂ and ¹⁸O₂ was used. Figure 5c depicts the absorption ratio of the

illuminated to the pristine unexposed sample, A/A_{ref} for FAPbI₃. It is important to note that an increase of the absorption ratio corresponds to a decrease of the corresponding molecular concentration. On the other hand, the absorption lines for water point downward. Since the unexposed sample (A_{ref}) did not show H₂O vibrational modes, their appearance indicates that water molecules escaped from the sample into the airtight sample compartment. The upper panel of figure 5c depicts the infrared absorption after an UV degradation in the presence of $^{16}\text{O}_2$ (800 mbar). The photon energy, photon flux, and degradation time amounted to hv = 3.4eV, photon flux $F \approx 5 \times 10^{17}$ s⁻¹ cm⁻², and t = 4.5 × 10³ s, respectively. The two asymmetric N-H stretching modes are located around 3130 and 3178 cm⁻¹.^[177] Clearly, the NH₃⁺ concentration in the perovskite sample is strongly reduced after the degradation. This is consistent with the dissociation of CH₃NH₃⁺ and the out-diffusion of the degradation products. Simultaneously, absorption lines of water vapor, pointing downwards, emerge. These H₂O molecules evolved from the MAPbI₃ thin films. Previously, it was suggested that H₂O forms from oxygen that diffuses into the specimens and protons from the organic cations. ^[53,57] To test this hypothesis the degradation experiment was repeated in the presence of isotopically enriched ${}^{18}O_2$ (80) mbar). If the light-induced degradation results in the formation of water molecules one would expect to observe vibrational modes of $H_2^{18}O$, that are shifted in frequency due to the heavier ¹⁸O atoms. However, the IR spectra do not show an isotope shift of the water lines (figure 5c lower panel). Hence, this experiment demonstrates that water is not a reaction product of the light-induced degradation process. Instead, perovskite samples already contain H_2O , most likely in form of crystallization water. ^[71,73,78] To show that the atmosphere did contain ¹⁸O the vibrational modes of CO_2 are depicted in the inset. The modes consist of $C^{16}O_2$ and $C^{18}O_2$. Hence, the data presented in figure 5c clearly demonstrate that oxygen acts as a catalyst only.

Molecular oxygen is known to diffuse within MAPbI₃. ^[178] Its charge state resides about 0.2 eV below the conduction band of MAPbI₃. ^[179] Under charge injection or illumination with $h\nu \ge E_{\rm G}$, free charge carriers are generated within MAPbI₃. ^[180] Consequently, an O₂ molecule can capture a free electron from the conduction band, forming superoxide in the ${}^{2}\Pi_{\rm g}$ configuration (equation 2). The O₂⁻ and CH₃NH₃⁺ ions are a Brønsted-Lowry base and acid, respectively. Therefore, in a further reaction step CH₃NH₃⁺ is deprotonated (see equation 3). The produced methylamine diffuses out of the sample and hence, impedes any back-reaction. In a final step, the produced HO₂ molecules dissociate into molecular hydrogen and oxygen and the H₂ molecules diffuse out of the sample. Then, the remaining O₂ molecules can act as catalyst for the next dissociation process (equation 4).

$$O_2 \xrightarrow[MAPbI_3]{h\nu \ge E_G} O_2^-$$
(2)

$$CH_3NH_3^+ + O_2^- \longrightarrow HO_2 + CH_3NH_2 \uparrow$$
(3)

$$2 HO_2 \longrightarrow 2O_2 + H_2 \uparrow \tag{4}$$

3.3.2. Photo-dissociation

Several groups observed significant photo-instabilities of MAPbI₃.^[55,56,58,173,181–186] However, most experiments were performed in air. Thus, the dominant degradation mechanism is a consequence of the combined action of oxygen and light. A few groups studied the effect of light soaking of MAPbI₃ in inert atmosphere. Astonishingly, a strong increase in photoluminescence intensity^[178,185,187], open circuit voltage, and short circuit current ^[188] was shown. As result, power conversion efficiencies often increase within the first seconds to minutes upon illumination. ^[56,172,188] Recently, this so called *photo-curing* behavior has been associated with a reduction of the density of trapping sites from initially $\approx 10^{17}$ cm⁻³ to about 10^{16} cm⁻³. ^[187] It was suggested that illumination annihilates iodine vacancy/interstitial Frenkel pairs. ^[187] These defects are relatively common in MAPbI₃ and are shown to diffuse easily.^[187,189] Huang et al.^[172] further proved that the positive effect of photo curing-vanishes upon storage in the dark. In fact, the efficiency η decreases to less than 10 % of the maximal value after 12 hours in darkness. Upon light exposure, the efficiency subsequently recovers repeatedly. Interestingly, a photo-induced degradation was not observed in this study. Moreover, Ummadisingu et al. ^[190] reported on a light-activated crystallization of MAPbI₃.

Conversely, several other groups report on the commencement of a photo-induced degradation after several minutes of illumination ^[56,158,181], as shown in **figure 6**a. ^[56] Within the first minutes of illumination, the authors observe an increase in η . Then however, light-induced degradation commences. When the solar cells are kept at V = V_{OC} (condition A) Nie et al. ^[56] show a 30 % degradation over the course of 2 hrs under illumination. Near infrared absorption indicates the formation of deep trapping levels. ^[56] Interestingly, the degradation decelerates when the devices are operated at V = 0 V (condition B). The solar cells recover in the dark, as the majority of trapping states dissipates. This recovery is strongly temperature dependent. Figure 6b depicts the degradation and self-healing scheme. Within a microscopic picture, it is suggested that localized polarons form upon illumination. They reside deep

within the band gap and act as trapping states. The corresponding structural distortions are corroborated by DFT calculations ^[56] and an increase of distinct Raman modes ^[56,191]. A deeper insight into the governing mechanism can be found elsewhere. ^[56]

Domanski et al.^[158] observed a very similar photo-induced degradation while operating devices at the maximum power point. This is depicted in figure 6c. In general, they differentiate two distinct degradation mechanisms: a swift degradation within the first hours and a somewhat slower mechanism. They show that the fast degradation process is fully reversible upon storing devices in the dark. This is identical to the study by Nie et al.^[56] (figure 6a). Conversely, Domanski et al. ^[158] attribute the effect to the migration of ionic vacancies. As depicted in figure 6d, they suggest an initially random distribution of anion and cation vacancies, which migrate and form Debye layers at the charge selective contacts under working conditions. The generated interface charge causes a band bending, which facilitates or hinders charge extraction.^[192] This induces reversible performance losses. A strong manifestation of this effect has been observed on devices following the layer stack glass/ITO/PEDOT:PSS/CH₃NH₃PbI₃/MoO_x/Al. There, pre-biasing allows Zhao et al. ^[193] to switch the polarity. In general, ion migration is widely accepted in organic-inorganic perovskites. ^[194,195] Activation energies for migration are in the range of 0.1 to 0.8 eV in MAPbI₃. ^[182] This is a low value range even compared to oxide ion conductors used in fuel cells, such as LaFeO₃ where typical activation energies are reported in the range of 0.5 to 2 eV^[196]. The migration of ions or ionic defects is believed to be responsible for slow photoconductivity responses ^[182], solar cell hysteresis ^[197] or switchable photocurrents. ^[66] A detailed report on ion migration in MAPbI₃ can be found elsewhere. ^[194] It must be noted, however, that Domanski et al. ^[158] investigated a mixed cation, mixed halide (FAPbI₃)_x(MAPbBr₃)_y perovskite. Such alloyed perovskites are prone to photo-induced phase segregation, which might affect the reported results. This behavior is discussed in more detail in section 3.3.3.

Well-encapsulated MAPbI₃ layer were recently studied under intense illumination by steady state and time resolved PL measurements. Most notably, Chen et al. ^[198], find a swift quenching of the PL intensity, which is strongly dependent on the initial defect density. Photo-induced phase segregation, discussed in section 3.3.3, is also retarded in perovskites that exhibit a long-range crystalline order.^[199–201] It is likely that this behavior is related to a suppressed ionic vacancy migration in absence of grain boundaries. ^[194] Certainly, this can also be related to a complex interplay of pre-existing and light-induced defects. In this context

it is important to note that electric noise spectroscopy by Landi et al.^[202] suggests that the capture of charge carriers takes place in two steps: First electrons/holes are trapped by shallow defects, and second they either are re-emitted or thermalize down to a deeper state, where they recombine subsequently.

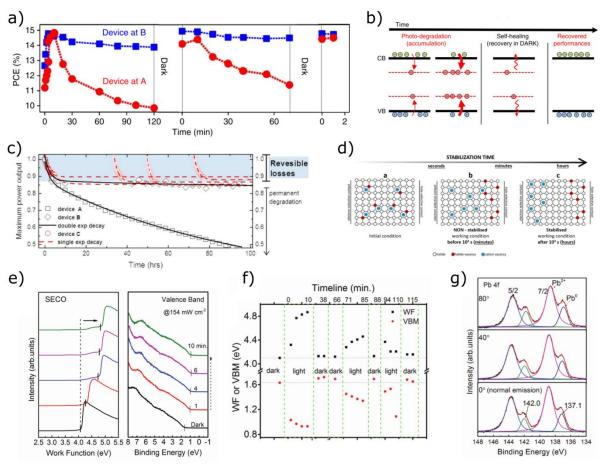


Figure 6. Photo degradation of organic-inorganic perovskites: (a) Power conversion efficiency of ITO/PEDOT:PSS/MAPbI₃/PCBM/Al devices under constant illumination. Devices were stressed either at J = 0 and $V = V_{OC}$ (condition A) or $J = J_{SC}$, V = 0 (condition B). (b) Proposed mechanism of photo-degradation and self-healing processes. (a, b) Reproduced with permission. ^[56] Copyright 2016, Nature

(c) Maximum power-point tracking of FTO/TiO₂/mp-TiO₂/(FAPbI₃)_x(MAPbBr₃)_y/spiro-OMeTAD/Au devices under continuous UV filtered LED illumination. Devices A and B were continuously tracked, while device C was cyclically tracked for 5 hours and then stored in the dark. (d) Schematic depiction of the evolution of the ion distribution within the perovskite layer. (c, d) Reproduced with permission. ^[158] Copyright 2017, Wiley. (e) Time dependent UPS spectra of MAPbI_{3-x}Cl_x thin films under white LED illumination (154 mW cm⁻²) (f) Evolution of the work function and VBM over multiple on-off illumination cycles. (g) Angle dependent XPS spectra of the thin film after multiple on-off illumination cycles. (e, f, g) Reproduced with permission. ^[173] Copyright 2016, Wiley.

As indicated before, Domanski et al. observed a second somewhat slower mechanism, which leads to permanent degradation. As the kinetics of the slow process vary vastly between

identically prepared devices (device A & B in figure 6c), they propose that one or more device components are involved.

Recently, Zu et al. ^[173] observed permanent degradation on MAPbI_{3-x}Cl_x films and MAPbI₃ single crystals. They perform ultraviolet photoelectron spectroscopy (UPS) measurements prior to, during, and after exposure to white light. A gradual shift of the valence band maximum (VBM) and the work function due to surface photovoltage under illumination was reported (see figure 6e). For low photon doses, these shifts are, as expected, reversible. However, for high photon doses of about 154 Wcm⁻² and long-term illumination the surface photovoltage vanishes. Using angle dependent XPS they identify the formation of Pb⁽⁰⁾ after intense prolonged illumination. The authors suggested that the formed Pb⁽⁰⁾ induces a strong pinning of the Fermi energy close to the conduction band maximum. Simultaneously, a different absorption was observed that was attributed to PbI₂, which formed during the illumination in high vacuum.

It appears that light induced degradation of MAPbI₃ based solar cells involves two distinct processes: a swift but reversible degradation process and a slower degradation process that causes irreversible changes of the samples. The latter irreversible degradation process has not been subject to investigations so far. To elucidate the irreversible degradation of organic inorganic perovskites light-induced gas-effusion experiments have been performed. ^[155] For this purpose polycrystalline MAPbI₃ thin films were placed into a gas effusion setup with a base-pressure of 10⁻⁸ mbar. The residual gas composition was measured as a function of time using a quadrupole mass spectrometer. Figure 7a depicts the ion current as a function of time for H₂, HD, CH₂, NH₂, CH₃NH₂, and CH₃NH₃; D represents the hydrogen isotope deuterium (²H). Without illumination, the ion currents are constant. After 6200 s a blue LED with an excitation energy of hv = 2.72 eV and a photon flux of $F \approx 3 \times 10^{17}$ s⁻¹cm⁻² is switched on (see red arrow). After a short delay of about 200 s the ion currents of all monitored masses increase. It has to be noted, that all masses reach a maximum ion current at $t \approx 7860$ s except for CH₃NH₃, which exhibits its maximum ion current already at $t \approx 6880$ s. Light-induced gas effusion experiments were repeated using an ultraviolet LED ($hv = 3.4 \text{ eV}, F \approx 5 \times 10^{17} \text{ s}^{-1} \text{ cm}^{-1}$ ²) and a red LED ($hv = 3.4 \text{ eV}, F \approx 5 \times 10^{17} \text{ s}^{-1} \text{cm}^{-2}$). The obtained results are shown in figures 7b and 7c. For clarity ion currents for HD, CH₂, and NH₂ are omitted. Their time dependence is the same as for the other molecules.

When illuminating with UV or blue light all ion currents increase. However, in case of red light the ion currents do not increase with time. Hence, the excitation energy for the light induced degradation of MAPbI₃ has to be equal to or exceed 2.72 eV. The observation of H₂ and HD out-diffusion from the perovskite layer upon illumination implies that N – H/D bonds dissociate. Further insight into the microscopic mechanism can be obtained from the onset of the ion current traces. In particular, the ion current for CH₃NH₃, increases first and decreases to its background value before H₂ and CH₃NH₃ reach their maximum. It is likely that out-diffusion of CH₃NH₃ is limited to a small amount that originates from a surface near region.^[155] The subsequent increase of the ion currents for H₂, HD, and CH₃NH₂ molecules implies that illumination results in the deprotonation of CH₃NH₃ in the perovskite lattice.

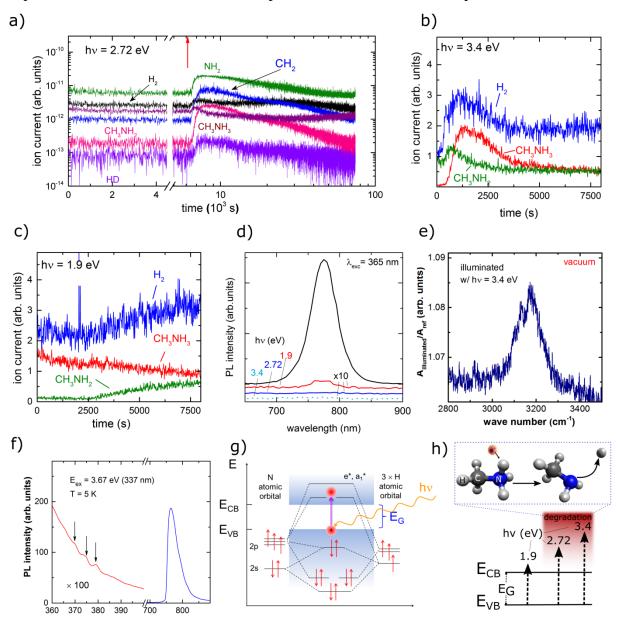


Figure 7. Light-induced dissociation of organic-inorganic perovskites. (a) Ion current of H₂, HD, CH₂, NH₂, CH₃NH₂, and CH₃NH₃ as a function of time. At t = 6200s the MAPbI₃ sample was illuminated with hv = 2.72 eV and a photon flux $F \approx 3 \times 10^{17}$ s⁻¹cm⁻². (b) Ion current as a function of time under illumination with UV light; hv = 3.4 eV, $F \approx 5 \times 10^{17}$ s⁻¹cm⁻². (c) Ion current as a function of time under illumination with red light; hv = 1.9 eV, $F \approx 5 \times 10^{17}$ s. (d) Photoluminescence spectra of MAPbI₃ samples before (black curve) and after illumination for about $t = 7.4 \times 10^4$ s with a photon energy of hv = 1.9 eV (red curve), 2.72 eV (blue curve), and 3.4 eV (dashed curve). (e) FT-IR spectra of MAPbI₃ after exposure to UV light (hv = 3.4 eV, $t_{ill} = 2.82 \times 10^4$ s) in vacuum. The data are normalized to the untreated specimen prior to illumination. (f) Photoluminescence spectrum is magnified by a factor of 100. (g) Schematic depiction of the molecular orbitals of ammonia in respect to the perovskite conduction and valence bands. (h) Schematic depiction of the dissociation mechanism of CH₃NH₃⁺ due to the capture of a photo-generated electron. (a, d, g, and h) Reproduced with permission. ^[155] Copyright 2017, John Wiley and Sons.

The proposed photo-induced deprotonation was corroborated using FT-IR measurements before, during, and after illumination in vacuum.^[155] Figure 7e depicts the ratio of the illuminated to the pristine unexposed sample, A/A_{ref}. The observed change in A/A_{ref} points upwards. This corresponds to a decrease of amine concentration in the lattice after illumination. This is consistent with the light-induced effusion measurements showing that illumination results in the dissociation of N – H bonds. Consequently, the opto-electronic properties of MAPbI₃ are deteriorated vastly. Figure 7d depicts PL measurements before and after prolonged illumination. Note, that the yield of the PL intensity is inversely proportional to the concentration of localized defects.^[203,204] Already after illumination with *hv* = 1.9 eV for *t* = 7.4×10^4 s in vacuum, the PL yield decreased significantly. However, as soon as the photon energy exceeds 2.7 eV a vast number of localized defects is created, which results in a complete quenching of the band-to-band photoluminescence. Most likely, the localized defects are fragments of the methylammonium ions that have energy levels in the band gap of the perovskite. ^[205,206]

Considering the observed photo-dissociation of N – H bonds in MAPbI₃ a microscopic mechanism was proposed. ^[155] Upon illumination with light of $h\nu \ge 2.7$ eV the dissociation of N – H bonds occurs according to the reaction:

$$2(CH_3NH_3^+) + 2e^- \to 2(CH_3NH_2) + H_2 \qquad . \tag{5}$$

This is schematically depicted in figure 7h. In general, photo-dissociation of NH₃ can occur if an electron is placed in one of the antibonding orbitals. Figure 7g illustrates the molecular orbitals of NH₃ according to molecular orbital theory. Most notably, the antibonding molecular orbitals e^* and a_1^* , are high in energy and unoccupied.^[155] It is interesting to note,

that photo-dissociation of isolated NH₃ requires an energy of around 6.4 eV.^[207] This is far beyond the photon energies used above and usually takes place in Jovian's atmosphere only.^[208] In the hybrid semiconductor molecular orbitals from CH₃NH₃⁺ form energy bands which reside deep in the valence and conduction bands. *Ab-initio* methods predicted a minimal energy of 8 eV, which is needed for a direct transition from the valence band maximum into empty conduction bands associated to CH₃NH₃⁺. ^[209,210] In principle such a transition should *not* occur, as excited electrons usually thermalize down to the conduction band-edge within 10⁻¹¹ s. However, it has been shown recently, that the hot carrier lifetime in MAPbI₃ is about 3 orders of magnitude higher as compared to conventional semiconductors. ^[211,212]

Figure 7f shows photoluminescence measurements of MAPbI₃ taken at 5 K. In addition to the band gap luminescence at 1.62 eV, Nickel et al ^[155] observe three hitherto unknown PL peaks located at 3.27, 3.31, and 3.34 eV. It was concluded, that these states correspond to the antibonding orbitals of NH₃. As they give rise to luminescence, it is clear that they are *not* in resonance with the conduction bands of the crystal lattice of the MAPbI₃. Therefore, the data suggest that the dissociation of $CH_3NH_3^+$ commences by trapping charge in the antibonding states of N – H, which requires a minimal energy of about 2.72 eV.

Light-induced degradation of perovskites in the absence of O_2 , Ti O_2 , moisture, and heat is broadly discussed in the literature. While not all reported changes of device parameters and materials properties are supported with exact models the data presented in figure 7 establish that deprotonation of the organic cation occurs for photon energies ≥ 2.7 eV. In fact, it is conceivable that even hot carriers with sufficient energy can be trapped by the N-H antibonding states, which ultimately results in the dissociation of the cation. This seems highly relevant for degradation due to high-energy irradiation as described in chapter 3.4.

3.3.3 Photo-Induced Phase Segregation of Alloyed Perovskites

Alloying ABX₃ perovskites proved to be the route to tailor opto-electronic properties and in particular the band gap. Monovalent cations $[CH_3NH_3^+, HC(NH_2)_2^+, Cs^+, or Rb^+]$, divalent cations $[Pb^{2+} \text{ or } Sn^{2+}]$ and halide anions $[Br^- \text{ or } I^-]$ can be used for A, B and X respectively. Combining halides, for example, allows a continuous tuning of the band-gap energy between

1.5 and 2.2 eV or between 1.8 and 2.4 eV for MAPb($I_{1-x}Br_x$)₃ ^[42] and CsPb($I_{1-x}Br_x$)₃ ^[169], respectively. This is an important prerequisite for tandem solar cell optimization. However, the majority of these alloys are not stable under illumination. Photoluminescence measurements reported by Hoke et al. (**figure 8**a), reveal a swift rise of a low-energy emission peak under continuous illumination for MAPb($I_{1-x}Br_x$)₃ based alloys. ^[213] A photo-induced phase segregation into iodide rich and bromide rich domains, frequently referred to as the Hoke effect, ^[199,214] has been suggested. ^[213] This was corroborated by XRD measurements that revealed the coexistence of both phases after illumination. As a result, photo-induced charge carriers are trapped in the low-energy iodide rich phase. Interestingly, the effect is fully reversible upon thermal annealing in the dark.^[213] Figure 8b depicts the temporal evolution of the photo-induced phase segregation on the nanoscale. Bischak et al.^[215] overlaid secondary electron (SE) and a cathodoluminescence (CL) images of MAPb($I_{0.9}Br_{0.1}$)₃ before and after illumination. The data show that photo-induced iodine rich clusters grow quickly in size and somewhat slower in number.

The driving force for the photo-induced phase segregation is subject of ongoing research.^[199,213–216] It has been proposed, that the coexistence of both phases corresponds to a local minimum of the Helmholtz free energy. ^[217] Remarkably, photo-induced phase segregation does not occur when a pulsed excitation is used with a repetition rate of less than 1 kHz. This effect is independent of the excitation energy.^[216] This indicates that the phase segregation is governed by a two-step process. ^[216] To account for the degradation a model has been suggested that postulates the existence of long-lived polaronic states that drive ion migration locally.^[215,216]

As a result, iodine rich clusters form that have a lower band gap. These clusters are the origin of the low-energy PL emission line. It is interesting to note that segregation of iodine into nanodomains also was observed in highly efficient (FAPbI₃)_{0.85}(MAPbBr₃)_{0.15} solar cells. ^[218]

As a matter of fact, photo-induced phase segregation is not restricted to MAPb(I_{1-x}Br_x)₃ based alloys. It has also been reported for CsPb(I_{1-x}Br_x)₃, FAPb(I_{1-x}Br_x)₃, (MA,FA)Pb(I_{1-x}Br_x)₃, (Cs_xFA_{1-x})PbI₃, and (Cs_{1-y}FA_y)Pb(I_{1-x}Br_x)₃ alloys. As example figure 8c depicts X-ray diffractograms of (Cs_{1-x}FA_x)PbI₃ before and after irradiation. ^[219] In case of (Cs_{0.3}FA_{0.7})PbI₃ a splitting of the XRD reflex of the 110 planes at 2 $\Theta \approx 14^{\circ}$ indicates the presence of two coexisting phases after irradiation. Recently, Unger et al. ^[214] published an overview of the Hoke effect using all available compositional data from literature.^{[40-42,103,146,169,200,213,216,220-^{230]} As a general trend, substitution of iodine with bromine forms materials with a higher band}

gap. This is shown in figure 8d where the initial and final PL emission after light soaking is depicted. Above a certain threshold all alloys exhibit a lower energy PL line. The onset is located at ≈ 1.7 eV and ≈ 1.9 eV for MAPb(I_{1-x}Br_x)₃ and CsPb(I_{1-x}Br_x)₃, respectively. In addition, the formation of a lower energy emission line is accompanied by a pinning of the open circuit voltage of perovskite solar cells. This is observed for all alloys with an absorption onset above the aforementioned threshold (see figure 8e). ^[214]

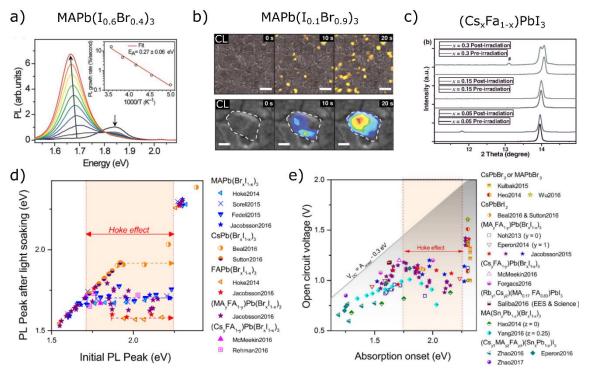


Figure 8. Photo-induced phase segregation of mixed perovskites. (a) Photoluminescence spectra of MAPb($I_{0.6}Br_{0.4}$)₃ under illumination (15 mWcm⁻², 457 nm, for 45 s). Reproduced with permission. ^[213] Copyright 2015, The Royal Society of Chemistry. (b) Cathodoluminescence image series after light soaking using a 405 nm LED with a power density of 50 mWcm⁻². Scale bars are 2 µm (top) and 200 nm (bottom) Reproduced with permission. ^[215] Copyright 2017, American Chemical Society. (c) X-ray diffractogram of (Cs_xFA_{1-x})PbI₃ before and after irradiation. Reproduced with permission. ^[219] Copyright 2017, John Wiley and Sons.(d) PL emission after light soaking as a function of initial PL peak energy (e) Open circuit voltage of reported perovskite solar cells as a function of perovskite band gap. (d, e) Reproduced with permission. ^[214] Copyright 2015, The Royal Society of Chemistry.

The photo-induced phase segregation limits the performance of higher band gap perovskite solar-cells fundamentally. Apparently, a continuous tuning of the band-gap energy is not achievable. However, a band gap of 1.73 eV, which is needed for efficient tandem solar-cells with crystalline Si ^[46], is achievable using, for example, CsPb(I_{1-x}Br_x)₃ or (Cs_{1-y}FA_y)Pb(I_{1-x}Br_x)₃ alloys. ^[146,169,200]

The photo-induced phase segregation is retarded in perovskites that exhibit a long-range crystalline order.^[199–201] However, the question arises whether light-induced defect generation and/or the presence of heat and moisture etc. will affect the crystalline structure, thereby promoting phase segregation.

3.4. High Energy Irradiation

In addition to the degradation mechanisms discussed above the irradiation with high-energy rays and particles, such as X-rays, γ -rays, electrons, neutrons or protons, can result in detrimental effects on all known inorganic and organic materials. On the other hand, various characterization methods rely on the interaction of a probe beam with the specimen. Any degradation of the specimen during a measurement alters the physical properties of the sample and may produces artifacts that lead to unreliable information and subsequently to a misinterpretation of the data.

Rutherford backscattering spectroscopy, for example, is widely used to obtain depth dependent concentration profiles. ^[231] Yet the irradiation of a MAPbI₃ thin-films with ⁴He⁺ ions of an energy of 1.4 MeV causes a swift degradation to PbI₂.^[232] Apparently, two or more complementary techniques with orthogonal analysis beam direction have to be used to ensure correct measurements.^[232] The chemical structure of hybrid perovskites is commonly determined from X-ray photoelectron spectroscopy (XPS) data. ^[69,76,87,88,233–245] As example, figure 9a depicts XPS spectra of the Pb4f levels of MAPbI₃, MAPbI_{3-x}Cl_x, MAPbCl₃, and PbCl₂ thin films. ^[246] The X-ray energy used for data acquisition amounts to 4 keV. All spectra are dominated by two intense peaks, which can be assigned to the Pb4 $f_{5/2}$ and Pb4 $f_{7/2}$ levels of Pb²⁺. At lower binding energies of 137.0 eV and 141.85 eV all perovskite thin-films exhibit two additional peaks that are reflecting the presence of Pb⁽⁰⁾. On the one hand, Pb⁽⁰⁾ is often observed in XPS spectra of hybrid perovskites throughout literature. [69,247,248] On the other hand, metallic Pb⁽⁰⁾ is not found in spectra of lead halides (figure 9a). Clearly, metallic Pb⁽⁰⁾ can be a remnant of an improper preparation and/or annealing procedure. ^[249] However, it has been shown that the peak attributed to metallic Pb⁽⁰⁾ increases in intensity with prolonged X-ray exposure. ^[246] By contrast, a variety of reports use thick MAPbI₃ ^[250] films, or single crystals of MAPbI₃^[251], MAPbBr₃^[252] or Cs_xFA_{1-x}PbI_{3-y}Br_y^[253] to detect γ- and Xrays. Perovskite single crystals combine high stopping powers for X- and γ -rays with a large

charge-carrier mobility-lifetime product. Hence, very high sensitivities of 80 μ C Gy_{air}⁻¹cm⁻² can be achieved. ^[252] This sensitivity is four times higher than the sensitivity of commercially used α -Se X-ray detectors^[252], and might allow medical diagnostics to work with much lower doses. However, none of the reports published so far paid attention to the long-term stability of perovskite based γ - and X-ray detectors. ^[254]

Other common tools to study the microstructure and composition of hybrid perovskite films on the nanoscale are electron microscopy and energy dispersive X-ray spectroscopy. Both techniques are widely used but only a few studies paid attention to perovskite degradation during the exposure to low-energy electrons. ^[255–258] For example, Edri et al. ^[256] measured electron beam induced current (EBIC) spectra on the cross section of MAPbI₃ based solar cells. They observed a strong fading of the EBIC contrast after only 3 consecutive measurements.^[256] This low-energy electron-beam-induced degradation of MAPbI₃ and FAPbI₃ was studied in depth using cathodoluminescence (CL) spectroscopy. ^[90] Xia et al. ^[90] showed that an intense electron irradiation with energies ranging from 2 to 10 keV can alter the properties of the perovskites. CL spectra measured with an excitation of 5 keV for a variety of electron currents and irradiation times are depicted in figure 9b. The observed changes of the CL spectra are ascribed to two different mechanisms: (i) defect formation caused by radiation damage, and (ii) electron-beam induced heating in case of elevated beam currents. The former mechanism causes quenching and broadening of the excitonic peaks in the CL spectra, whereas the latter one results in the formation of new peaks with higher emission photon energies (see figure 9b).^[90]

A further increase of the excitation energy to values in the MeV regime alters the interaction entirely. The penetration depth increases to values in the μ m regime and beyond. Nevertheless, electron and proton irradiation in the MeV regime deteriorate inorganic solar cells and electronic devices based on Si, InGaP, GaAs, or Ge vastly. ^[259–267] Astonishingly, recent findings revealed an unexpectedly high radiation hardness of MAPbI₃ based solar cells under high-energetic electron and proton irradiation. ^[89,91,92] Figure 9c shows *in-situ* measurements of the fill factor, FF, short circuit current, *J*_{SC}, open circuit voltage, *V*_{OC}, and power conversion efficiency, η , of a perovskite solar cell as a function of the proton dose, ϕ . The protons had a kinetic energy of 68 MeV. For comparison, the blue line in figure 9c shows *J*_{SC} of a single crystal silicon (c-Si) diode. Interestingly, for the MAPbI₃ based solar cell FF and V_{OC} remain unchanged over the whole course of the proton irradiation. Only *J*_{SC} decreases to

25

a values 0.6 for a proton dose of $\phi = 10^{13}$ p cm⁻². Note that J_{SC} of the c-Si diode decreases to a value of less than 0.6 for a lower proton dose of $\phi = 10^{12}$ p cm⁻². For the perovskite solar cell, it was shown that a major part of the reduction of J_{SC} is due to shading caused by the formation of color centers in the glass substrate. ^[89] It was concluded that the radiation hardness of MAPbI₃ exceeds the radiation hardness of c-Si by almost three orders of magnitude. ^[89]

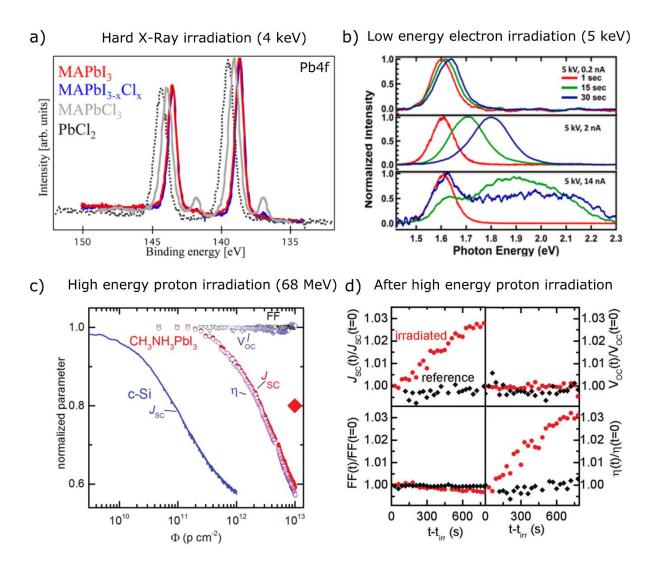


Figure 9. Degradation of hybrid perovskites due to high-energy irradiation. a) Hard X-ray photoelectron spectra of the Pb4f levels of MAPbI₃, MAPbI_{3-x}Cl_x, MAPbCl₃ and PbCl₂ thin films on glass/FTO/TiO₂ substrates using a photon energy of 4 keV. Reproduced with permission. ^[246] Copyright 2015, American Chemical Society. b) Cathodoluminescence spectra of MAPbI₃ thin films under excitation with high-energy electrons ($E_e = 5 \text{ kV}$). The irradiation times and electron currents are indicated in the figure. Reproduced with permission. ^[90] Copyright 2015, American Chemical Society. c) Normalized solar cell parameters of a CH₃NH₃PbI₃ solar cell under high-energy proton irradiation ($E_P = 68 \text{ MeV}$) as a function of the proton dose, ϕ . The red diamond depicts the normalized value of *J*_{SC} after correcting for shading losses due to the creation of color centers in the substrate. For comparison, the

evolution of a c-Si photo-diode is shown by the blue solid curve. d) Self-healing of a protonirradiated perovskite solar cell (red circles) after terminating the proton irradiation. Data points are normalized to unity at $t - t_{irr} = 0$. (c, d) Reproduced with permission. ^[89] Copyright 2016, John Wiley and Sons.

After termination of the proton irradiation the MAPbI₃ based solar cells were stored in the dark for 2 weeks to allow the induced radioactivity to decrease. After that time the solar cells show a partial recover of J_{SC} and η (see figure 9d). In fact, it has been shown that the density of recombination centers in the MAPbI₃ absorber layer is significantly reduced after these 2 weeks. ^[92] As a result, V_{OC}, FF, and the recombination lifetime of photo-generated charge carries are higher in proton-irradiated devices compared to non-irradiated reference devices. ^[89,92]

The question arises if the self-healing is caused by the same mechanism that is responsible for the recovery of perovskite solar cells after prolonged illumination with photons in the UV/Vis energy range (see figure 6a and 6c in chapter 3.3.2). It is known that high-energy electron and proton irradiation generates cascades of secondary electrons with high energies. It is conceivable, that the secondary electrons can cause the breaking of N-H and C-H bonds of the organic cations prior to thermalization to the conduction band-edge.^[89] The dissociation of C-H bonds upon irradiation is well-known in organic semiconductors such as PCDTBT or P3HT. ^[268,269] The governing microscopic mechanism could be similar to the light-induced dissociation of N-H bonds observed for MAPbI₃ and FAPbI₃ (see chapter 3.3.2 and figure 7). In fact, fragmentation of CH₃NH₃⁺ gives rise to localized states in the band gap of MAPbI₃ and therefore, might account for the degradation of solar cells due to proton irradiation.^[205,206] The subsequent self-healing of the solar cells indicates that some of the newly created defects are passivated or neutralized. This may occur as a back reaction of the defect creation process and/or the neutralization of defects due to the formation of neutral complexes e.g.: the capture of a proton.

The superior radiation hardness of MAPbI₃ may lead to practical applications. As an example, it is conceivable to integrate perovskite solar cells in tandem devices with radiation hard CIGS [copper indium gallium di-selenide] devices for space applications. Such tandem solar cells would be lightweight, thin, and flexible and have a high power conversion efficiency. Another example would be radiation hard sensors for the nuclear energy sector, medical imaging, and radiotherapy treatments. Even applications for low-energy electron radiation are

conceivable despite of the reported swift degradation of the perovskite (see figure 9b), since low-energy particles can be shielded with a few millimeters of material.

4. Beyond Organic-Inorganic Perovskites

As discussed above, hybrid perovskites are suffering from different types of degradation mechanisms. On the one hand, the temperature and humidity induced degradation pathways can be easily overcome. The former can be avoided by introducing the more robust formamidinium cation into the perovskite lattice, and the latter by introducing different encapsulation layers. On the other hand, photo-induced dissociation of the organic cation is a major drawback (see section 3.3.2). In fact, any photo-instability may impede commercialization of this promising material.

In general, there are various routes to stabilize hybrid perovskites. For example, one could replace the organic cation with a more stable cation. This is possible using cesium based perovskites, following the general formula $C_{s}BX_{3}$, were X = I, Br and B = Pb, Sn. Interestingly, all inorganic ABX₃ perovskites were introduced even before the organic/inorganic prototypes were available. Early works by Wells and Haupt report on the synthesis and the properties of Cesium, Potassium, and Rubidium lead halide perovskites.^[270,271] Nowadays, Cesium containing perovskites are successfully implemented in solar cells. Moreover, cesium lead trihalides have demonstrated a number of properties similar to those of organic/inorganic perovskites. In these perovskites the band gap can be tuned over a wide range. They have a high tolerance for defects, show a high charge carrier mobility, and a long carrier diffusion length. ^[228,272] Moreover, the inorganic perovskites can be operated at elevated temperatures and do not exhibit light induced degradation. For example, nanocrystals of CsPbX₃ perovskite show an extreme luminescence quantum yield at room temperature combined with a broadly tunable band gap (Figure 10a).^[273,274] In addition, CsPbI₃ QDs exhibit strong quantum confinement and allow to obtain solar cells with a stabilized power conversion efficiency of $\eta \approx 8$ % .^[275] Recently, Kulbak et al. demonstrated solar cells with open circuit voltages as high as 1.3 V using bulk CsPbBr₃ layers that were sandwiched between FTO/TiO₂ and a hole-transport-layer/Au contacts (figure 10b).^[224] Subsequently, several groups reported on efficient cesium based perovskite solar-cells. ^[228,276,277] Using CsPbI₃, the best performing devices had an efficiency of more than 9 %. ^[228] Unfortunately, none of the previously reported materials of the APbX₃ structure form a thermodynamically stable perovskite phase, which is photoactive and typically called "black

phase". The reason for this instability was found in the slightly smaller tolerance factor, due to the mismatch of the ionic radius of the inorganic A-site (Cs, Rb, K) cation within the lead iodide octahedral structure.

Nevertheless, early work on inorganic perovskites by Wells and Haupt revealed that iodoplumbates of Cs, Rb, and potassium form stable materials that adopt a different stoichiometry, namely $A_m B_n X_{2n+m}$, where m = 2 and n = 1. Latest results of computational screening revealed that potential candidates for optoelectronic devices could be found among a broader range of stoichiometric ratio - $A_m B_n X_y$. Moreover, hypothetical $A_m B_n X_y$ perovskites have less restrictions regarding the valences of the metallic cation, namely $B = B^{2+}$ (Sn, Ge), B^{3+} (Bi, Sb), or B^{4+} (Figure 10c).^[278,279]

Recently, inorganic alloys Ag₂PbI₄ and AgBi₂I₇ were for the first time implemented into devices and yielded an efficiency of $\eta = 3.9$ and 1.22 %, respectively (Figure 10 (d)).^[280,281] A further alternative approach to obtain stable perovskites is the heterovalent substitution of the B-site cation with B¹⁺ and B³⁺, which allows to obtain "double perovskite" structures.^[282–284] The material Cs₂AgBiX₆ shows a moderately long recombination lifetime of $\tau \ge 660$ ns and a band gap of $E_{\rm G} \sim 1.95$ eV for Cs₂AgBiBr₆ and $E_{\rm G} \sim 2.7$ eV was observed for Cs₂AgBiI₆.^[285,286]

Recent computational research to design and identify new potential hybrid perovskites is based mainly on quaternary ammonium compounds.^[287,288] However, N – H bonds play a major role in the light-induced degradation mechanism of organic/inorganic perovskites (see chapter 3.3.2). Therefore, we would like to point out that for the design of stable hybrid perovskites organic cations might have to be considered that do not contain weak bonds. Altogether, these interesting and diverse approaches have the potential to overcome the deficiencies described above and lead to new perovskites with a long-time stability and the potential for high efficiency solar cells.

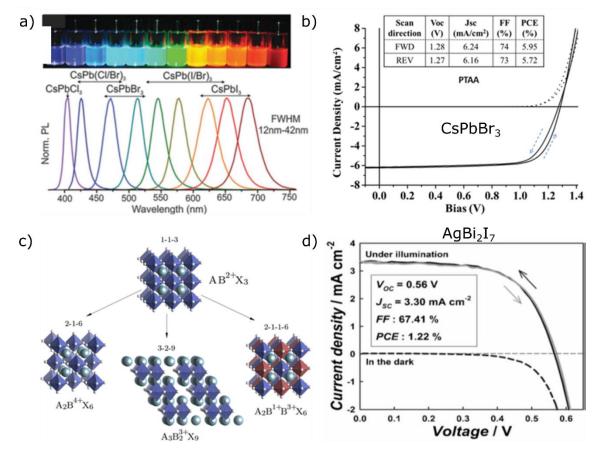


Figure 9. Beyond organic-inorganic perovskites: a) Colloidal perovskite CsPbX₃ nanocrystals (X = Cl, Br, I) in toluene under UV excitation and corresponding PL spectra. Reproduced with permission.^[274] Copyright 2015, American Chemical Society. b) Light and dark J-V curves of mp-TiO₂/CsPbBr₃ solar cells with PTAA as a hole transport layer (100 mW/cm² simulated solar irradiation, masked cell area 0.16 cm², scan rate 0.06 V/s). Reproduced with permission.^[224] Copyright 2015, American Chemical Society c) Schematic depiction of the relation between the crystal structures of Pb perovskites and A_mB_nX_y perovskites, where B-site cation has the different oxidation states, and double perovskites. Reproduced with permission.^[284] Copyright 2016, American Chemical Society d) J-V curves of amp-TiO₂/AgBi₂I₇/P3HT/Ag solar cell in the dark and under illumination with a light density of 100 mWcm² (area = 0.049 cm2). Reproduced with permission.^[281] Copyright 2016, John Wiley and Sons.

5. Conclusion

In conclusion, we thoroughly analyzed T_{80} lifetimes of various perovskite solar cells by comparing over 50 publications reported in literature. Often lifetimes are given for storing devices in the dark and a considerable reduction of the lifetime is reported under illumination conditions. However, in the past years considerable progress was made regarding the stabilities of solar cells. The known mechanisms that are responsible for the degradation of perovskite solar cells were reviewed and their impact was evaluated. Among the degradation mechanisms are the influence of moisture, heat, photo-induced phase segregation in mixed

cation mixed halide perovskites, oxygen catalyzed degradation, photo-induced dissociation of organic cations, and high-energy irradiation. Alternative approaches to overcome materials and device instabilities were suggested. They comprise the use of inorganic perovskites, crystal structures of the type $A_m B_n X_y$, double perovskites, and approaches to eliminate weak bonds in the organic cations.

Acknowledgements

The authors would like to thank E. Unger, S. Albrecht, C. Hartmann and A. Abate for helpful discussions. One of the authors, V.V.B., acknowledges financial support from the Alexander-von-Humboldt Foundation. O. S. is pleased to acknowledge financial support from the graduate school 'Hybrid4Energy'.

Received: ((will be filled in by the editorial staff)) Revised: ((will be filled in by the editorial staff)) Published online: ((will be filled in by the editorial staff))

References

- [1] C. C. Stoumpos, C. D. Malliakas, M. G. Kanatzidis, *Inorg. Chem.* 2013, 52, 9019.
- [2] G. Xing, N. Mathews, S. S. Lim, N. Yantara, X. Liu, D. Sabba, M. Grätzel, S. Mhaisalkar, T. C. Sum, *Nat. Mater.* 2014, *13*, 476.
- S. D. Stranks, G. E. Eperon, G. Grancini, C. Menelaou, M. J. P. Alcocer, T. Leijtens, L.
 M. Herz, A. Petrozza, H. J. Snaith, *Science* 2013, *342*, 341.
- [4] T. Dittrich, F. Lang, O. Shargaieva, J. Rappich, N. H. Nickel, E. Unger, B. Rech, *Appl. Phys. Lett.* 2016, *109*, 73901.
- [5] Y.-H. Kim, H. Cho, J. H. Heo, T.-S. Kim, N. Myoung, C.-L. Lee, S. H. Im, T.-W. Lee, *Adv. Mater.* 2015, 27, 1248.
- [6] S. A. Veldhuis, P. P. Boix, N. Yantara, M. Li, T. C. Sum, N. Mathews, S. G. Mhaisalkar, *Adv. Mater.* 2016, 28, 6804.
- Z.-K. Tan, R. S. Moghaddam, M. L. Lai, P. Docampo, R. Higler, F. Deschler, M. Price,
 A. Sadhanala, L. M. Pazos, D. Credgington, F. Hanusch, T. Bein, H. J. Snaith, R. H.
 Friend, *Nat. Nanotechnol.* 2014, 9, 687.
- [8] O. A. Jaramillo-Quintero, R. S. Sanchez, M. Rincon, I. Mora-Sero, J. Phys. Chem. Lett. 2015, 6, 1883.
- [9] D. Liang, Y. Peng, Y. Fu, M. J. Shearer, J. Zhang, J. Zhai, Y. Zhang, R. J. Hamers, T.
 L. Andrew, S. Jin, ACS Nano 2016, 10, 6897.
- [10] J. Li, S. G. R. Bade, X. Shan, Z. Yu, Adv. Mater. 2015, 27, 5196.
- [11] J. Pan, L. N. Quan, Y. Zhao, W. Peng, B. Murali, S. P. Sarmah, M. Yuan, L. Sinatra, N. M. Alyami, J. Liu, E. Yassitepe, Z. Yang, O. Voznyy, R. Comin, M. N. Hedhili, O. F. Mohammed, Z. H. Lu, D. H. Kim, E. H. Sargent, O. M. Bakr, *Adv. Mater.* 2016, 28, 8718.
- [12] Z. Chen, C. Zhang, X.-F. Jiang, M. Liu, R. Xia, T. Shi, D. Chen, Q. Xue, Y.-J. Zhao, S.

Su, H.-L. Yip, Y. Cao, Adv. Mater. 2017, 29, 1603157.

- [13] X. Y. Chin, D. Cortecchia, J. Yin, A. Bruno, C. Soci, Nat. Commun. 2015, 6, 1.
- [14] X. Liu, L. Niu, C. Wu, C. Cong, H. Wang, Q. Zeng, H. He, Q. Fu, W. Fu, T. Yu, C. Jin,
 Z. Liu, T. C. Sum, *Adv. Sci.* 2016, *3*, 1600137.
- [15] T. S. Kao, Y.-H. Chou, C.-H. Chou, F.-C. Chen, T.-C. Lu, *Appl. Phys. Lett.* 2014, 105, 231108.
- [16] R. Dhanker, A. N. Brigeman, A. V. Larsen, R. J. Stewart, J. B. Asbury, N. C. Giebink, *Appl. Phys. Lett.* 2014, 105, 151112.
- [17] Y. Fu, H. Zhu, A. W. Schrader, D. Liang, Q. Ding, P. Joshi, L. Hwang, X.-Y. Y. Zhu,
 S. Jin, *Nano Lett.* 2016, *16*, 1000.
- [18] S. Yakunin, L. Protesescu, F. Krieg, M. I. Bodnarchuk, G. Nedelcu, M. Humer, G. De Luca, M. Fiebig, W. Heiss, M. V. Kovalenko, *Nat. Commun.* 2015, 6, 8056.
- [19] F. Deschler, M. Price, S. Pathak, L. E. Klintberg, D.-D. Jarausch, R. Higler, S. Hüttner,
 T. Leijtens, S. D. Stranks, H. J. Snaith, M. Atatüre, R. T. Phillips, R. H. Friend, *J. Phys. Chem. Lett.* 2014, *5*, 1421.
- [20] B. R. Sutherland, S. Hoogland, M. M. Adachi, C. T. O. Wong, E. H. Sargent, ACS Nano 2014, 8, 10947.
- [21] S. Chen, K. Roh, J. Lee, W. K. Chong, Y. Lu, N. Mathews, T. C. Sum, A. Nurmikko, ACS Nano 2016, 10, 3959.
- [22] Y. Fu, H. Zhu, C. C. Stoumpos, Q. Ding, J. Wang, M. G. Kanatzidis, X. Zhu, S. Jin, ACS Nano 2016, 10, 7963.
- [23] M. Saliba, S. M. Wood, J. B. Patel, P. K. Nayak, J. Huang, J. A. Alexander-Webber, B. Wenger, S. D. Stranks, M. T. Hörantner, J. T.-W. Wang, R. J. Nicholas, L. M. Herz, M. B. Johnston, S. M. Morris, H. J. Snaith, M. K. Riede, *Adv. Mater.* 2016, 28, 923.
- [24] K. C. Kwon, K. Hong, Q. Van Le, S. Y. Lee, J. Choi, K.-B. Kim, S. Y. Kim, H. W. Jang, Adv. Funct. Mater. 2016, 26, 4213.

- [25] K. Domanski, W. Tress, T. Moehl, M. Saliba, M. K. Nazeeruddin, M. Grätzel, Adv. Funct. Mater. 2015, 25, 6936.
- [26] Y. Lee, J. Kwon, E. Hwang, C.-H. Ra, W. J. Yoo, J.-H. Ahn, J. H. Park, J. H. Cho, Adv. Mater. 2015, 27, 41.
- [27] B. R. Sutherland, A. K. Johnston, A. H. Ip, J. Xu, V. Adinolfi, P. Kanjanaboos, E. H. Sargent, ACS Photonics 2015, 2, 1117.
- [28] Y. Fang, Q. Dong, Y. Shao, Y. Yuan, J. Huang, *Nat. Photonics* **2015**, *9*, 679.
- [29] C. Liu, K. Wang, P. Du, E. Wang, X. Gong, A. J. Heeger, *Nanoscale* 2015, 7, 16460.
- [30] Z. Lian, Q. Yan, Q. Lv, Y. Wang, L. Liu, L. Zhang, S. Pan, Q. Li, L. Wang, J.-L. L. Sun, Sci. Rep. 2015, 5, 16563.
- [31] R. Dong, Y. Fang, J. Chae, J. Dai, Z. Xiao, Q. Dong, Y. Yuan, A. Centrone, X. C. Zeng,
 J. Huang, *Adv. Mater.* 2015, 27, 1912.
- [32] M. He, Y. Chen, H. Liu, J. Wang, X. Fang, Z. Liang, Chem. Commun. 2015, 51, 9659.
- [33] V. Adinolfi, O. Ouellette, M. I. Saidaminov, G. Walters, A. L. Abdelhady, O. M. Bakr,E. H. Sargent, *Adv. Mater.* 2016, 28, 7264.
- [34] S. Chen, C. Teng, M. Zhang, Y. Li, D. Xie, G. Shi, Adv. Mater. 2016, 28, 5969.
- [35] Y. Wang, Z. Xia, S. Du, F. Yuan, Z. Li, Z. Li, Q. Dai, H. Wang, S. Luo, S. Zhang, H.
 Zhou, *Nanotechnology* 2016, 27, 175201.
- [36] C. Gao, J. Liu, C. Liao, Y. Q. Ye, Y. Zhang, X. He, X. Guo, J. Mei, L. W. M. Lau, *RSC Adv.* 2015, 5, 26175.
- [37] C. Liu, K. Wang, P. Du, E. Wang, X. Gong, A. J. Heeger, *Nanoscale* 2015, 7, 16460.
- [38] H. Rao, W. Li, B. Chen, D. Kuang, C. Su, Adv. Mater. 2017, 29, 1602639.
- [39] Y. Lee, J. Kwon, E. Hwang, C.-H. Ra, W. J. Yoo, J.-H. Ahn, J. H. Park, J. H. Cho, Adv. Mater. 2015, 27, 41.
- [40] M. Saliba, T. Matsui, K. Domanski, J.-Y. Seo, A. Ummadisingu, S. M. Zakeeruddin,J.-P. Correa-Baena, W. R. Tress, A. Abate, A. Hagfeldt, M. Grätzel, *Science* 2016, 354,

206.

- [41] M. Saliba, T. Matsui, J.-Y. Seo, K. Domanski, J.-P. Correa-Baena, M. K. Nazeeruddin, S. M. Zakeeruddin, W. Tress, A. Abate, A. Hagfeldt, M. Grätzel, *Energy Environ. Sci.* 2016, *9*, 1989.
- [42] J. H. Noh, S. H. Im, J. H. Heo, T. N. Mandal, S. Il Seok, Nano Lett. 2013, 13, 1764.
- [43] S. Albrecht, M. Saliba, J. P. Correa Baena, F. Lang, L. Kegelmann, M. Mews, L. Steier,
 A. Abate, J. Rappich, L. Korte, R. Schlatmann, M. K. Nazeeruddin, A. Hagfeldt, M.
 Grätzel, B. Rech, *Energy Environ. Sci.* 2016, 9, 81.
- [44] F. Lang, M. A. Gluba, S. Albrecht, J. Rappich, L. Korte, B. Rech, N. H. Nickel, J. Phys. Chem. Lett. 2015, 6, 2745.
- [45] K. A. Bush, A. F. Palmstrom, Z. J. Yu, M. Boccard, R. Cheacharoen, J. P. Mailoa, D. P. McMeekin, R. L. Z. Hoye, C. D. Bailie, T. Leijtens, I. M. Peters, M. C. Minichetti, N. Rolston, R. Prasanna, S. Sofia, D. Harwood, W. Ma, F. Moghadam, H. J. Snaith, T. Buonassisi, Z. C. Holman, S. F. Bent, M. D. McGehee, *Nat. Energy* 2017, *2*, 17009.
- [46] S. Albrecht, M. Saliba, J.-P. Correa-Baena, K. Jäger, L. Korte, A. Hagfeldt, M. Grätzel,
 B. Rech, J. Opt. 2016, 18, 64012.
- [47] R. G. Wilks, M. Bär, *Nat. Energy* **2017**, *2*, 16204.
- [48] S. Yang, W. Fu, Z. Zhang, H. Chen, C.-Z. Li, J. Mater. Chem. A 2017, 5, 11462.
- [49] A. H. Slavney, R. W. Smaha, I. C. Smith, A. Jaffe, D. Umeyama, H. I. Karunadasa, *Inorg. Chem.* 2017, 56, 46.
- [50] X. Zhao, N.-G. Park, *Photonics* **2015**, *2*, 1139.
- [51] J. Schoonman, Chem. Phys. Lett. 2015, 619, 193.
- [52] J. Yang, T. L. Kelly, *Inorg. Chem.* **2017**, *56*, 92.
- [53] N. Aristidou, I. Sanchez-Molina, T. Chotchuangchutchaval, M. Brown, L. Martinez, T. Rath, S. A. Haque, *Angew. Chemie Int. Ed.* 2015, *54*, 8208.
- [54] W.-J. Yin, T. Shi, Y. Yan, Appl. Phys. Lett. 2014, 104, 63903.

- [55] A. Merdasa, M. Bag, Y. Tian, E. Källman, A. Dobrovolsky, I. G. Scheblykin, J. Phys. Chem. C 2016, 120, 10711.
- [56] W. Nie, J. C. Blancon, A. J. Neukirch, K. Appavoo, H. Tsai, M. Chhowalla, M. A. Alam, M. Y. Sfeir, C. Katan, J. Even, S. Tretiak, J. J. Crochet, G. Gupta, A. D. Mohite, *Nat Commun* 2016, 7, 11574.
- [57] D. Bryant, N. Aristidou, S. Pont, I. Sanchez-Molina, T. Chotchunangatchaval, S.Wheeler, J. R. Durrant, S. A. Haque, *Energy Environ. Sci.* 2016, *9*, 1655.
- [58] G. Murugadoss, S. Tanaka, G. Mizuta, S. Kanaya, H. Nishino, T. Umeyama, H. Imahori, S. Ito, *Jpn. J. Appl. Phys.* 2015, *54*, 08KF08.
- [59] A. Dualeh, P. Gao, S. Il Seok, M. K. Nazeeruddin, M. Grätzel, *Chem. Mater.* 2014, 26, 6160.
- [60] R. K. Misra, S. Aharon, B. Li, D. Mogilyansky, I. Visoly-Fisher, L. Etgar, E. A. Katz, J. Phys. Chem. Lett. 2015, 6, 326.
- [61] B. Conings, J. Drijkoningen, N. Gauquelin, A. Babayigit, J. D'Haen, L. D'Olieslaeger,
 A. Ethirajan, J. Verbeeck, J. Manca, E. Mosconi, F. De Angelis, H.-G. Boyen, *Adv. Energy Mater.* 2015, *5*, 1500477.
- [62] D. P. Nenon, J. A. Christians, L. M. Wheeler, J. L. Blackburn, E. M. Sanehira, B. Dou,
 M. L. Olsen, K. Zhu, J. J. Berry, J. M. Luther, *Energy Environ. Sci.* 2016, 9, 2072.
- [63] Y.-B. Cheng, Y. Han, S. Meyer, Y. Dkhissi, K. Weber, J. M. J. Pringle, U. Bach, L. Spiccia, Y.-B. Cheng, J. Mater. Chem. A 2015, 3, 8139.
- [64] Z. Song, S. C. Watthage, A. B. Phillips, B. L. Tompkins, R. J. Ellingson, M. J. Heben, *Chem. Mater.* 2015, 27, 4612.
- [65] T. Supasai, N. Rujisamphan, K. Ullrich, A. Chemseddine, T. Dittrich, *Appl. Phys. Lett.***2013**, *103*, 183906.
- [66] Y. Yuan, J. Chae, Y. Shao, Q. Wang, Z. Xiao, A. Centrone, J. Huang, Adv. Energy Mater. 2015, 5, 1500615.

- [67] Q. Jeangros, M. Duchamp, J. Werner, M. Kruth, R. E. Dunin-Borkowski, B. Niesen, C. Ballif, A. Hessler-Wyser, *Nano Lett.* 2016, *16*, 7013.
- [68] Y. Yuan, Q. Wang, Y. Shao, H. Lu, T. Li, A. Gruverman, J. Huang, *Adv. Energy Mater.* **2016**, *6*, 1501803.
- [69] W. Huang, J. S. Manser, P. V. Kamat, S. Ptasinska, Chem. Mater. 2016, 28, 303.
- [70] X. Dong, X. Fang, M. Lv, B. Lin, S. Zhang, J. Ding, N. Yuan, J. Mater. Chem. A 2015, 3, 5360.
- [71] X. Gong, M. Li, X.-B. Shi, H. Ma, Z.-K. Wang, L.-S. Liao, Adv. Funct. Mater. 2015, 25, 6671.
- [72] W. Xiang, Q. Chen, Y. Wang, M. Liu, F. Huang, T. Bu, T. Wang, Y.-B. Cheng, X.
 Gong, J. Zhong, P. Liu, X. Yao, X. Zhao, *J. Mater. Chem. A* 2017, *5*, 5486.
- [73] E. Mosconi, J. M. Azpiroz, F. De Angelis, Chem. Mater. 2015, 27, 4885.
- [74] L. Zhang, P. H.-L. Sit, J. Phys. Chem. C 2015, 119, 22370.
- [75] C. Müller, T. Glaser, M. Plogmeyer, M. Sendner, S. Döring, A. A. Bakulin, C. Brzuska,
 R. Scheer, M. S. Pshenichnikov, W. Kowalsky, A. Pucci, R. Lovrinčić, *Chem. Mater.*2015, 27, 7835.
- [76] Y. Li, X. Xu, C. Wang, C. Wang, F. Xie, J. Yang, Y. Gao, J. Phys. Chem. C 2015, 119, 23996.
- [77] Z. Song, A. Abate, S. C. Watthage, G. K. Liyanage, A. B. Phillips, U. Steiner, M. Graetzel, M. J. Heben, *Adv. Energy Mater.* 2016, *6*, 1600846.
- [78] W. Zhou, Y. Zhao, C. Shi, H. Huang, J. Wei, R. Fu, K. Liu, D. Yu, Q. Zhao, J. Phys. Chem. C 2016, 120, 4759.
- [79] C. Qin, T. Matsushima, T. Fujihara, W. J. Potscavage, C. Adachi, *Adv. Mater.* 2016, 28, 466.
- [80] J.-W. Lee, D.-H. Kim, H.-S. Kim, S.-W. Seo, S. M. Cho, N.-G. Park, Adv. Energy Mater. 2015, 5, 1501310.

- [81] G. E. Eperon, S. N. Habisreutinger, T. Leijtens, B. J. Bruijnaers, J. J. van Franeker, D.
 W. DeQuilettes, S. Pathak, R. J. Sutton, G. Grancini, D. S. Ginger, R. A. J. Janssen, A.
 Petrozza, H. J. Snaith, ACS Nano 2015, 9, 9380.
- [82] C.-J. Tong, W. Geng, Z.-K. Tang, C.-Y. Yam, X. Fan, J. Liu, W.-M. Lau, L.-M. Liu, J. Phys. Chem. Lett. 2015, 6, 3289.
- [83] I. Hwang, I. Jeong, J. Lee, M. J. Ko, K. Yong, ACS Appl. Mater. Interfaces 2015, 7, 17330.
- [84] Q. Jiang, D. Rebollar, J. Gong, E. L. Piacentino, C. Zheng, T. Xu, *Angew. Chemie Int.Ed.* 2015, 54, 7617.
- [85] J. Yang, B. D. Siempelkamp, D. Liu, T. L. Kelly, ACS Nano 2015, 9, 1955.
- [86] K. K. Bass, R. E. McAnally, S. Zhou, P. I. Djurovich, M. E. Thompson, B. C. Melot, *Chem. Commun.* 2014, 50, 15819.
- [87] A. Alberti, I. Deretzis, G. Pellegrino, C. Bongiorno, E. Smecca, G. Mannino, F. Giannazzo, G. G. Condorelli, N. Sakai, T. Miyasaka, C. Spinella, A. La Magna, *ChemPhysChem* 2015, 16, 3064.
- [88] Y. Li, X. Xu, C. Wang, C. Wang, F. Xie, J. Yang, Y. Gao, J. Phys. Chem. C 2015, 119, 23996.
- [89] F. Lang, N. H. Nickel, J. Bundesmann, S. Seidel, A. Denker, S. Albrecht, V. V. Brus, J. Rappich, B. Rech, G. Landi, H. C. Neitzert, *Adv. Mater.* 2016, 28, 8726.
- C. Xiao, Z. Li, H. Guthrey, J. Moseley, Y. Yang, S. Wozny, H. Moutinho, B. To, J. J.
 Berry, B. Gorman, Y. Yan, K. Zhu, M. Al-Jassim, J. Phys. Chem. C 2015, 119, 26904.
- [91] Y. Miyazawa, M. Ikegami, T. Miyasaka, T. Ohshima, M. Imaizumi, K. Hirose, in 2015 IEEE 42nd Photovolt. Spec. Conf., IEEE, 2015, pp. 1–4.
- [92] V. V Brus, F. Lang, J. Bundesmann, S. Seidel, A. Denker, B. Rech, G. Landi, H. C. Neitzert, J. Rappich, N. H. Nickel, *Adv. Electron. Mater.* 2017, *3*, 1600438.
- [93] T. A. Berhe, W.-N. Su, C.-H. Chen, C.-J. Pan, J.-H. Cheng, H.-M. Chen, M.-C. Tsai,

L.-Y. Chen, A. A. Dubale, B.-J. Hwang, Energy Environ. Sci. 2016, 9, 323.

- [94] J.-H. Im, C.-R. Lee, J.-W. Lee, S.-W. Park, N.-G. Park, Nanoscale 2011, 3, 4088.
- [95] H.-S. Kim, C.-R. Lee, J.-H. Im, K.-B. Lee, T. Moehl, A. Marchioro, S.-J. Moon, R. Humphry-Baker, J.-H. Yum, J. E. Moser, M. Grätzel, N.-G. Park, *Sci. Rep.* 2012, 2, 591.
- [96] Y. S. Kwon, J. Lim, H.-J. Yun, Y.-H. Kim, T. Park, Energy Environ. Sci. 2014, 7, 1454.
- [97] L. Chen, X. Xie, Z. Liu, E.-C. Lee, J. Mater. Chem. A 2017, 5, 6974.
- [98] W. Yan, Y. Li, Y. Li, S. Ye, Z. Liu, S. Wang, Z. Bian, C. Huang, Nano Res. 2015, 8, 2474.
- [99] J. Ye, X. Zhang, L. Zhu, H. Zheng, G. Liu, H. Wang, T. Hayat, X. Pan, S. Dai, Sustain. Energy Fuels 2017, 1, 907.
- [100] J. H. Heo, M. H. Lee, H. J. Han, B. R. Patil, J. S. Yu, S. H. Im, J. Mater. Chem. A 2016, 4, 1572.
- [101] J. H. Heo, S. H. Im, D. H. Song, H. J. Han, S. Y. Kim, J. H. Kim, D. Kim, H. W. Shin,
 T. K. Ahn, C. Wolf, T.-W. Lee, S. H. Im, *Adv. Mater.* 2015, *27*, 3424.
- [102] L. Zheng, Y.-H. Chung, Y. Ma, L. Zhang, L. Xiao, Z. Chen, S. Wang, B. Qu, Q. Gong, *Chem. Commun.* 2014, 50, 11196.
- [103] J. Zhao, K. O. Brinkmann, T. Hu, N. Pourdavoud, T. Becker, T. Gahlmann, R. Heiderhoff, A. Polywka, P. Görrn, Y. Chen, B. Cheng, T. Riedl, *Adv. Energy Mater*. 2017, 1602599.
- [104] H. Choi, C.-K. Mai, H.-B. Kim, J. Jeong, S. Song, G. C. Bazan, J. Y. Kim, A. J. Heeger, *Nat. Commun.* 2015, 6, 7348.
- [105] Q. Tai, P. You, H. Sang, Z. Liu, C. Hu, H. L. W. Chan, F. Yan, *Nat. Commun.* 2016, 7, 11105.
- [106] S. N. Habisreutinger, T. Leijtens, G. E. Eperon, S. D. Stranks, R. J. Nicholas, H. J. Snaith, Nano Lett. 2014, 14, 5561.

- [107] W. Chen, Y. Wu, Y. Yue, J. Liu, W. Zhang, X. Yang, H. Chen, E. Bi, I. Ashraful, M. Grätzel, L. Han, *Science* 2015, 350, 944.
- [108] M. Zhang, M. Lyu, H. Yu, J.-H. Yun, Q. Wang, L. Wang, Chem. A Eur. J. 2015, 21, 434.
- [109] Y. Hou, H. Zhang, W. Chen, S. Chen, C. O. R. Quiroz, H. Azimi, A. Osvet, G. J. Matt,
 E. Zeira, J. Seuring, N. Kausch-Busies, W. Lövenich, C. J. Brabec, *Adv. Energy Mater.*2015, *5*, 1500543.
- [110] H. Back, G. Kim, J. Kim, J. Kong, T. K. Kim, H. Kang, H. Kim, J. Lee, S. Lee, K. Lee, *Energy Environ. Sci.* 2016, 9, 1258.
- [111] G.-W. Kim, G. Kang, J. Kim, G.-Y. Lee, H. Il Kim, L. Pyeon, J. Lee, T. Park, *Energy Environ. Sci.* 2016, 9, 2326.
- [112] J. A. Christians, P. A. Miranda Herrera, P. V. Kamat, J. Am. Chem. Soc. 2015, 137, 1530.
- [113] Q. Luo, Y. Zhang, C. Liu, J. Li, N. Wang, H. Lin, J. Mater. Chem. A 2015, 3, 1.
- [114] C.-Y. Y. Chang, K.-T. T. Lee, W.-K. K. Huang, H.-Y. Y. Siao, Y.-C. C. Chang, *Chem. Mater.* 2015, 27, 5122.
- [115] C.-G. Wu, C.-H. Chiang, Z.-L. Tseng, M. K. Nazeeruddin, A. Hagfeldt, M. Grätzel, *Energy Environ. Sci.* 2015, 8, 2725.
- [116] X. Zhao, H. Shen, Y. Zhang, X. Li, X. Zhao, M. Tai, J. Li, J. Li, X. Li, H. Lin, ACS Appl. Mater. Interfaces 2016, 8, 7826.
- [117] M. R. Leyden, M. V. Lee, S. R. Raga, Y. Qi, J. Mater. Chem. A 2015, 3, 16097.
- [118] N. Lin, J. Qiao, H. Dong, F. Ma, L. Wang, J. Mater. Chem. A 2015, 3, 22839.
- [119] C.-Y. Chang, Y.-C. Chang, W.-K. Huang, K.-T. Lee, A.-C. Cho, C.-C. Hsu, Chem. Mater. 2015, 27, 7119.
- [120] J. Cao, Y.-M. M. Liu, X. Jing, J. Yin, J. Li, B. Xu, Y.-Z. Z. Tan, N. Zheng, J. Am. Chem. Soc. 2015, 137, 10914.

- [121] A. Mei, X. Li, L. Liu, Z. Ku, T. Liu, Y. Rong, M. Xu, M. Hu, J. Chen, Y. Yang, M. Gratzel, H. Han, *Science* 2014, 345, 295.
- [122] J. Zhao, X. Zheng, Y. Deng, T. Li, Y. Shao, A. Gruverman, J. Shield, J. Huang, *Energy Environ. Sci.* 2016, 9, 3650.
- [123] Q. Wang, C. C. Chueh, M. Eslamian, A. K. Y. Jen, ACS Appl. Mater. Interfaces 2016, 8, 32068.
- [124] J. Ye, H. Zheng, L. Zhu, X. Zhang, L. Jiang, W. Chen, G. Liu, X. Pan, S. Dai, Sol. Energy Mater. Sol. Cells 2017, 160, 60.
- [125] W. Ke, D. Zhao, C. R. Grice, A. J. Cimaroli, G. Fang, Y. Yan, J. Mater. Chem. A 2015, 3, 23888.
- [126] L. K. Ono, S. R. Raga, M. Remeika, A. J. Winchester, A. Gabe, Y. Qi, J. Mater. Chem. A 2015, 3, 15451.
- [127] J. P. Correa Baena, L. Steier, W. Tress, M. Saliba, S. Neutzner, T. Matsui, F. Giordano,
 T. J. Jacobsson, A. R. Srimath Kandada, S. M. Zakeeruddin, A. Petrozza, A. Abate, M.
 K. Nazeeruddin, M. Grätzel, A. Hagfeldt, *Energy Environ. Sci.* 2015, *8*, 2928.
- [128] Z. Zhu, Y. Bai, X. Liu, C.-C. Chueh, S. Yang, A. K.-Y. Jen, Adv. Mater. 2016, 28, 6478.
- [129] Z. Yu, B. Chen, P. Liu, C. Wang, C. Bu, N. Cheng, S. Bai, Y. Yan, X. Zhao, Adv. Funct. Mater. 2016, 26, 4866.
- [130] M. Kaltenbrunner, G. Adam, E. D. Głowacki, M. Drack, R. Schwödiauer, L. Leonat, D. H. Apaydin, H. Groiss, M. C. Scharber, M. S. White, N. S. Sariciftci, S. Bauer, *Nat. Mater.* 2015, 14, 1032.
- [131] X. Bao, Y. Wang, Q. Zhu, N. Wang, D. Zhu, J. Wang, A. Yang, R. Yang, J. Power Sources 2015, 297, 53.
- [132] J. H. Heo, H. J. Han, D. Kim, T. K. Ahn, S. H. Im, Energy Environ. Sci. 2015, 8, 1602.
- [133] J. Song, E. Zheng, J. Bian, X.-F. Wang, W. Tian, Y. Sanehira, T. Miyasaka, J. Mater.

Chem. A **2015**, *3*, 10837.

- [134] S. Guarnera, A. Abate, W. Zhang, J. M. Foster, G. Richardson, A. Petrozza, H. J. Snaith, J. Phys. Chem. Lett. 2015, 6, 432.
- [135] F. K. Aldibaja, L. Badia, E. Mas-Marzá, R. S. Sánchez, E. M. Barea, I. Mora-Sero, J. Mater. Chem. A 2015, 3, 9194.
- [136] J. H. Kim, P.-W. Liang, S. T. Williams, N. Cho, C.-C. Chueh, M. S. Glaz, D. S. Ginger,
 A. K.-Y. Jen, *Adv. Mater.* 2015, 27, 695.
- [137] Y.-Y. Yu, R.-S. Chiang, H.-L. Hsu, C.-C. Yang, C.-P. Chen, Nanoscale 2014, 6, 11403.
- [138] T. Leijtens, G. E. Eperon, S. Pathak, A. Abate, M. M. Lee, H. J. Snaith, *Nat. Commun.***2013**, *4*, 2885.
- [139] Z. Ku, Y. Rong, M. Xu, T. Liu, H. Han, Sci. Rep. 2013, 3, 3132.
- [140] J. Burschka, N. Pellet, S.-J. Moon, R. Humphry-Baker, P. Gao, M. K. Nazeeruddin, M. Grätzel, *Nature* 2013, 499, 316.
- [141] M. O. Reese, S. A. Gevorgyan, M. Jørgensen, E. Bundgaard, S. R. Kurtz, D. S. Ginley, D. C. Olson, M. T. Lloyd, P. Morvillo, E. A. Katz, A. Elschner, O. Haillant, T. R. Currier, V. Shrotriya, M. Hermenau, M. Riede, K. R. Kirov, G. Trimmel, T. Rath, O. Inganäs, F. Zhang, M. Andersson, K. Tvingstedt, M. Lira-Cantu, D. Laird, C. McGuiness, S. (Jimmy) Gowrisanker, M. Pannone, M. Xiao, J. Hauch, R. Steim, D. M. DeLongchamp, R. Rösch, H. Hoppe, N. Espinosa, A. Urbina, G. Yaman-Uzunoglu, J.-B. Bonekamp, A. J. J. M. van Breemen, C. Girotto, E. Voroshazi, F. C. Krebs, *Sol. Energy Mater. Sol. Cells* 2011, 95, 1253.
- [142] NREL, "Research Cell Efficiency Records," can be found under http://www.nrel.gov/ncpv/images/efficiency_chart.jpg, 2016.
- [143] F. Bella, G. Griffini, J.-P. Correa-Baena, G. Saracco, M. Grätzel, A. Hagfeldt, S. Turri,C. Gerbaldi, *Science* 2016, *354*, 203.
- [144] T. Malinauskas, D. Tomkute-Luksiene, R. Sens, M. Daskeviciene, R. Send, H.

Wonneberger, V. Jankauskas, I. Bruder, V. Getautis, *ACS Appl. Mater. Interfaces* **2015**, 7, 11107.

- [145] C. Yi, J. Luo, S. Meloni, A. Boziki, N. Ashari-Astani, C. Grätzel, S. M. Zakeeruddin,
 U. Röthlisberger, M. Grätzel, *Energy Environ. Sci.* 2016, 9, 656.
- [146] D. P. McMeekin, G. Sadoughi, W. Rehman, G. E. Eperon, M. Saliba, M. T. Horantner,
 A. Haghighirad, N. Sakai, L. Korte, B. Rech, M. B. Johnston, L. M. Herz, H. J. Snaith,
 Science 2016, *351*, 151.
- [147] S. S. Shin, E. J. Yeom, W. S. Yang, S. Hur, M. G. Kim, J. Im, J. Seo, J. H. Noh, S. Il Seok, *Science* 2017, 356, 167.
- [148] D. Joly, L. Pellejà, S. Narbey, F. Oswald, J. Chiron, J. N. Clifford, E. Palomares, R. Demadrille, Sci. Rep. 2015, 4, 4033.
- [149] Q. Yu, D. Zhou, Y. Shi, X. Si, Y. Wang, P. Wang, Energy Environ. Sci. 2010, 3, 1722.
- [150] F. Sauvage, S. Chhor, A. Marchioro, J.-E. Moser, M. Graetzel, J. Am. Chem. Soc. 2011, 133, 13103.
- [151] R. Harikisun, H. Desilvestro, Sol. Energy 2011, 85, 1179.
- [152] C. H. Peters, I. T. Sachs-Quintana, J. P. Kastrop, S. Beaupré, M. Leclerc, M. D. McGehee, Adv. Energy Mater. 2011, 1, 491.
- [153] P. Cheng, X. Zhan, Chem. Soc. Rev. 2016, 45, 2544.
- [154] R. Cheacharoen, W. R. Mateker, Q. Zhang, B. Kan, D. Sarkisian, X. Liu, J. A. Love, X. Wan, Y. Chen, T. Q. Nguyen, G. C. Bazan, M. D. McGehee, *Sol. Energy Mater. Sol. Cells* 2017, 161, 368.
- [155] N. H. Nickel, F. Lang, V. V. Brus, O. Shargaieva, J. Rappich, *Adv. Electron. Mater.*2017, under revision.
- [156] W. Chen, Y. Wu, Y. Yue, J. Liu, W. Zhang, X. Yang, H. Chen, E. Bi, I. Ashraful, M. Grätzel, L. Han, *Science* 2015, 350, 944.
- [157] G. Grancini, C. Roldán-Carmona, I. Zimmermann, E. Mosconi, X. Lee, D. Martineau,

S. Narbey, F. Oswald, F. De Angelis, M. Graetzel, M. K. Nazeeruddin, *Nat. Commun.* **2017**, *8*, 15684.

- [158] K. Domanski, B. Roose, T. Matsui, M. Saliba, S.-H. Turren-Cruz, J.-P. Correa-Baena,
 C. R. Carmona, G. Richardson, J. M. Foster, F. De Angelis, J. M. Ball, A. Petrozza, N.
 Mine, M. K. Nazeeruddin, W. Tress, M. Grätzel, U. Steiner, A. Hagfeldt, A. Abate, *Energy Environ. Sci.* 2017, 10, 604.
- [159] T. Matsui, I. Petrikyte, T. Malinauskas, K. Domanski, M. Daskeviciene, M. Steponaitis,
 P. Gratia, W. Tress, J. P. Correa-Baena, A. Abate, A. Hagfeldt, M. Gr??tzel, M. K.
 Nazeeruddin, V. Getautis, M. Saliba, *ChemSusChem* 2016, 9, 2567.
- [160] T. Dittrich, Materials Concepts for Solar Cells, Imperial College Press, London, 2014.
- [161] N. Onoda-Yamamuro, O. Yamamuro, T. Matsuo, H. Suga, J. Phys. Chem. Solids 1992, 53, 277.
- [162] F. Brivio, J. M. Frost, J. M. Skelton, A. J. Jackson, O. J. Weber, M. T. Weller, A. R.
 Goñi, A. M. A. Leguy, P. R. F. Barnes, A. Walsh, *Phys. Rev. B* 2015, *92*, 144308.
- [163] H. Zhang, X. Qiao, Y. Shen, T. Moehl, S. M. Zakeeruddin, M. Grätzel, M. Wang, J. Mater. Chem. A 2015, 3, 11762.
- [164] C. Quarti, E. Mosconi, J. M. Ball, V. D'Innocenzo, C. Tao, S. Pathak, H. J. Snaith, A. Petrozza, F. De Angelis, *Energy Environ. Sci.* 2016, 9, 155.
- [165] T. J. Jacobsson, W. Tress, J.-P. Correa-Baena, T. Edvinsson, A. Hagfeldt, J. Phys. Chem. C 2016, 120, 11382.
- [166] Y. Rong, L. Liu, A. Mei, X. Li, H. Han, Adv. Energy Mater. 2015, 5, 1.
- [167] R. L. Milot, G. E. Eperon, H. J. Snaith, M. B. Johnston, L. M. Herz, *Adv. Funct. Mater.* **2015**, *25*, 6218.
- [168] Q. Ma, S. Huang, X. Wen, M. A. Green, A. W. Y. Ho-Baillie, *Adv. Energy Mater*.**2016**, *6*, 2.
- [169] R. E. Beal, D. J. Slotcavage, T. Leijtens, A. R. Bowring, R. A. Belisle, W. H. Nguyen,

G. F. Burkhard, E. T. Hoke, M. D. McGehee, J. Phys. Chem. Lett. 2016, 746.

- [170] A. Binek, F. C. Hanusch, P. Docampo, T. Bein, J. Phys. Chem. Lett. 2015, 6, 1249.
- [171] N. J. Jeon, J. H. Noh, W. S. Yang, Y. C. Kim, S. Ryu, J. Seo, S. Il Seok, *Nature* 2015, 517, 476.
- [172] F. Huang, L. Jiang, A. R. Pascoe, Y. Yan, U. Bach, L. Spiccia, Y. B. Cheng, *Nano Energy* 2016, 27, 509.
- [173] F. Zu, P. Amsalem, I. Salzmann, R. Wang, M. Ralaiarisoa, S. Kowarik, S. Duhm, N. Koch, Adv. Opt. Mater. 2017, 1700139.
- [174] V. Duzhko, V. Y. Timoshenko, F. Koch, T. Dittrich, Phys. Rev. B 2001, 64, 75204.
- [175] J.-P. Correa-Baena, W. Tress, K. Domanski, E. H. Anaraki, S.-H. Turren-Cruz, B.
 Roose, P. P. Boix, M. Grätzel, M. Saliba, A. Abate, A. Hagfeldt, S. M. Zakeeruddin, U.
 Steiner, M. Grätzel, A. Abate, M. K. Nazeeruddin, *Energy Environ. Sci.* 2017, 1, 556.
- [176] K. Jemli, H. Diab, F. Lédée, G. Trippé-Allard, D. Garrot, B. Geffroy, J.-S. Lauret, P. Audebert, E. Deleporte, *Molecules* 2016, 21, 885.
- [177] M. A. Pérez-Osorio, R. L. Milot, M. R. Filip, J. B. Patel, L. M. Herz, M. B. Johnston, F. Giustino, J. Phys. Chem. C 2015, 119, 25703.
- [178] Y. Tian, M. Peter, E. L. Unger, M. Abdellah, K. Zheng, T. T. Pullerits, A. Yartsev, V. Sundstrom, I. G. Scheblykin, V. Sundström, I. G. Scheblykin, *Phys. Chem. Chem. Phys.* 2015, 17, 24978.
- [179] A. J. Bard, R. Parsons, J. Jordan, *Standard Potentials in Aqueous Solution*, Taylor & Francis, **1985**.
- [180] A. Miyata, A. Mitioglu, P. Plochocka, O. Portugall, J. T.-W. Wang, S. D. Stranks, H. J. Snaith, R. J. Nicholas, *Nat. Phys.* 2015, *11*, 582.
- [181] P. H. Joshi, L. Zhang, I. M. Hossain, H. A. Abbas, R. Kottokkaran, S. P. Nehra, M. Dhaka, M. Noack, V. L. Dalal, AIP Adv. 2016, 6, 115114.
- [182] J. M. Azpiroz, E. Mosconi, J. Bisquert, F. De Angelis, Energy Environ. Sci. 2015, 8,

2118.

- [183] F. Matsumoto, S. M. Vorpahl, J. Q. Banks, E. Sengupta, D. S. Ginger, J. Phys. Chem. C 2015, 119, 20810.
- [184] R. Sheng, X. Wen, S. Huang, X. Hao, S. Chen, Y. Jiang, X. Deng, M. A. Green, A. W.Y. Ho-Baillie, *Nanoscale* 2016, 8, 1926.
- [185] J. F. Galisteo-López, M. Anaya, M. E. Calvo, H. Míguez, J. Phys. Chem. Lett. 2015, 6, 2200.
- [186] X. Tang, M. Brandl, B. May, I. Levchuk, Y. Hou, M. Richter, H. Chen, S. Chen, S. Kahmann, A. Osvet, F. Maier, H.-P. Steinrück, R. Hock, G. J. Matt, C. J. Brabec, J. Mater. Chem. A 2016, 4, 15896.
- [187] E. Mosconi, D. Meggiolaro, H. J. Snaith, S. D. Stranks, F. De Angelis, *Energy Environ. Sci.* 2016, 9, 3180.
- [188] C. Zhao, B. Chen, X. Qiao, L. Luan, K. Lu, B. Hu, Adv. Energy Mater. 2015, 5, 1500279.
- [189] E. Mosconi, F. De Angelis, ACS Energy Lett. 2016, 1, 182.
- [190] A. Ummadisingu, L. Steier, J.-Y. Seo, T. Matsui, A. Abate, W. Tress, M. Grätzel, *Nature* 2017, 545, 208.
- [191] R. Gottesman, L. Gouda, B. S. Kalanoor, E. Haltzi, S. Tirosh, E. Rosh-Hodesh, Y.
 Tischler, A. Zaban, C. Quarti, E. Mosconi, F. De Angelis, *J. Phys. Chem. Lett.* 2015, 6, 2332.
- [192] H. min Zhang, C. Liang, Y. Zhao, M. Sun, H. Liu, J. Liang, D. Li, F. Zhang, Z. He, *Phys. Chem. Chem. Phys.* 2015, 17, 9613.
- [193] Y. Zhao, C. Liang, H. Zhang, D. Li, D. Tian, G. Li, X. Jing, W. Zhang, W. Xiao, Q. Liu, F. Zhang, Z. He, *Energy Environ. Sci.* 2015, *8*, 1256.
- [194] Y. Yuan, J. Huang, Acc. Chem. Res. 2016, 49, 286.
- [195] W. Tress, J. Phys. Chem. Lett. 2017, 3106.

- [196] F. H. Taylor, J. Buckeridge, C. R. A. Catlow, Chem. Mater. 2016, 28, 8210.
- [197] W. Tress, N. Marinova, T. Moehl, S. M. Zakeeruddin, N. Mohammad K., M. Grätzel,
 M. K. Nazeeruddin, M. Grätzel, *Energy Environ. Sci.* 2015, *8*, 995.
- [198] S. Chen, X. Wen, S. Huang, F. Huang, Y. Cheng, M. Green, A. Ho-baillie, 2017, 1.
- [199] D. J. Slotcavage, H. I. Karunadasa, M. D. McGehee, ACS Energy Lett. 2016, 1, 1199.
- [200] W. Rehman, D. P. McMeekin, J. B. Patel, R. L. Milot, M. B. Johnston, H. J. Snaith, L. M. Herz, *Energy Environ. Sci.* 2017, *10*, 361.
- [201] M. Hu, C. Bi, Y. Yuan, Y. Bai, J. Huang, Adv. Sci. 2016, 3, 1500301.
- [202] G. Landi, H. C. Neitzert, C. Barone, C. Mauro, F. Lang, S. Albrecht, B. Rech, S. Pagano, Adv. Sci. 2017, accepted.
- [203] T. Dittrich, J. Rappich, V. Y. Timoshenko, Appl. Phys. Lett. 1997, 70, 2705.
- [204] V. Y. Timoshenko, a. B. Petrenko, M. N. Stolyarov, T. Dittrich, W. Fuessel, J. Rappich, J. Appl. Phys. 1999, 85, 4171.
- [205] J. Wang, A. Zhang, J. Yan, D. Li, Y. Chen, Appl. Phys. Lett. 2017, 110, 123903.
- [206] P. Delugas, A. Filippetti, A. Mattoni, Phys. Rev. B 2015, 92, 45301.
- [207] V. M. Donnelly, A. P. Baronavski, J. R. McDonald, Chem. Phys. 1979, 43, 271.
- [208] S. S. Prasad, L. A. Capone, J. Geophys. Res. 1976, 81, 5596.
- [209] Y. H. Chang, C. H. Park, J. Korean Phys. Soc. 2004, 44, 889.
- [210] Y. Wang, T. Gould, J. F. Dobson, H. Zhang, H. Yang, X. Yao, H. Zhao, *Phys. Chem. Chem. Phys.* 2014, 16, 1424.
- [211] Z. Guo, Y. Wan, M. Yang, J. Snaider, K. Zhu, L. Huang, Science 2017, 356, 59.
- [212] Y. Yang, D. P. Ostrowski, R. M. France, K. Zhu, J. van de Lagemaat, J. M. Luther, M. C. Beard, *Nat. Photonics* 2015, *10*, 53.
- [213] E. T. Hoke, D. J. Slotcavage, E. R. Dohner, A. R. Bowring, H. I. Karunadasa, M. D. McGehee, *Chem. Sci.* 2015, 6, 613.
- [214] E. L. Unger, L. Kegelmann, K. Suchan, D. S?rell, L. Korte, S. Albrecht, J. Mater.

Chem. A 2017, 5, 11401.

- [215] C. G. Bischak, C. L. Hetherington, H. Wu, S. Aloni, D. F. Ogletree, D. T. Limmer, N. S. Ginsberg, *Nano Lett.* 2017, *17*, 1028.
- [216] X. Yang, X. Yan, W. Wang, X. Zhu, H. Li, W. Ma, C. X. Sheng, Org. Electron. physics, Mater. Appl. 2016, 34, 79.
- [217] F. Brivio, C. Caetano, A. Walsh, J. Phys. Chem. Lett. 2016, 7, 1083.
- [218] P. Gratia, G. Grancini, J.-N. Audinot, X. Jeanbourquin, E. Mosconi, I. Zimmermann, D. Dowsett, Y. Lee, M. Grätzel, F. De Angelis, K. Sivula, T. Wirtz, M. K. Nazeeruddin, J. Am. Chem. Soc. 2016, 138, 15821.
- [219] J. Huang, P. Xu, J. Liu, X.-Z. You, Small 2017, 13, 1603225.
- [220] D. Forgács, D. Pérez-del-Rey, J. Ávila, C. Momblona, L. Gil-Escrig, B. Dänekamp, M. Sessolo, H. J. Bolink, J. Mater. Chem. A 2017, 5, 3203.
- [221] W. Liao, D. Zhao, Y. Yu, N. Shrestha, K. Ghimire, C. R. Grice, C. Wang, Y. Xiao, A. J. Cimaroli, R. J. Ellingson, N. J. Podraza, K. Zhu, R. G. Xiong, Y. Yan, J. Am. Chem. Soc. 2016, 138, 12360.
- [222] G. E. Eperon, S. D. Stranks, C. Menelaou, M. B. Johnston, L. M. Herz, H. J. Snaith, *Energy Environ. Sci.* 2014, 7, 982.
- [223] Z. Wu, B. Wang, J. He, T. Chen, J. Colloid Interface Sci. 2016, 461, 162.
- [224] M. Kulbak, D. Cahen, G. Hodes, J. Phys. Chem. Lett. 2015, 6, 2452.
- [225] T. J. Jacobsson, L. J. Schwan, M. Ottosson, A. Hagfeldt, T. Edvinsson, *Inorg. Chem.***2015**, *54*, 10678.
- [226] J. T. Jacobsson, J. P. Correa Baena, M. Pazoki, M. Saliba, K. Schenk, M. Grätzel, A. Hagfeldt, *Energy Environ. Sci.* 2016, 41, 1.
- [227] G. E. Eperon, C. E. Beck, H. J. Snaith, Mater. Horiz. 2016, 3, 63.
- [228] R. J. Sutton, G. E. Eperon, L. Miranda, E. S. Parrott, B. A. Kamino, J. B. Patel, M. T.Hörantner, M. B. Johnston, A. A. Haghighirad, D. T. Moore, H. J. Snaith, *Adv. Energy*

Mater. 2016, 6, 1502458.

- [229] P. Fedeli, F. Gazza, D. Calestani, P. Ferro, T. Besagni, A. Zappettini, G. Calestani, E. Marchi, P. Ceroni, R. Mosca, J. Phys. Chem. C 2015, 119, 21304.
- [230] J. H. Heo, D. H. Song, S. H. Im, Adv. Mater. 2014, 26, 8179.
- [231] J. R. Tesmer, M. A. Nastasi, M. Tesmer, J R;Nastasi, Handbook of Modern Ion Beam Analysis, Materials Research Society, Pittsburgh, Pennsilvania, 1995.
- [232] F. Lang, A. Juma, V. Somsongkul, T. Dittrich, M. Arunchaiya, *Hybrid Mater.* 2014, 1, 52.
- [233] B. Li, Y. Chen, Z. Liang, D. Gao, W. Huang, RSC Adv. 2015, 5, 94290.
- [234] D. Wei, T. Wang, J. Ji, M. Li, P. Cui, Y. Li, G. Li, J. M. Mbengue, D. Song, J. Mater. Chem. A 2016, 4, 1991.
- [235] W. Qiu, U. W. Paetzold, R. Gehlhaar, V. Smirnov, H.-G. Boyen, J. G. Tait, B. Conings,
 W. Zhang, C. Nielsen, I. Mcculloch, L. Froyen, P. Heremans, D. Cheyns, *J. Mater. Chem. A* 2015, 22824.
- [236] H. Xie, X. Liu, L. Lyu, D. Niu, Q. Wang, J. Huang, Y. Gao, J. Phys. Chem. C 2016, 120, 215.
- [237] F. Matteocci, Y. Busby, J.-J. Pireaux, G. Divitini, S. Cacovich, C. Ducati, A. Di Carlo, ACS Appl. Mater. Interfaces 2015, 7, 26176.
- [238] A. Calloni, A. Abate, G. Bussetti, G. Berti, R. Yivlialin, F. Ciccacci, L. Duò, J. Phys. Chem. C 2015, 119, 21329.
- [239] H. Xu, Y. Wu, J. Cui, C. Ni, F. Xu, J. Cai, F. Hong, Z. Fang, W. Wang, J. Zhu, L. Wang, R. Xu, F. Xu, Phys. Chem. Chem. Phys. 2016, 18, 18607.
- [240] W. A. Dunlap-Shohl, R. Younts, B. Gautam, K. Gundogdu, D. B. Mitzi, J. Phys. Chem. C 2016, 120, 16437.
- [241] Y. Liu, L. A. Renna, M. Bag, Z. A. Page, P. Kim, J. Choi, T. Emrick, D. Venkataraman,T. P. Russell, ACS Appl. Mater. Interfaces 2016, 8, 7070.

- [242] U. Bansode, R. Naphade, O. Game, S. Agarkar, S. Ogale, J. Phys. Chem. C 2015, 119, 9177.
- [243] P. Luo, Z. Liu, W. Xia, C. Yuan, J. Cheng, Y. Lu, J. Mater. Chem. A 2015, 3, 12443.
- [244] J. A. McLeod, Z. Wu, P. Shen, B. Sun, L. Liu, J. Phys. Chem. Lett. 2014, 5, 2863.
- [245] J. Emara, T. Schnier, N. Pourdavoud, T. Riedl, K. Meerholz, S. Olthof, *Adv. Mater.***2016**, 28, 553.
- [246] B. Philippe, B.-W. Park, R. Lindblad, J. Oscarsson, S. Ahmadi, E. M. J. Johansson, H. Rensmo, *Chem. Mater.* 2015, 27, 1720.
- [247] R. Lindblad, D. Bi, B. Park, J. Oscarsson, M. Gorgoi, H. Siegbahn, M. Odelius, E. M. J. Johansson, H. Rensmo, J. Phys. Chem. Lett. 2014, 5, 648.
- [248] B. Conings, L. Baeten, C. De Dobbelaere, J. D'Haen, J. Manca, H.-G. Boyen, Adv. Mater. 2014, 26, 2041.
- [249] G. Sadoughi, D. E. Starr, E. Handick, S. D. Stranks, M. Gorgoi, R. G. Wilks, M. Bär, H.J. Snaith, ACS Appl. Mater. Interfaces 2015, 7, 13440.
- [250] S. Yakunin, M. Sytnyk, D. Kriegner, S. Shrestha, M. Richter, G. J. Matt, H. Azimi, C. J. Brabec, J. Stangl, M. V. Kovalenko, W. Heiss, *Nat. Photonics* 2015, 9, 444.
- [251] Q. Dong, Y. Fang, Y. Shao, P. Mulligan, J. Qiu, L. Cao, J. Huang, *Science* 2015, 347, 967.
- [252] H. Wei, Y. Fang, P. Mulligan, W. Chuirazzi, H.-H. Fang, C. Wang, B. R. Ecker, Y. Gao, M. A. Loi, L. Cao, J. Huang, *Nat. Photonics* 2016, *10*, 333.
- [253] O. Nazarenko, S. Yakunin, V. Morad, I. Cherniukh, M. V Kovalenko, NPG Asia Mater.2017, 9, e373.
- [254] W. Heiss, C. Brabec, Nat. Photonics 2016, 10, 288.
- [255] Y. Zhou, A. L. Vasiliev, W. Wu, M. Yang, S. Pang, K. Zhu, N. P. Padture, J. Phys. Chem. Lett. 2015, 6, 2292.
- [256] E. Edri, S. Kirmayer, S. Mukhopadhyay, K. Gartsman, G. Hodes, D. Cahen, Nat.

Commun. 2014, 5, 3461.

- [257] H. Yuan, E. Debroye, K. Janssen, H. Naiki, C. Steuwe, G. Lu, M. Moris, E. Orgiu, H. Uji-i, F. De Schryver, P. Samor?, J. Hofkens, M. Roeffaers, *J. Phys. Chem. Lett.* 2016, 7, 561.
- [258] N. Kedem, T. M. Brenner, M. Kulbak, N. Schaefer, S. Levcenko, I. Levine, D. Abou-Ras, G. Hodes, D. Cahen, J. Phys. Chem. Lett. 2015, 6, 2469.
- [259] B. Todd, S. Uznanski, CAS Cern Accel. Sch. Power Convert. 2015, 3, 1.
- [260] J. J. Loferski, P. Rappaport, Phys. Rev. 1958, 111, 432.
- [261] J. R. Srour, C. J. Marshall, P. W. Marshall, IEEE Trans. Nucl. Sci. 2003, 50, 653.
- [262] T. Markvart, J. Mater. Sci. Mater. Electron. 1990, 1, 1.
- [263] E. J. Daly, G. Drolshagen, A. Hilgers, H. D. R. Evans, *Environ. Model. Space-based Appl. Symp. Proc.* 1996, 15.
- [264] H. C. Neitzert, P. Spinillo, S. Bellone, G. D. Licciardi, M. Tucci, F. Roca, L. Gialanella, M. Romano, *Sol. Energy Mater. Sol. Cells* 2004, 83, 435.
- [265] S. R. Messenger, G. P. Summers, E. A. Burke, R. J. Walters, M. A. Xapsos, Prog. Photovoltaics Res. Appl. 2001, 9, 103.
- [266] M. Alurralde, M. J. L. Tamasi, C. J. Bruno, M. G. Martnez Bogado, J. Plá, J. Fernández Vázquez, J. Durán, J. Schuff, A. A. Burlon, P. Stoliar, A. J. Kreiner, Sol. Energy Mater. Sol. Cells 2004, 82, 531.
- [267] A. D. Verkerk, J. K. Rath, R. E. I. Schropp, *Proc. Inorg. Nanostructured Photovoltaics* 2010, 2, 221.
- [268] R. A. Street, J. E. Northrup, B. S. Krusor, *Phys. Rev. B Condens. Matter Mater. Phys.* **2012**, 85, 1.
- [269] F. Bebensee, J. Zhu, J. H. Baricuatro, J. A. Farmer, Y. Bai, H. P. Steinrück, C. T. Campbell, J. M. Gottfried, *Langmuir* 2010, 26, 9632.
- [270] H. L. Wells, Zeitschrift für Anorg. Chemie 1893, 3, 195.

- [271] H. Haupt, F. Huber, H. Preut, Zeitschrift für Anorg. und Allg. Chemie 1974, 408, 209.
- [272] I. Chung, J.-H. Song, J. Im, J. Androulakis, C. D. Malliakas, H. Li, A. J. Freeman, J. T. Kenney, M. G. Kanatzidis, J. Am. Chem. Soc. 2012, 134, 8579.
- [273] S. Bai, Z. Yuan, F. Gao, J. Mater. Chem. C 2016, 4, 3898.
- [274] L. Protesescu, S. Yakunin, M. I. Bodnarchuk, F. Krieg, R. Caputo, C. H. Hendon, R. X. Yang, A. Walsh, M. V. Kovalenko, *Nano Lett.* 2015, *15*, 3692.
- [275] M. C. Tathavadekar, S. A. Agarkar, O. S. Game, U. P. Bansode, S. A. Kulkarni, S. G. Mhaisalkar, S. B. Ogale, *Sol. Energy* 2015, *112*, 12.
- [276] J. Liang, C. Wang, Y. Wang, Z. Xu, Z. Lu, Y. Ma, H. Zhu, Y. Hu, C. Xiao, X. Yi, G.
 Zhu, H. Lv, L. Ma, T. Chen, Z. Tie, Z. Jin, J. Liu, J. Am. Chem. Soc. 2016, 138, 15829.
- [277] L. A. Frolova, D. V. Anokhin, A. A. Piryazev, S. Y. Luchkin, N. N. Dremova, K. J. Stevenson, P. A. Troshin, J. Phys. Chem. Lett. 2017, 8, 67.
- [278] R. E. Brandt, V. Stevanović, D. S. Ginley, T. Buonassisi, MRS Commun. 2015, 5, 1.
- [279] S. Chakraborty, W. Xie, N. Mathews, M. Sherburne, R. Ahuja, M. Asta, S. G. Mhaisalkar, ACS Energy Lett. 2017, 2, 837.
- [280] L. A. Frolova, D. V. Anokhin, A. A. Piryazev, S. Y. Luchkin, N. N. Dremova, P. A. Troshin, J. Phys. Chem. Lett. 2017, 8, 1651.
- [281] Y. Kim, Z. Yang, A. Jain, O. Voznyy, G.-H. Kim, M. Liu, L. N. Quan, F. P. García de Arquer, R. Comin, J. Z. Fan, E. H. Sargent, *Angew. Chemie Int. Ed.* 2016, 55, 9586.
- [282] G. Volonakis, M. R. Filip, A. A. Haghighirad, N. Sakai, B. Wenger, H. J. Snaith, F. Giustino, J. Phys. Chem. Lett. 2016, 7, 1254.
- [283] G. Volonakis, A. A. Haghighirad, R. L. Milot, W. H. Sio, M. R. Filip, B. Wenger, M. B. Johnston, L. M. Herz, H. J. Snaith, F. Giustino, J. Phys. Chem. Lett. 2017, 8, 772.
- [284] F. Giustino, H. J. Snaith, ACS Energy Lett. 2016, 1, 1233.
- [285] A. H. Slavney, T. Hu, A. M. Lindenberg, H. I. Karunadasa, J. Am. Chem. Soc. 2016, 138, 2138.

- [286] E. T. Mcclure, M. R. Ball, W. Windl, P. M. Woodward, Chem. Mater. 2016, 28, 1348.
- [287] W. Li, Z. Wang, F. Deschler, S. Gao, R. H. Friend, A. K. Cheetham, *Nat. Rev. Mater.* 2017, 2, 16099.
- [288] G. Kieslich, S. Sun, A. K. Cheetham, Chem. Sci. 2015, 6, 3430.

((**For Essays, Feature Articles, Progress Reports, and Reviews,** please insert up to three author biographies and photographs here, max. 100 words each))

Author Photograph(s) ((40 mm broad, 50 mm high, gray scale))



Felix Lang received his Bachelor degree in Applied and Engineering Physics in 2012 from the Technische Universität München, while working on organic bulk heterojunction solar cells. He received his Master degree in Physics in 2014 from the Freie Universität Berlin, working with Dr. T. Dittrich on perovskite solar cells. Currently he is a doctoral candidate under supervision of Prof. Dr. N. H Nickel at the Helmholtz-Zentrum Berlin für Materialien und Energie GmbH. His research interests range from functional materials, their physical properties to novel device concepts.



Oleksandra Shargaieva studied organic chemistry at Taras Shevchenko National University of Kyiv. She received her Master degree in Chemistry in 2014, working on supramolecular interactions of polymethine dyes. Currently she is a member of the graduate school Hybrid 4 Energy, working under supervision of Prof. Dr. N. H Nickel at the Helmholtz-Zentrum Berlin für Materialien und Energie GmbH. Her main scientific interests are synthesis, spectroscopy, and application of hybrid perovskites.



Viktor Brus received his PhD in Solid State Electronics from the Institute for Problems of Materials Science of NAS of Ukraine in 2012. During 2012 - 2015 he worked as an assistant professor at Chernivtsi National University and as postdoctoral fellow at University of California in Santa Barbara with Prof. T.-Q. Nguyen. Currently he is a George Forster Postdoctoral Fellow at Helmholtz-Zentrum Berlin für Materialien und Energie GmbH on behalf of the Alexander von Humboldt Foundation in the group of Prof. Dr. N.H. Nickel. His research interests are in the field of materials science and physics of semiconductor devices.



T₈₀ lifetimes of organic-inorganic perovskite solar cells are strongly reduced under illumination. Various degradation mechanisms are therefore discussed throughout literature. Degradation by moisture or heat is well understood and mitigation possible. Photo-induced phase segregation and photo-induced dissociation of the organic cation however remain unsolved. Recent observations enlighten the underlying microscopic mechanisms and may pave the way for stable perovskites.

Keyword

Felix Lang*, Oleksandra Shargaieva, Viktor V. Brus, Heinz C. Neitzert, Jörg Rappich, Norbert H. Nickel

Title Influence of Radiation on the Properties and the Stability of Hybrid Perovskites

ToC figure

